# Winds from pulsars: constraints on particle acceleration mechanism(s) at the termination shock through relativistic MHD simulations

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### Introduction to the problem & outline

Pulsar Wind Nebulae are powered by the spin-down of young neutron stars that lose their energy in the form of a magnetized, relativistic and cold wind. The fact that the wind is cold makes necessary to use numerical models to investigate it.

- How to construct the numerical model of the Pulsar Wind (PW)
- How to obtain synthetic emission maps from simulations
- Simulation overview: from the beginning to the final stage of the evolution
  - Some results from the simulation of the Crab nebula
- Investigating the particle acceleration mechanism(s) at the wind termination shock of the Crab nebula
  - What wisps observations tell us
  - The wisps multiwavelenghts analysis
  - Results

 $\sum M$ 

### How to fix your PW model

Information about the PW properties are obtained following the iterative scheme:



# Chose your parametrization

## **PW** modelization

From theoretical models we know that the main properties of the PW must be:

$$F(r,\theta) \propto \frac{\alpha + (1-\alpha)}{(2-\alpha)r^2} \Rightarrow F(r,\pi/2) \gg F(r,0) \Rightarrow \alpha \ll 1$$
 anisotropy parameter

The wind is striped in an equatorial belt of extension  $\approx 2x($  inclination of the magnetic axis  $\zeta$  )



$$B(r,\theta) \propto \sqrt{\sigma_0} \sin \theta \tanh \left[ b \left( \frac{\pi}{2} - \theta \right) \right]$$

Intensity + striped zone width

 $\alpha, \sigma, b$ 

Free parameters of the model

### How to obtain synthetic emission maps



(2)

(3)

Injecting the emitting particles in the nebula

Particles are injected at the termination shock with power law (in energy) distribution functions and injection energy  $\varepsilon \approx \infty$  (2 families: 1 for radio and 1 for optical/X-ray particles)

Let it evolves

The local particle distribution function in the nebula is evaluated by evolving the distribution function at the TS in time and space, taking into account adiabatic and syncrotron losses. The emitting properties at each point are described by dedicated particle tracers (from Del Zanna:2006 and Camus:2009).

### Compute emission properties

By applying standard formula for synchrotron emission taking into account relativistic effects (Doppler boosting: enanched emission from fluid elements which move towards the observer and vice versa).

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# How to obtain the synthetic Nebula

## Simulation overview

### STEP 0 Initialization of the system @ t=0 yr



1 week with 64 CPUs.

System evolved up to t=t<sub>Crab</sub><sup>≅</sup>950 yr, solving numerically the eqs of the RMHD 2d axysimmetric model + adiabatic EoS with the shockcapturing numerical code ECHO [Del Zanna 2007].



Typical Run:  $N_r = 800, N_{\theta} = 400, 0 \le t \le 1000 \text{ yr}$ 

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# Fix the best set of free parameters

## Results @ 950 yr

Dynamics





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## Investigating the particle acceleration mechanism

RADIO particles ORIGIN

Properties are only due to the fluid and magnetic field structure. see Olmi et al. 2014

### **OBSERVATIONS:**

Wisps

- Wisps are seen at multiwavelengths
- Not coicident locations
  - Different outward velocities, mildly relativistic  $0.1c \lesssim v \lesssim 0.4c$

To have not coincident wisps at different wavelengths particles must have different acceleration mechanisms! Bietenholz et al. 2004 RADIO + OPTICAL WISPS





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### [Olmi et al. 2015] Multi wavelengths wisp analysis

Radio and X-ray families injected at different locations at the TS. Optical emission = mixed origin (radio + X-ray contribution at optical frequencies)

- 1. Uniform injection  $\rightarrow \theta \in [0^{\circ}, 90^{\circ}]$
- 2. Wide equatorial zone  $\rightarrow$  Equatorial:  $\theta \in [20^\circ, 90^\circ]$ Polar:  $\theta \in [0^\circ, 20^\circ]$
- 3. Narrow equatorial zone  $\rightarrow$  Equatorial:  $\theta \in [70^\circ, 90^\circ]$ Polar:  $\theta \in [0^\circ, 70^\circ]$

Wisp profiles extracted as in Schweizer et al 2013 from a 3" slice in the upper hemisphere of intensity maps. Data extracted with monthly frequency during 10 yr. Results shown as plots of wisp local maxima radial positions vs time.



# (1). Uniform injection:







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## (3). Narrow equatorial zone + wide polar cone



# Best case for X-ray wisps: equatorial injection in a narrow region $\rightarrow \theta \in [70^{\circ}, 90^{\circ}]$

In order to have not coincident radio wisps  $\rightarrow$  injection in

complementary zone or a wider equatorial one (difficult to make stronger constraints).

### Possible scenario:

X-ray particles accelerated in a narrow, low σ, equatorial belt via Fermi I. Radio particles accelerated in a wider equatorial zone via driven magnetic reconnection.

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POLAR

 $0.08c \le v \le 0.38c$ 

 $0.1c \leq v \leq 0.4c$ 

From observations

### Summary and conclusions

### GOALS of NUMERICAL MHD 2D AXISYMM. MODELS

- High energy emission well reproduced
- ✓ Analysis of low energy emission → not strong constraints on radio particles nature, but shows that wisps arise as a fluid property
- ✓ Multi band wisps analysis → Best case for X-ray wisps is injection in a narrow equatorial belt (→ Fermi I acceleration?), no strong constraint for radio particles. They can be accelerated in a wider equatorial zone via driven magnetic reconnection.
  OPEN PROBLEMS:
- □ Radio particles origin → no strong constraints are found
  □ Magnetic field compression around polar axis → only due to simulation dimensionality?

### Complete 3D models needed!

First attempt by Porth et al 2014 (but only 70 yr of the Crab evolution are simulated...)

### Summary and conclusions

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 Multi band wisi
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 Radio particles of Magnetic field condition



### **Preliminary!**

Results of complete Crab's 3D simulation with PLUTO<sup>®</sup> are coming soon

