

Magnetar models for GRBs

Maxim Lyutikov (Purdue U)

Magnetar models for GRBs, and more

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Galilei to Cosimo II Medici

- It is impossible to obtain wages from a republic [...] without having duties attached. [...] so long as I am capable of lecturing and serving, no one in the republic can exempt me from duty while I receive pay. I can hope to enjoy these benefits only from an absolute ruler.

Galilei, Opere x, 348 ff

GRBMAG14 - The magnetic fields!

Magnetars (Thompson & Duncan)

- Magnetars are powered by **dissipation of $\sim 10^{15}$ G B-field**

$$E_B \sim 10^{47} b_{15}^2 \text{ erg}$$

- B-field determines the available energy
- Questions:
 - How B-field is generated and evolves in the crust/core
 - Properties of the crust (plastic or brittle deformations) - How crusts gates the flares

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- More like **Pulsar - PWN**, just more powerful

$$L_{\text{dipole}} \sim B^2 R^2 c \left(\frac{R\Omega}{c} \right)^4 = I\Omega\dot{\Omega}$$

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- Formed in
 - core collapse
 - NS-NS merger
 - AIC of a WD
- spindown: $\tau \sim 100 \text{ sec } P_{\text{msec}}^2 b_{15}^{-2}$
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Not very high. Can explode a SN (LeBlanc & Wilson)? Probably not: SN explosions are not magnetically-driven (nu-driven), but jet can be.

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B-field determines the rate of energy release, NOT the energy content (rotational)

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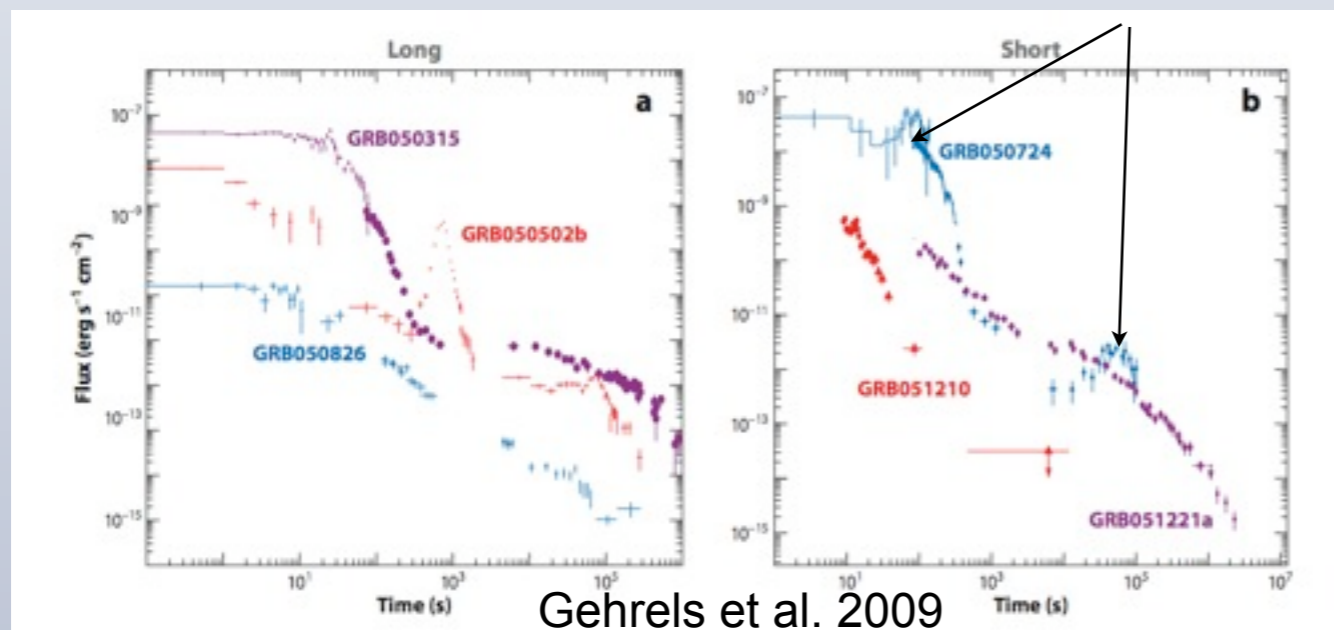
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GRBs: magnetically driven

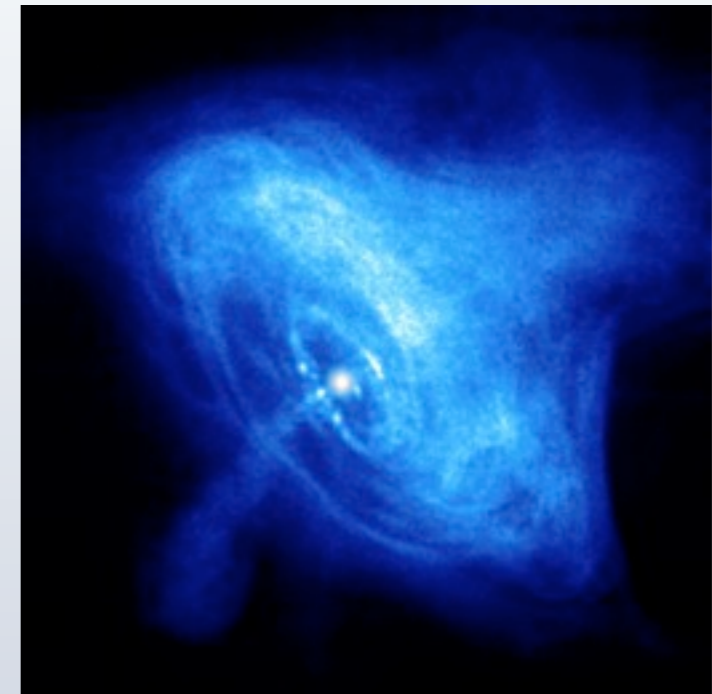
- Long term activity: but nu-fluxes are short lived, ~ seconds
- neutrinos drive baryon contamination
- Colliding shells? - Really fine tuned

Flare @ 100- 10⁵ sec in Short!

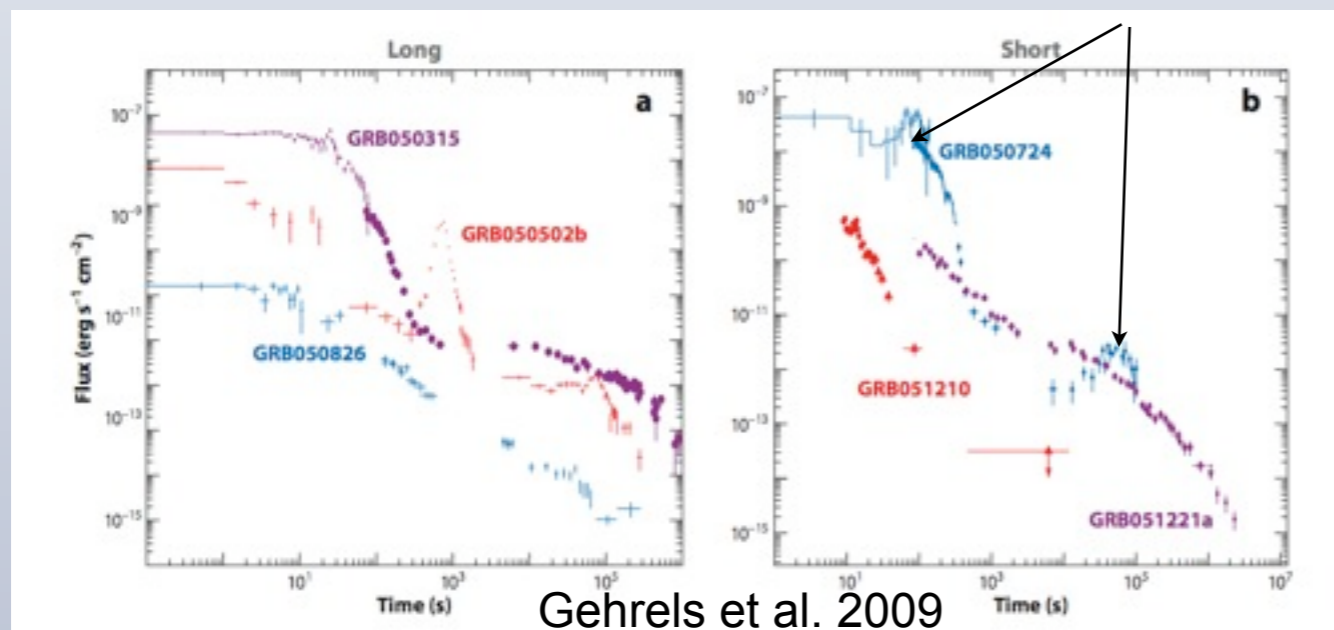


GRBs: magnetically driven, as PWNe

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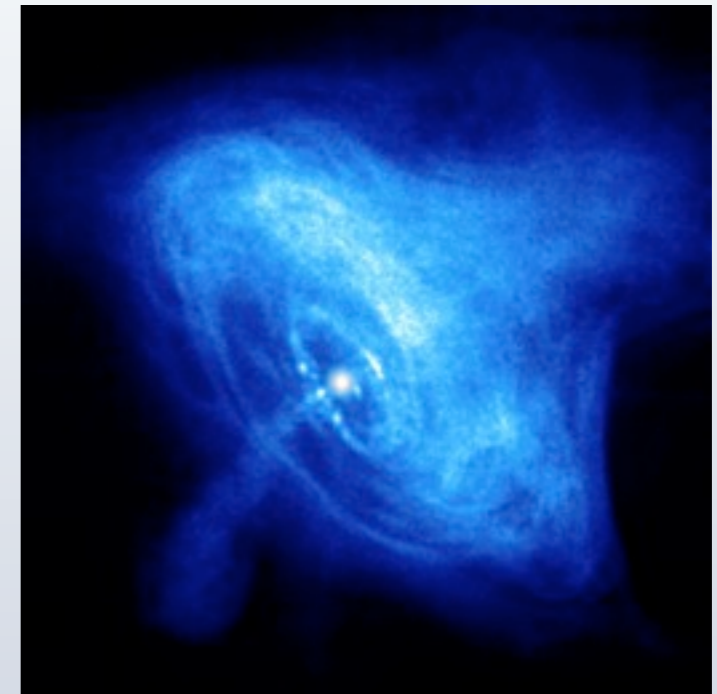


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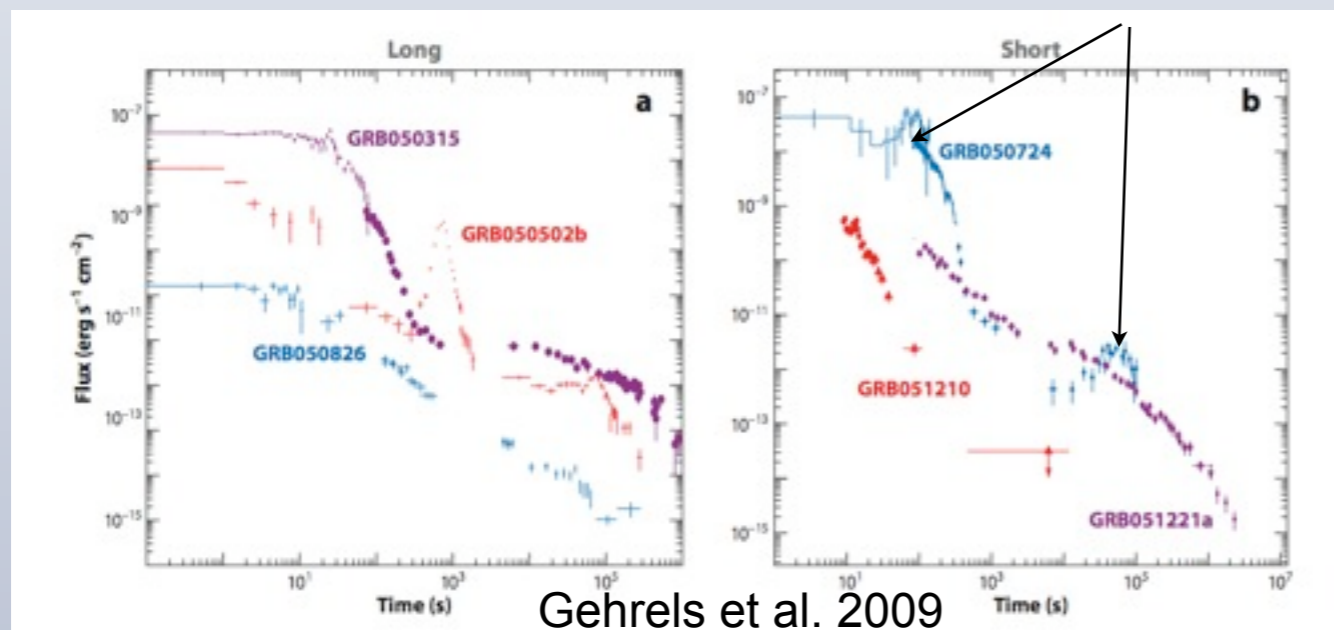


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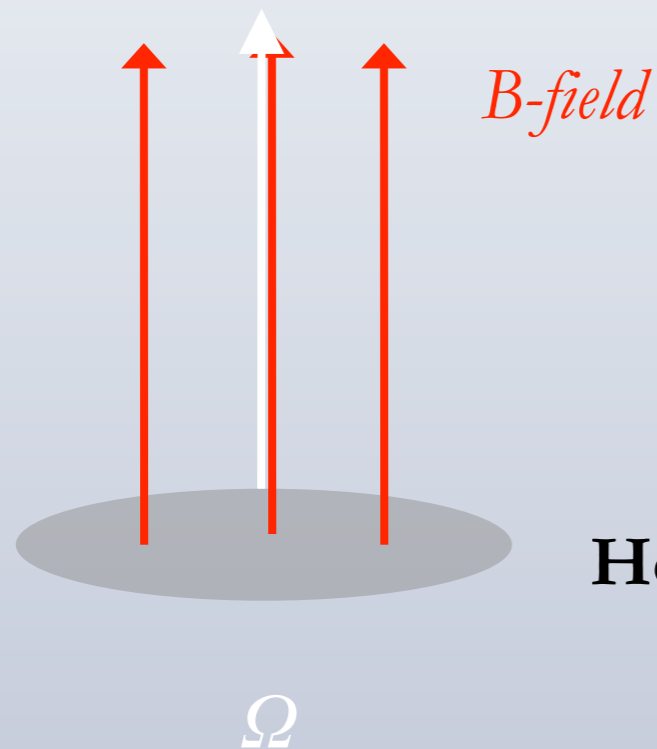


Which flare as well: constant energy supply produces bursts on sub-dynamical scales

Unipolar dynamo (Faraday wheel)

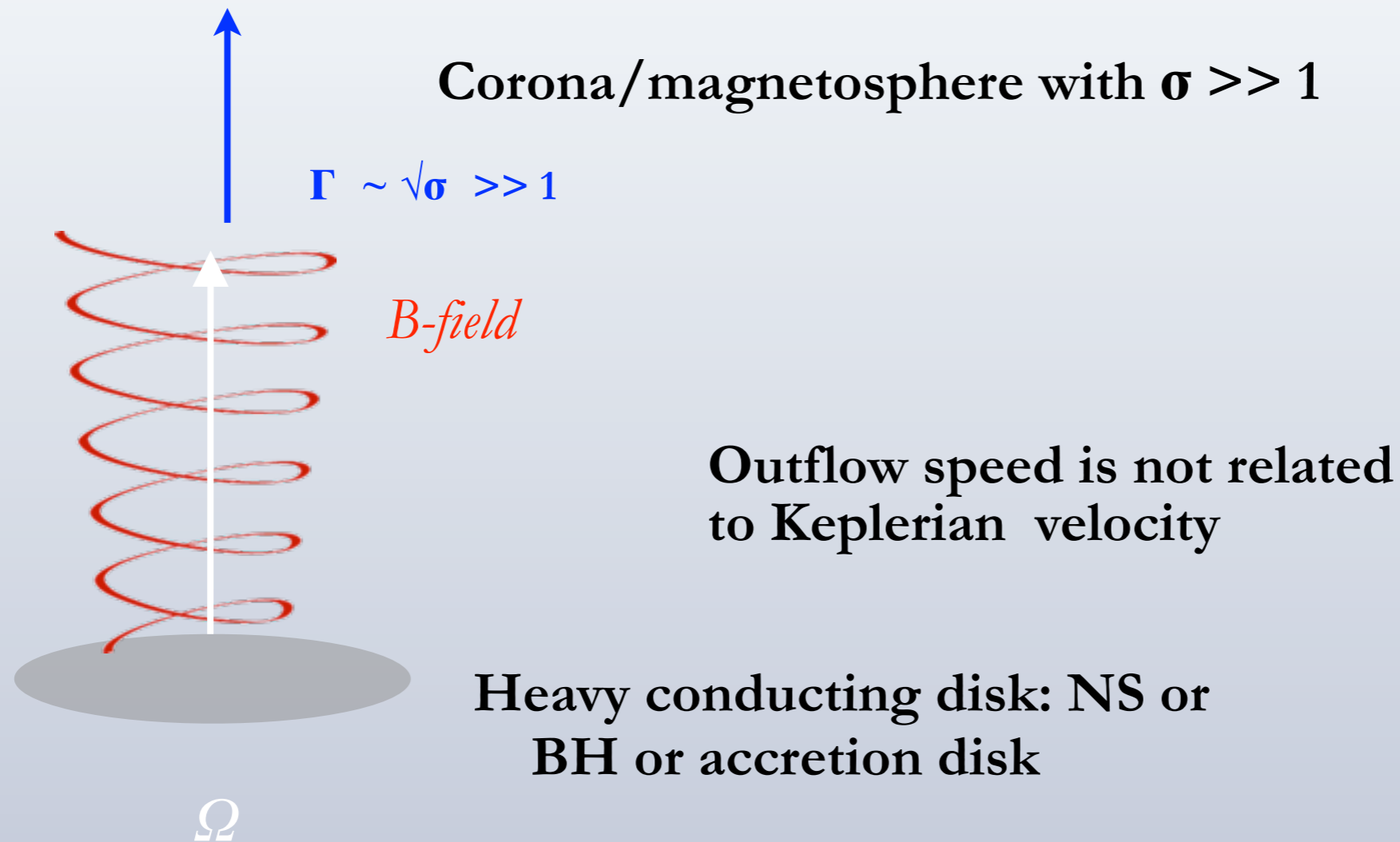
Corona/magnetosphere with $\sigma \gg 1$

$$\sigma = \frac{B^2}{4\pi\rho c^2}$$



Heavy conducting disk: NS or
BH or accretion disk

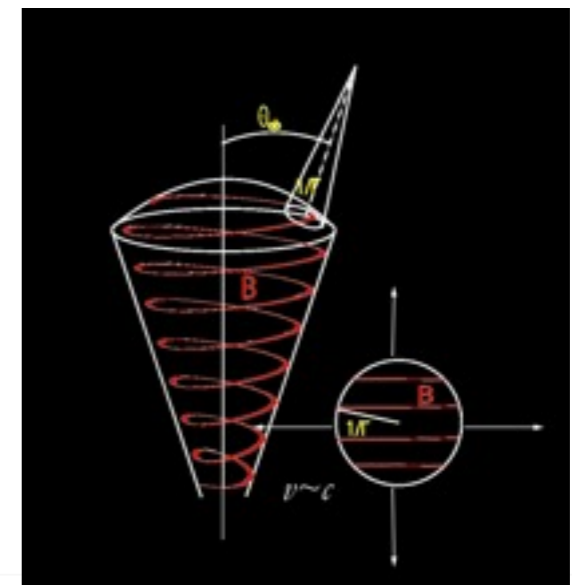
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Interesting concept... does it work?

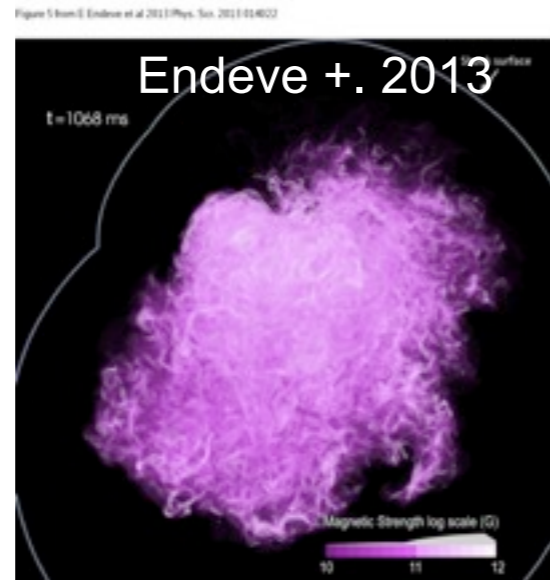
- Can $\sigma \gg 1$ magnetosphere be realized?
 - hot plasma (core collapse or NS-NS merger) - nu-driven contamination by baryons, but only for few seconds
 - dissipation inside a star (later in the talk)
 - GRB outflows must be clean - yes, it can.
- How B-fields accelerate and collimate the flow
- Do B-fields continue into the outflow?
 - fireball model: no, but are recreated at matter-dominated shocks
 - EM model (Lyutikov & Blandford 2003): yes, **dissipation & acceleration is magnetic** (not shocks)
- Are there evidence of large-scale B-fields? - Polarization
 - prompt (Coburn & Boggs 2003, others)
 - optical afterglows (e.g., Mundell+ 2013, others)
- How B-field dissipates and accelerates particles



B-field generation (need $\sim 10^{15}$ G)

Core-collapse (Long GRBs)

- Compression, shear, turbulent, MRI and/or alpha-Omega dynamo operate during collapse and core bounce



- Most MHD core-collapse simulations do not treat B-field generation, but start with huge magnetic fluxes.
- Even non-magnetic explosions are not settled...

NS-NS mergers (Short GRBs)

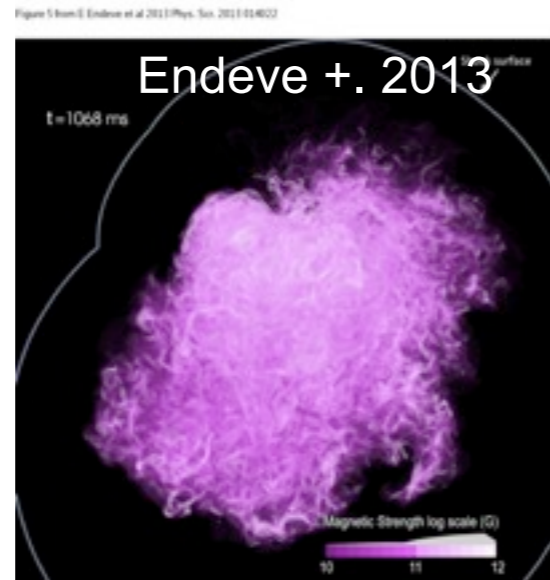


- dynamo in the supermassive NS (Price & Rosswog)?
- shear in the torus (Rezzolla+)?
- (Both saw amplification to $> 10^{15}$ G)

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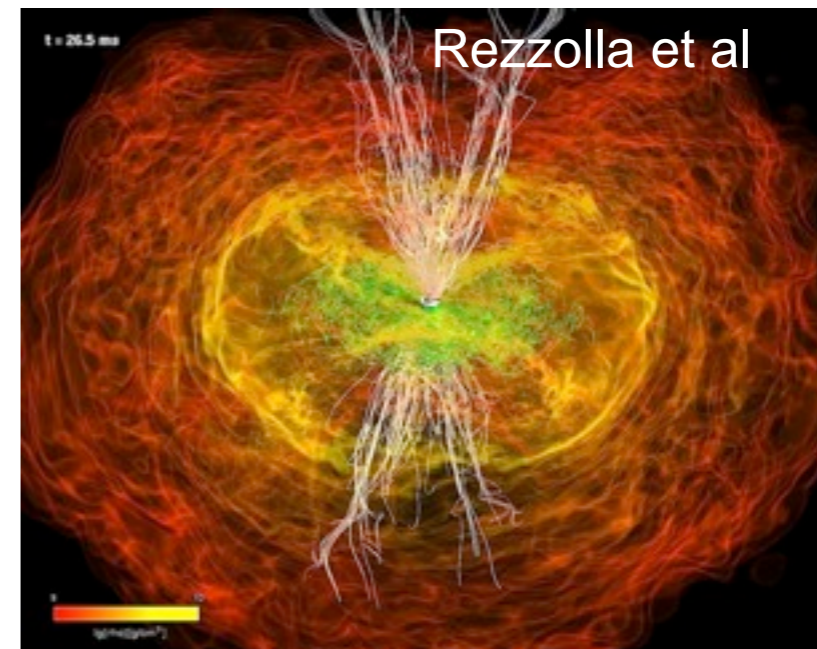
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Extremely challenging simulations

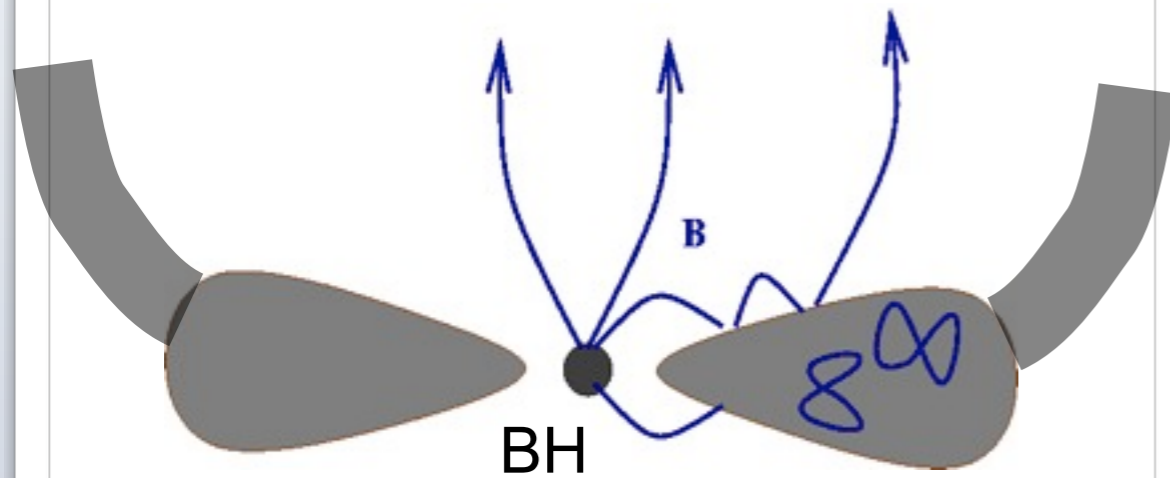
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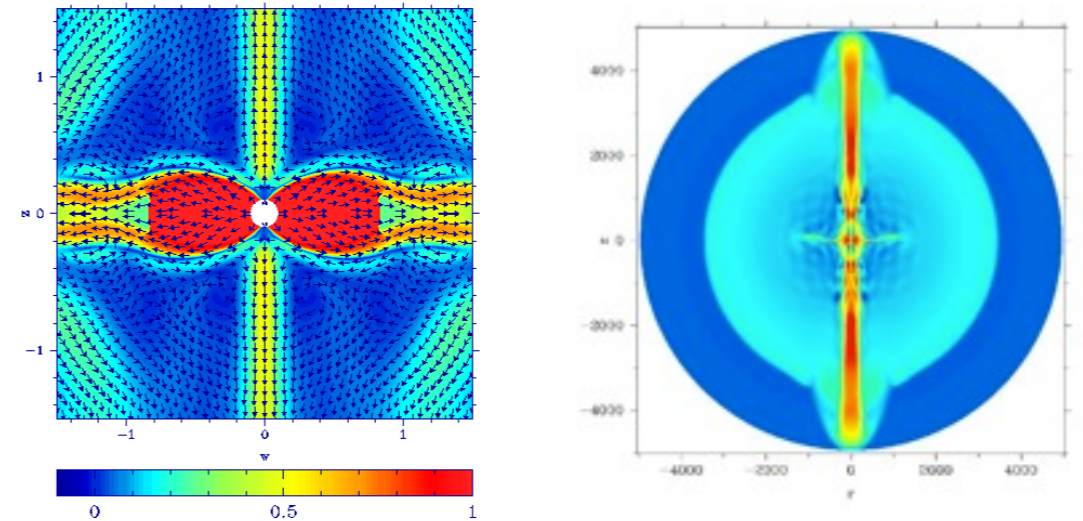
Two types of Faraday wheel in GRBs

- Collapsar: AGN-like (BZ)



- **B-fields are externally supplied and confined to a BH** (Blandford-Znajek)
- Confining walls (Lynden-Bell, Uzdensky)
- High-Gamma jet, superfast
- Perhaps a weak oblique collimation shock, but mostly continuous nozzle-like acceleration

- Millisecond magnetar: **PWN-like**



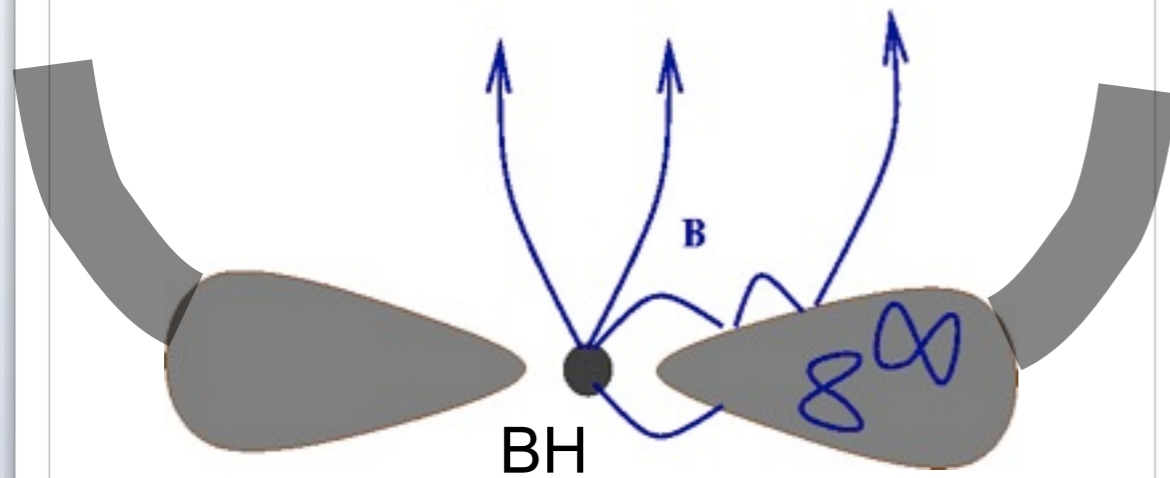
- **B-fields are intrinsic**
- Equatorially collimated initially

$$L \propto \sin^2 \theta$$
- Strong shock - stop!
- “Slowly” collimated by hoop stresses, sub-fast/super-Alfvenic plume (later nozzles out to super-fast)

Komissarov+,
Bucciantini+

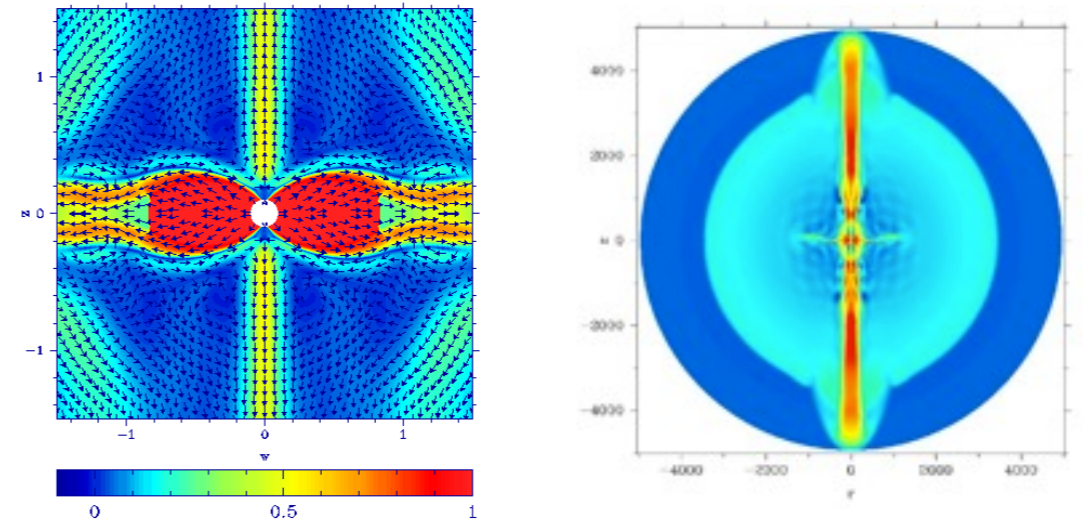
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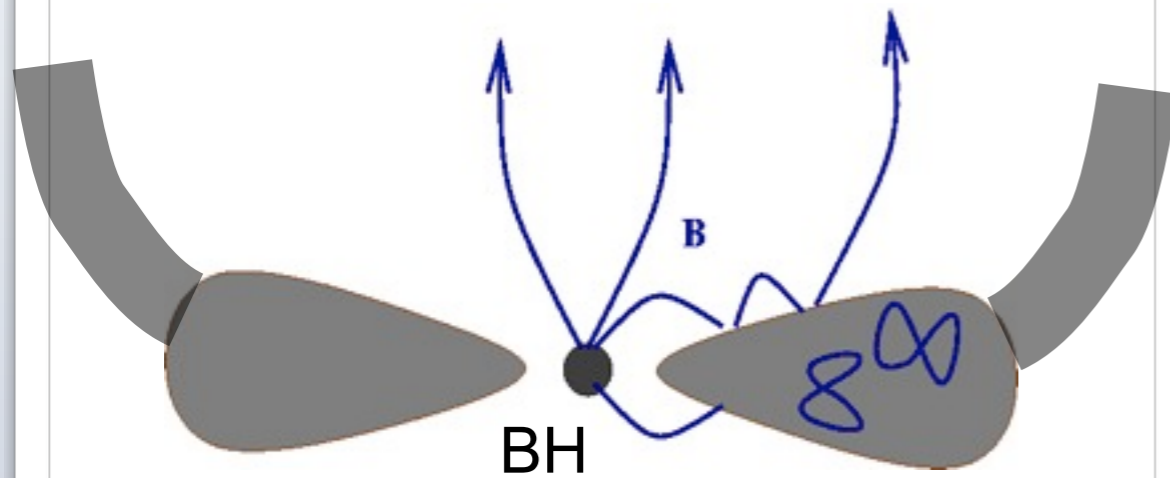


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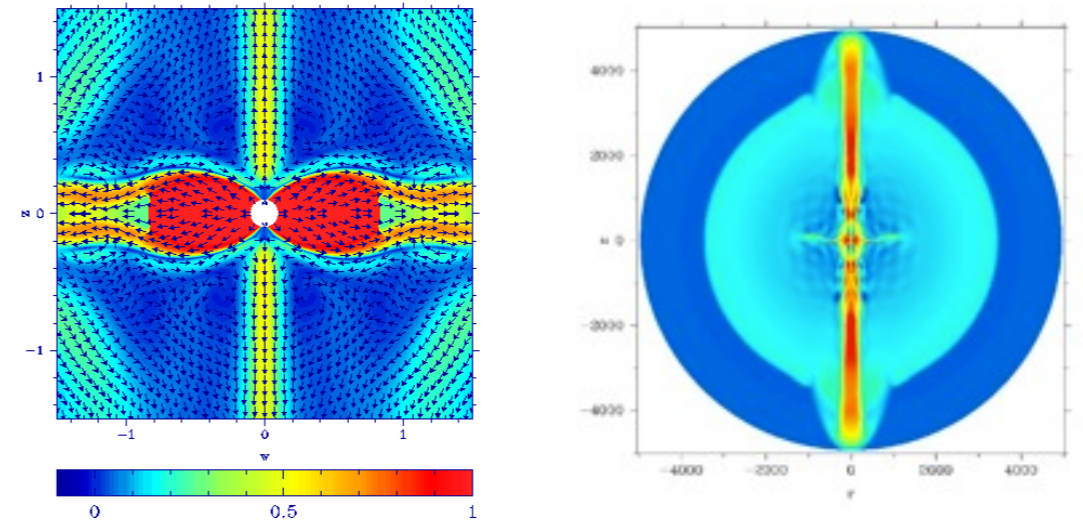
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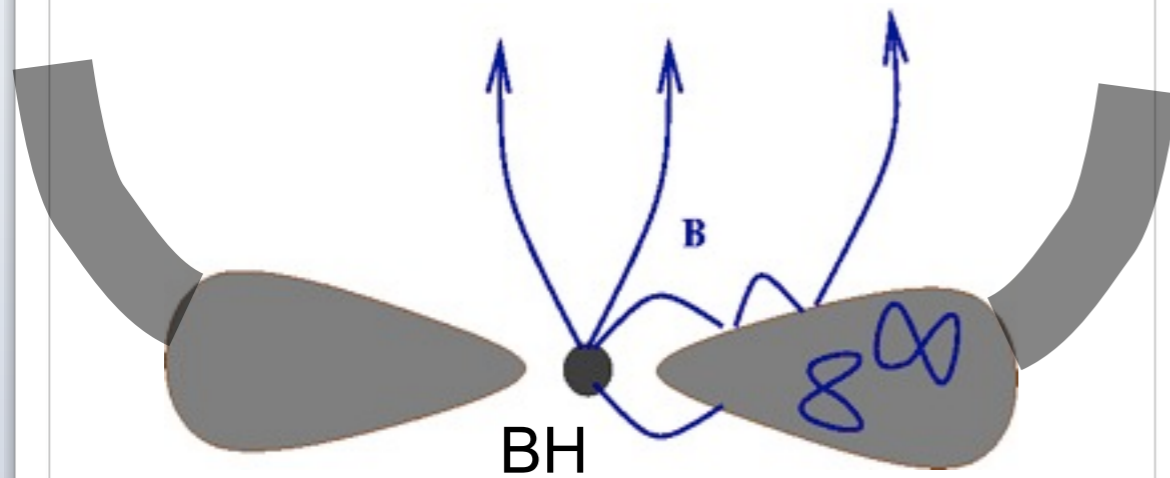


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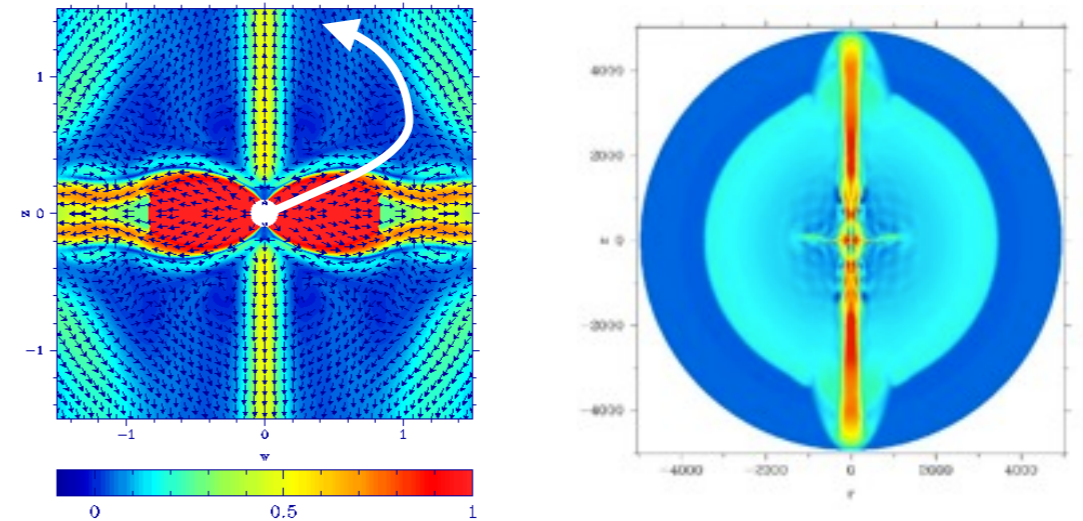
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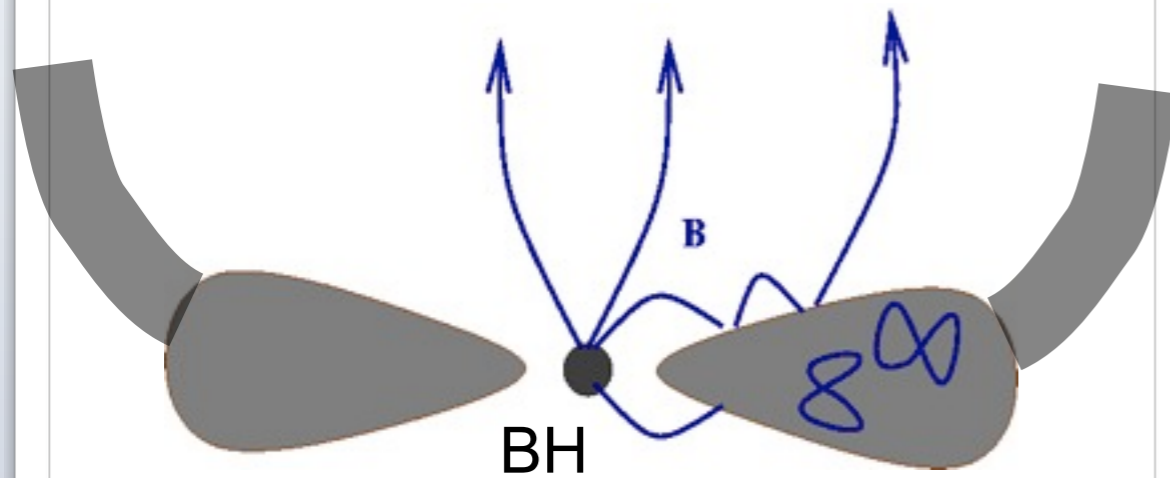


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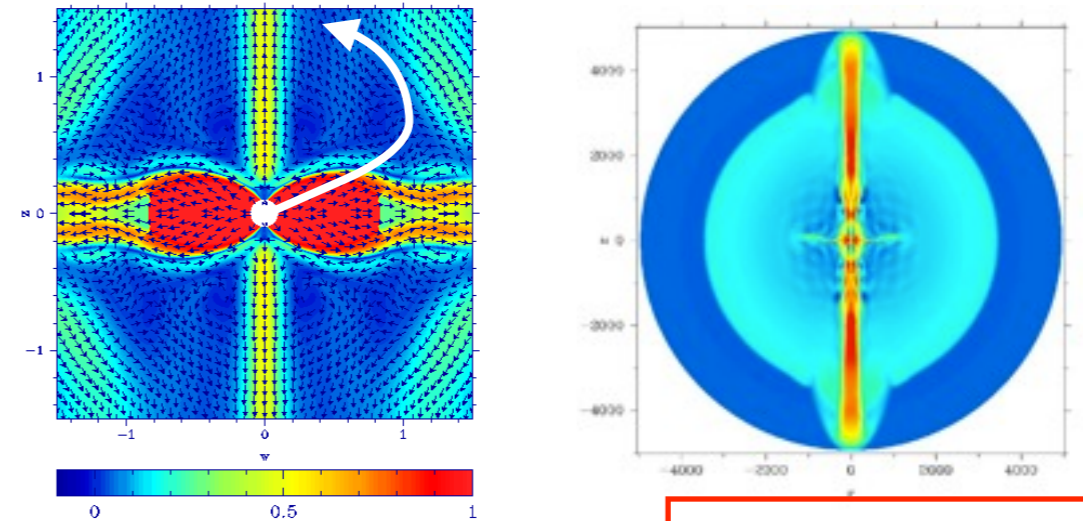
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Looks good! ?

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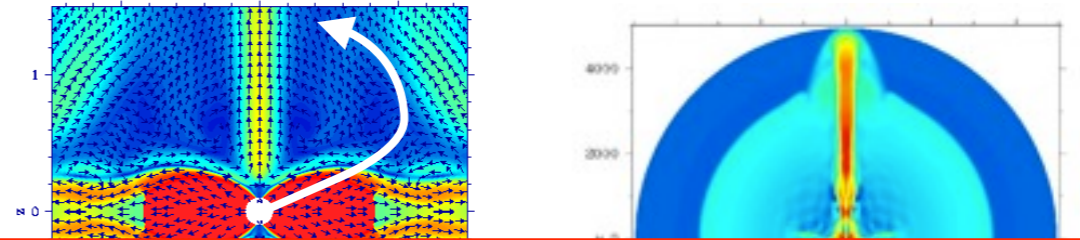
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- In application to Long GRBs central engine, both types of model have two related problems in making a jet
 - the sigma problem
 - stability

and confined to a BH
(Blandford-Znajek)

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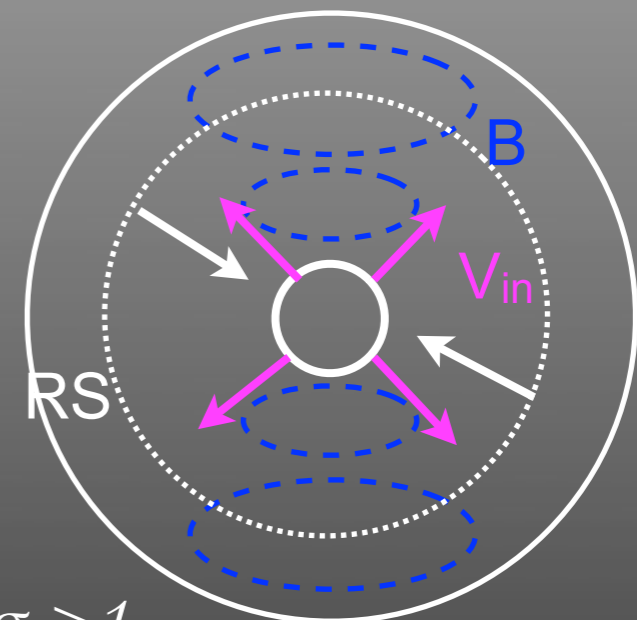
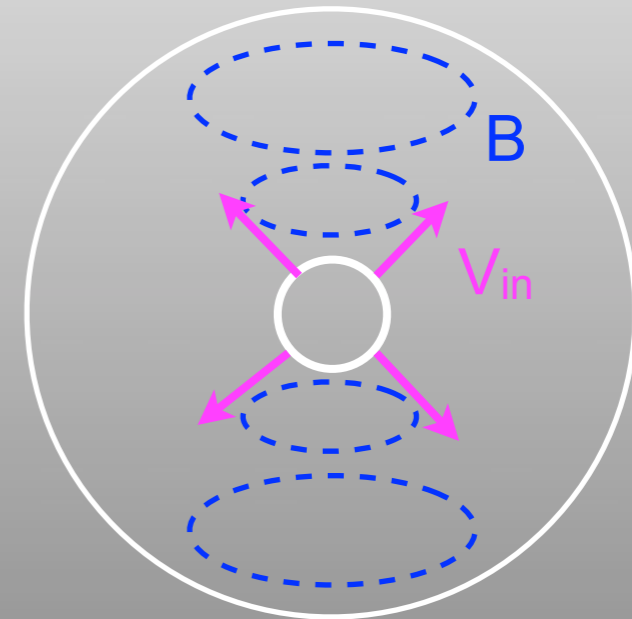
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The sigma problem

Rees & Gunn
Kennel & Coroniti

Consider a fixed cavity into which a central source injects **energy and magnetic flux** linearly with time. E.g. magnetar cavity is nearly constant on light travel time.

Stored B-flux $\sim t$, toroidal B-field $\sim t$, stored energy $\sim B^2 \sim t^2$???



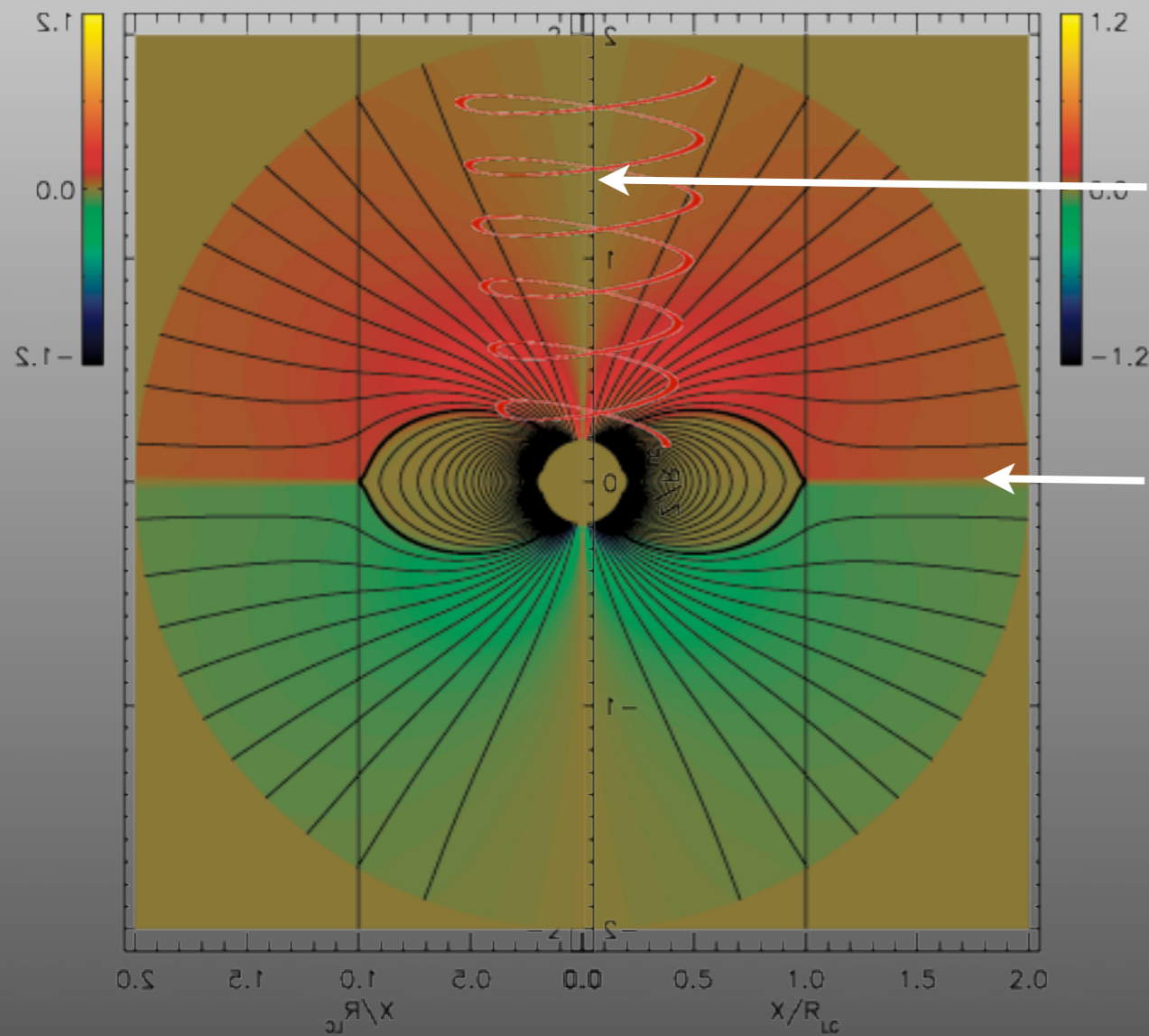
Kennel&Coroniti: in $\sigma > 1$, reverse shock would reach the central engine in light crossing time and model breaks down

$$V_{RS} \sim c \text{ for } \sigma > 1$$

Possible resolutions

- Kennel & Coroniti: σ must change to $\ll 1$ on the way.
(NB: this is a requirement of the self-consistency of the model, not a measured parameter within the model).
- σ remains high, but shock is not MHD (kinetic effects dominate, Lyubarsky). Unlikely in magnetars, too dense.
- **Most of magnetic flux should be destroyed between the source and the boundary. The flow must become dissipative.**

Need to destroy magnetic flux: reconnection.



O-point reconnection

X-point reconnection

Plasma will flow towards reconnection sites

Ideal flow in the bulk, dissipation on the axis & equator

$$2\nabla I^2 = r^2 \sin^2 \theta \Delta \Phi \nabla \Phi$$

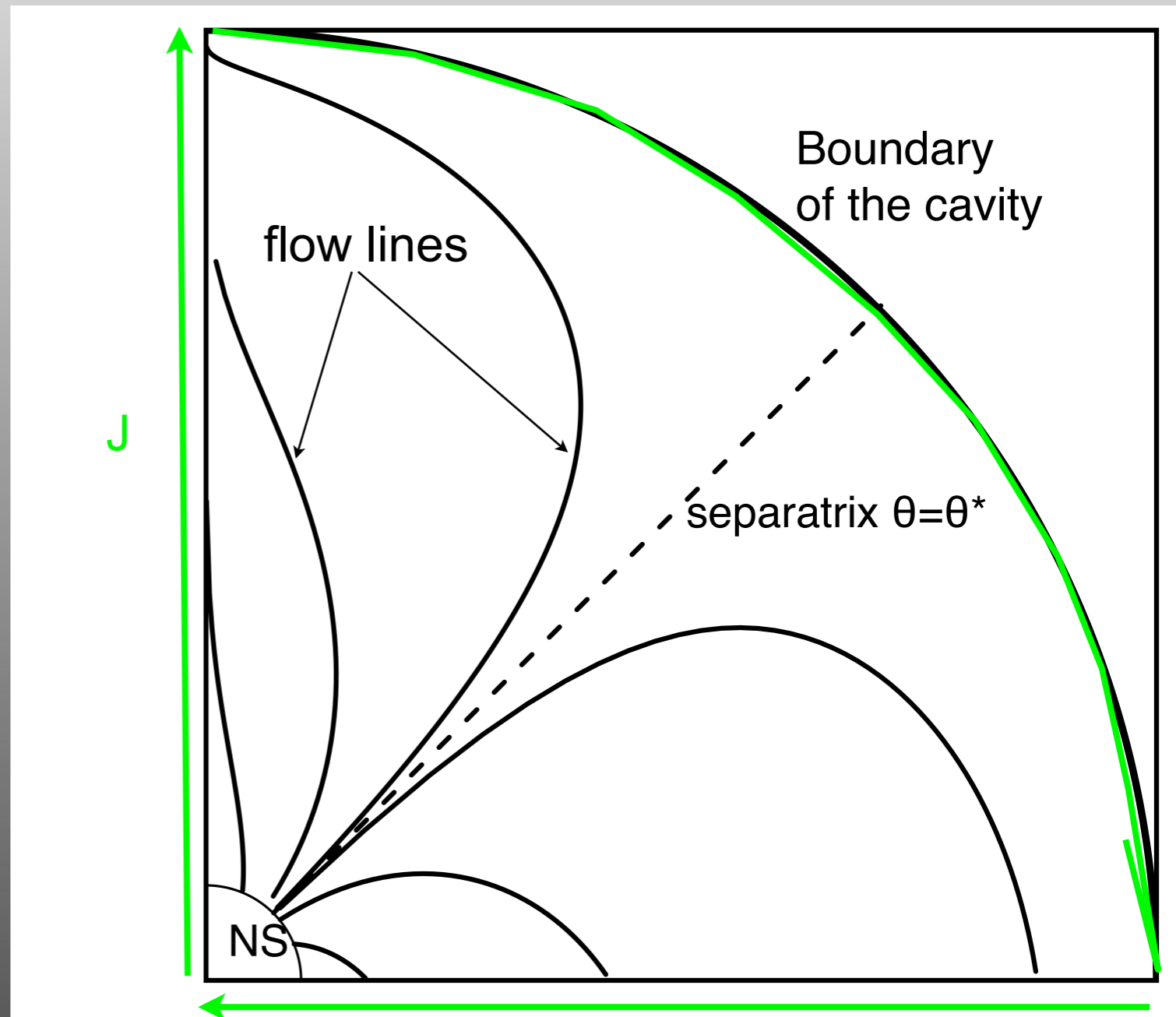
- Current and charge distributions are related

$$\mathbf{v} = \frac{\mathbf{E} \times \mathbf{B}}{B^2} = \frac{r \sin \theta (\mathbf{e}_\phi \times \nabla \Phi)}{2I}$$

- Vanishing charge and current densities in the bulk.

- $I=I_0$ - current on the axis

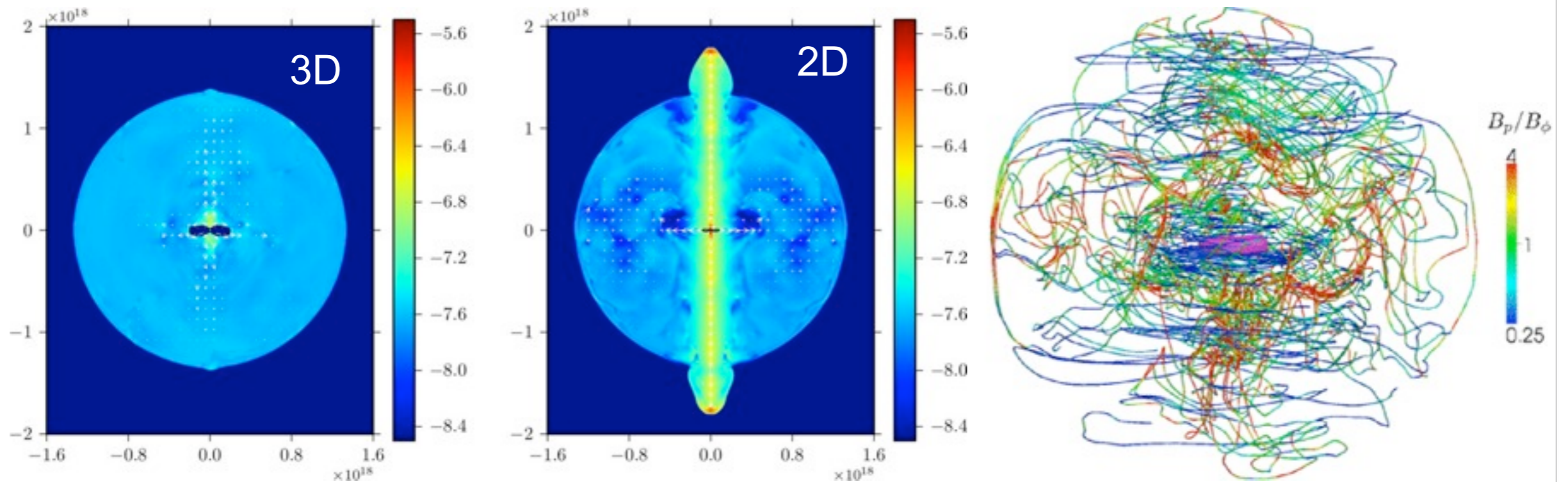
$$\Phi = \Phi_0 \left(1 - \frac{R}{r} \right) \ln \frac{\tan \frac{\theta}{2}}{\tan \frac{\theta^*}{2}}$$



On the axis: toothpaste tube effect

The stability problem

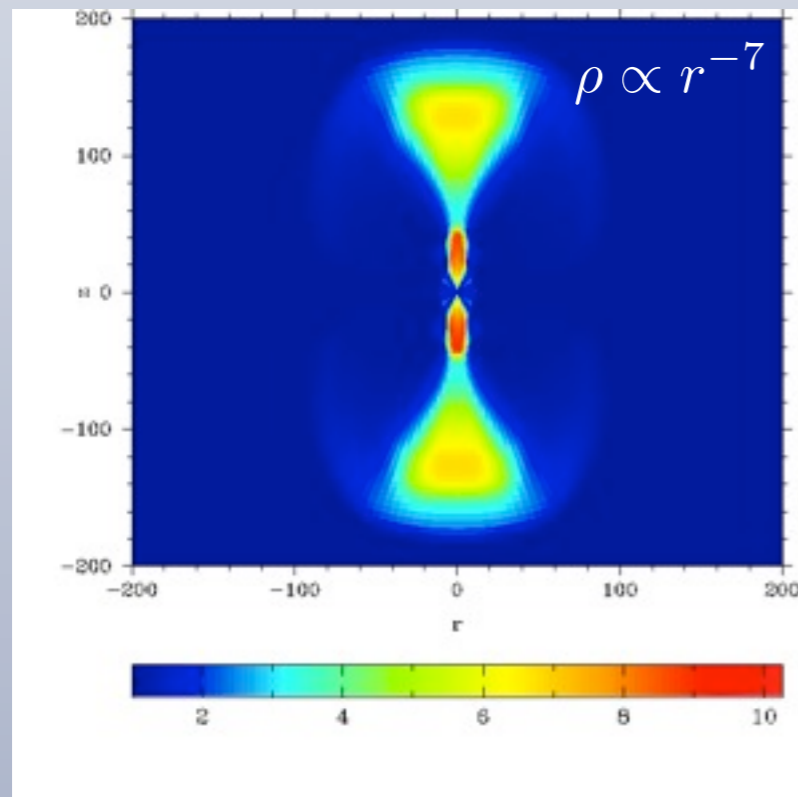
- **How to kill flux: twisted B-field susceptible to kink instabilities**
- First 3D relativistic MHD simulations of PWN, $\sigma \gg 1$



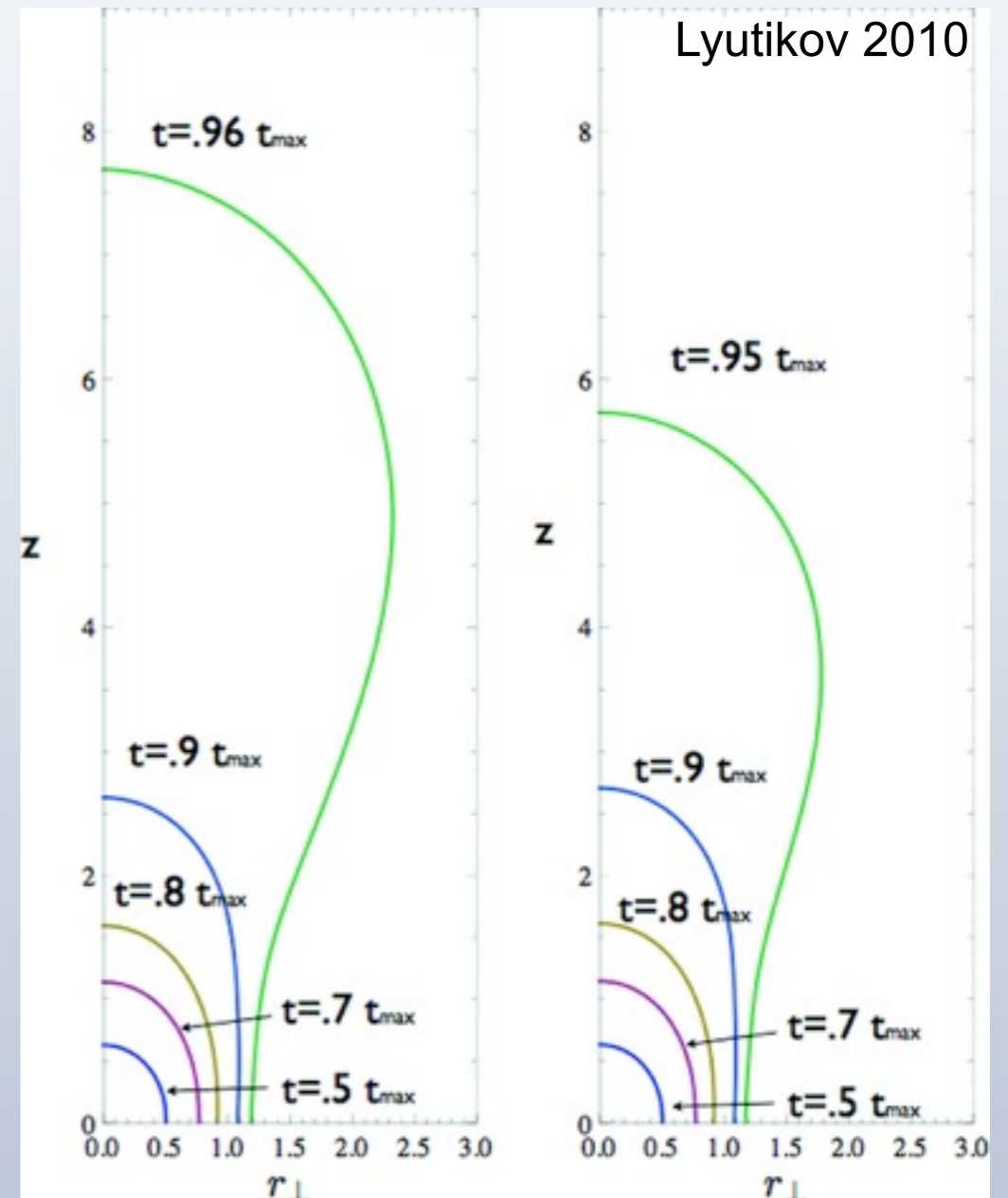
- **Magnetic flux is destroyed. Sigma problem solved.**
- axially-symmetric simulations way overestimated the stability and Lorentz factor of the jets: 3D jets are slow and susceptible to instabilities (in BH-driven jets as well).
- But there is no jet left, only a plume.... Sigma problem became the no-jet problem (~OK for PWNe, not OK for GRBs)

And what to do?... Wait a second

- Wait a few seconds for neutrinos to do the hard part, explode the star.
- B-field is amplified on a contracting proto-NS. Launch a **slightly non-spherical** shock inside a star.
- FS propagates in **expanding envelope** ($v \sim r$), with **sharp density gradient** - shock accelerates, makes a key hole (chimney)



Lyutikov & Komissarov, in prep



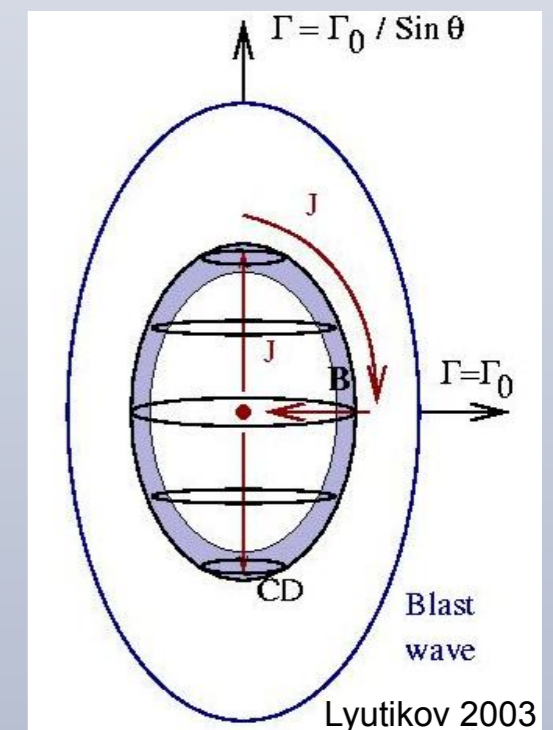
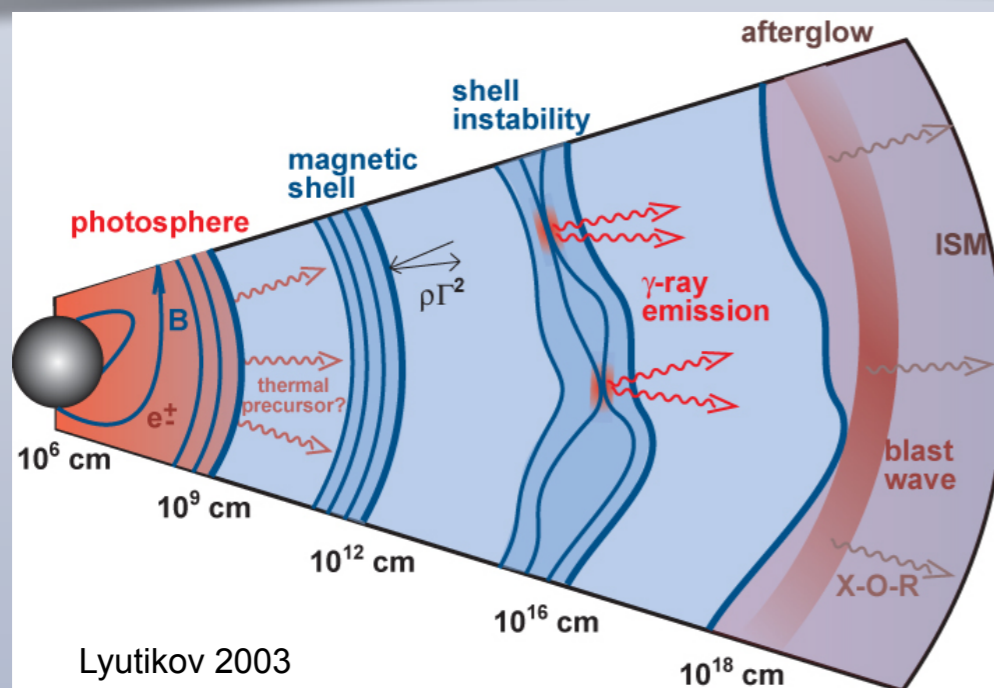
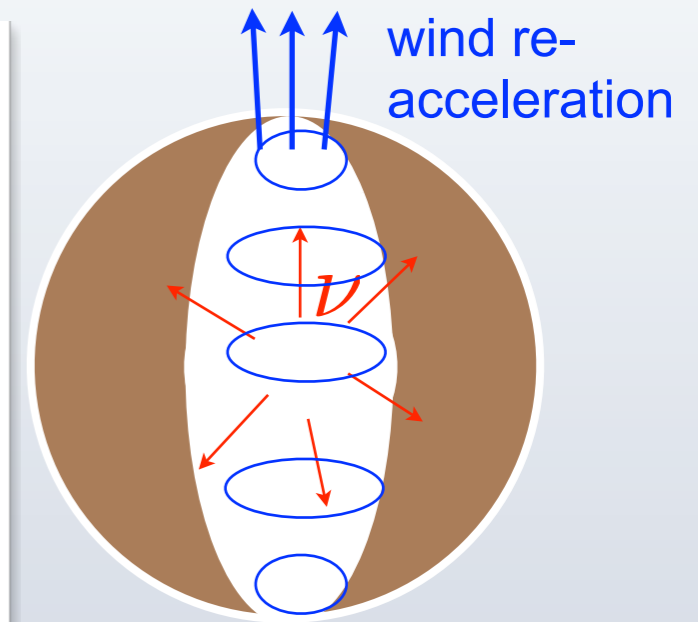
Anisotropic driving

Anisotropic r_0

$$1 + \cos^2 \theta$$

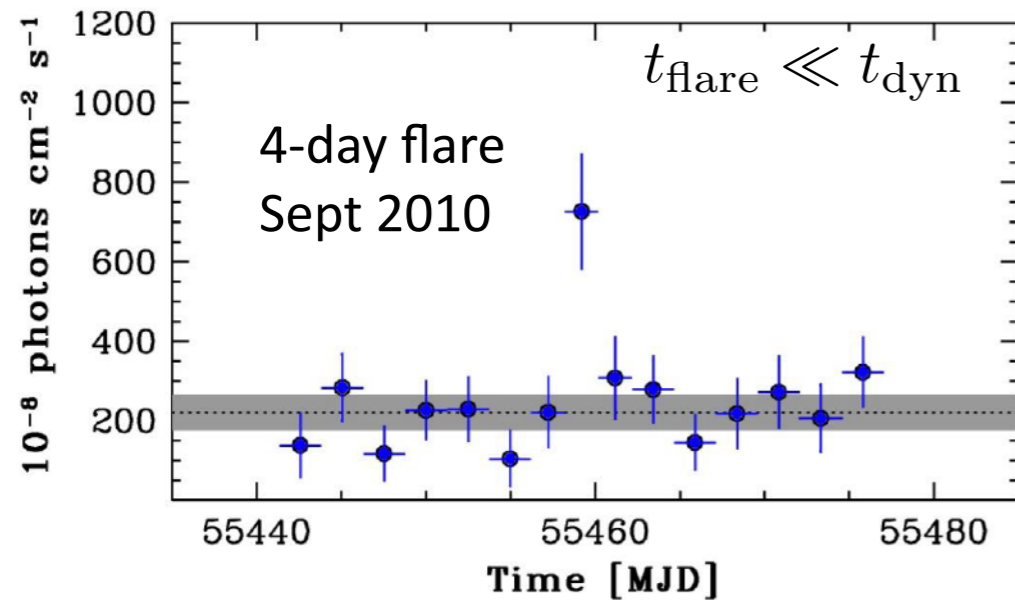
After the break-out the wind reaccelerates

- Force balance: $\Gamma \sim \Gamma_0 / \sin \theta$ (but: instabilities?)
- Engine active for 100s secs: magnetars are good for **late activity** (no need for long accretion)
 - not by powering the FS (little energy), but by **internal dissipation** in the long-lasting wind
- But magnetar emission is smooth? - **Peaks in the profile are signatures of the bursty dissipation in the wind, not the central engine activity** (Lyutikov & Blandford 2003, Lyutikov 2006)

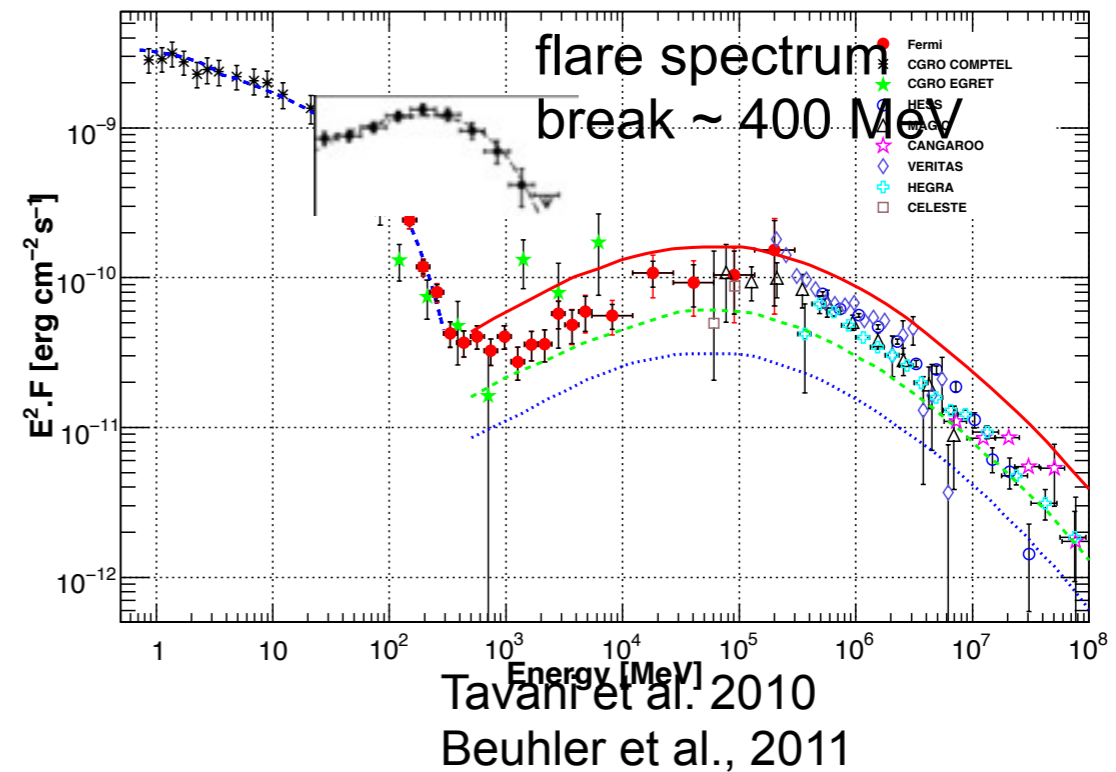


Bursty dissipation in the wind: Crab flares!

Recall: magnetar models of
GRBs ~ models of PWNe



Tavani et al. 2011



Accelerating E-field < B-field

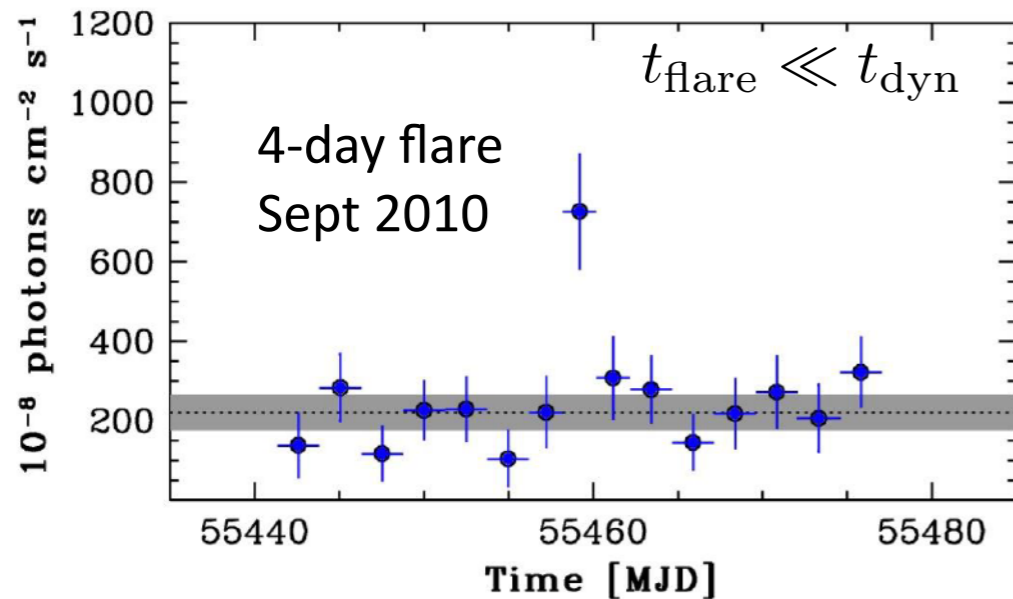
$$eEc = \eta eBc = \frac{4e^4}{9m^2c^3} B^2 \gamma^2$$

$$E_p = \frac{27}{16\pi} \eta \frac{mhc^3}{e^2} = 236 \eta \text{ MeV.}$$

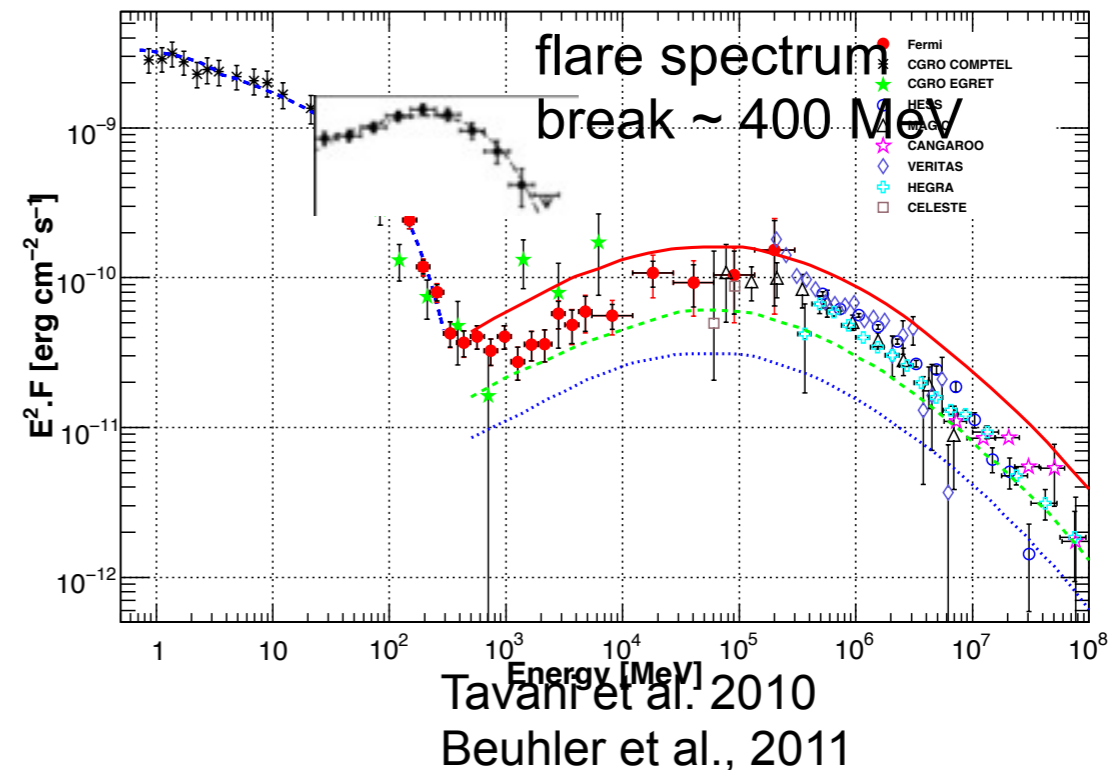
Lyutikov '10,
Komissarov & Lyutikov '11
de Jager '98 (for shocks)

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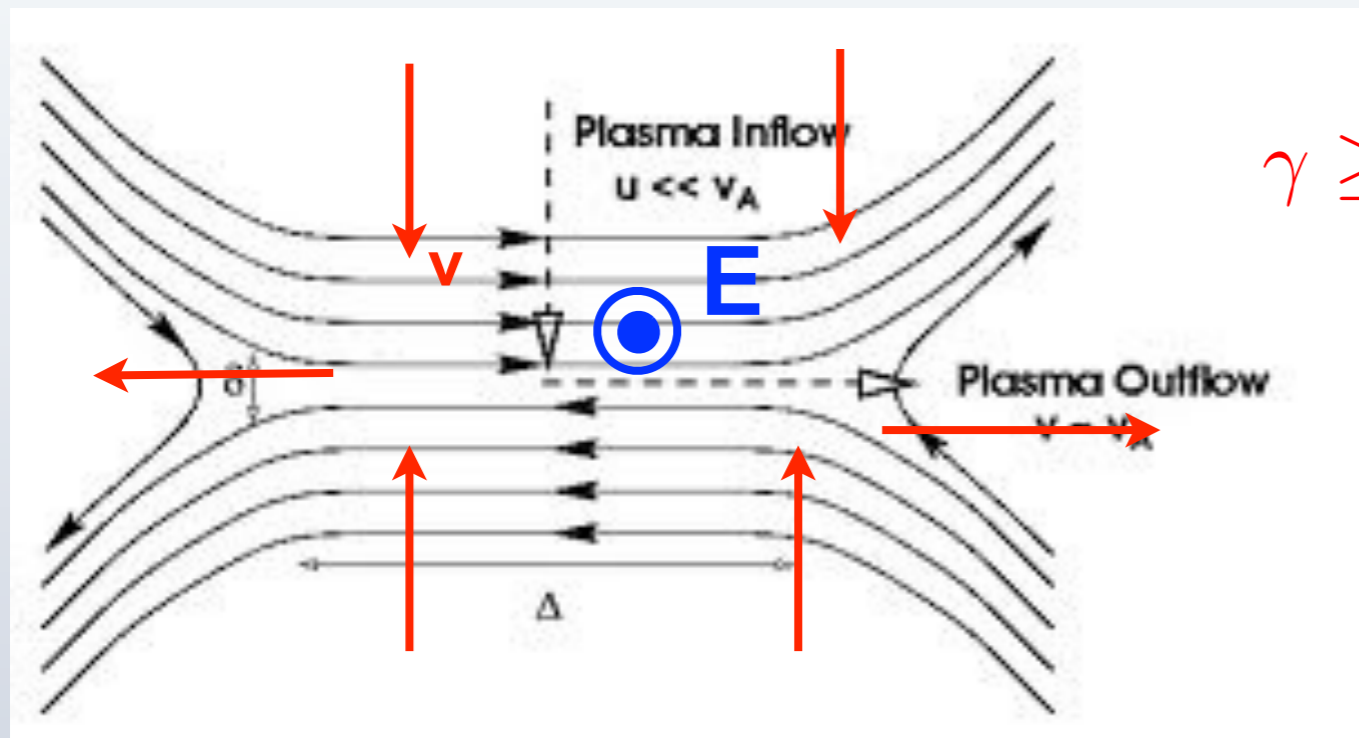
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Lyutikov '10,
Komissarov & Lyutikov '11
de Jager '98 (for shocks)

- Highly magnetized, $\sigma \gg 1$, shocks are weak, not likely to be efficient accelerators.
- All the energy in the B-field: accelerate particles directly via **reconnection**.
- **Paradigm change (?): some (most?) particles are accelerated by magnetic reconnection (and not shocks)**

Reconnection: efficient, non-stationary

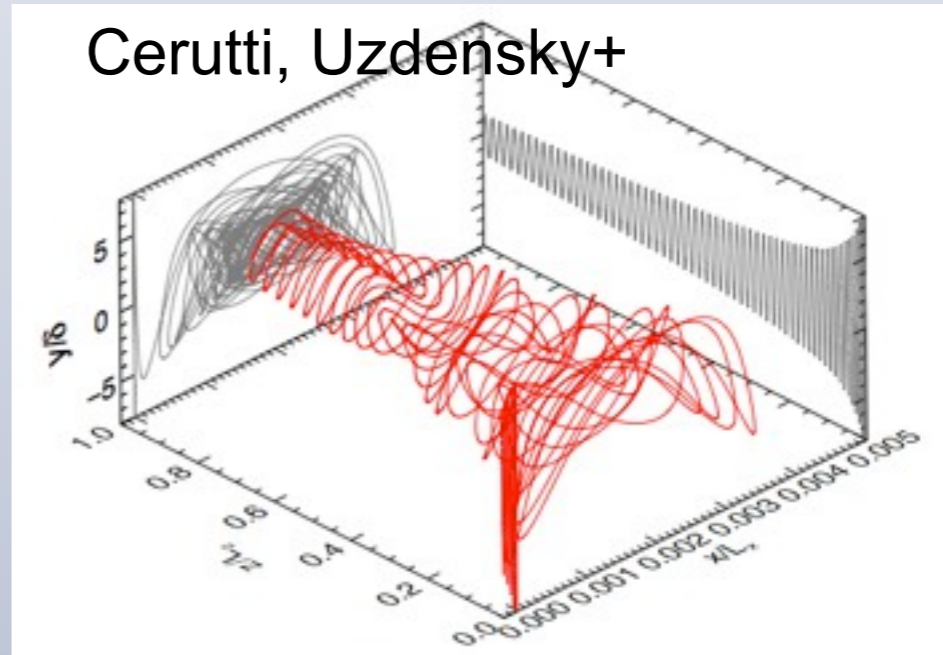


Reconnection in $\sigma \gg 1$ plasma: outflow can be relativistic (Lyutikov & Uzdensky 2002, Lyubarsky)

New plasma physics regime: $\sigma \gg 1$ plasma.

- **What are dynamic and dissipative properties of such plasmas? - very different from laboratory and space plasmas.**
- Pulsar winds, AGN & GRB jets and magnetospheres of BHs
- Alfvén velocity is highly relativistic
 - E-field is dynamically important
 - charge density is important

Cerutti, Uzdensky+



Reconnection can be bursty from smooth conditions

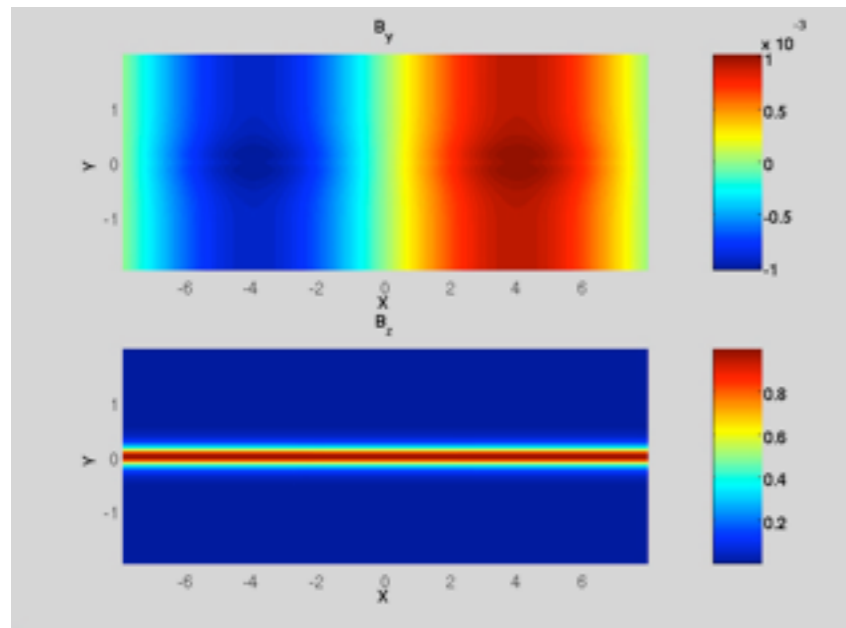
- Current sheet can be unstable to tearing

-

Lyutikov 2003, Komissarov+ 2007

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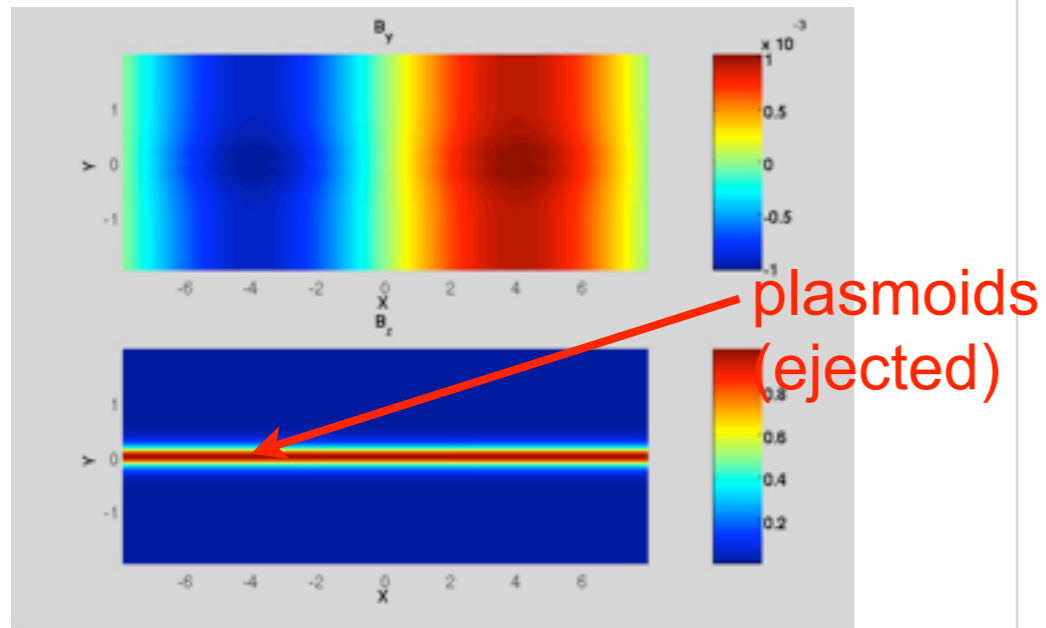
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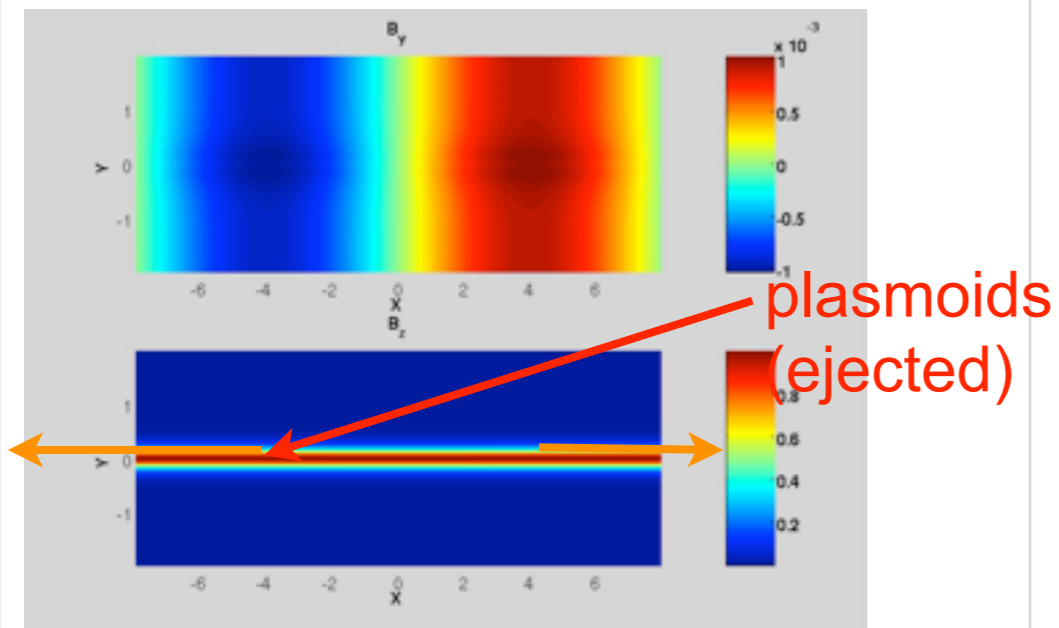
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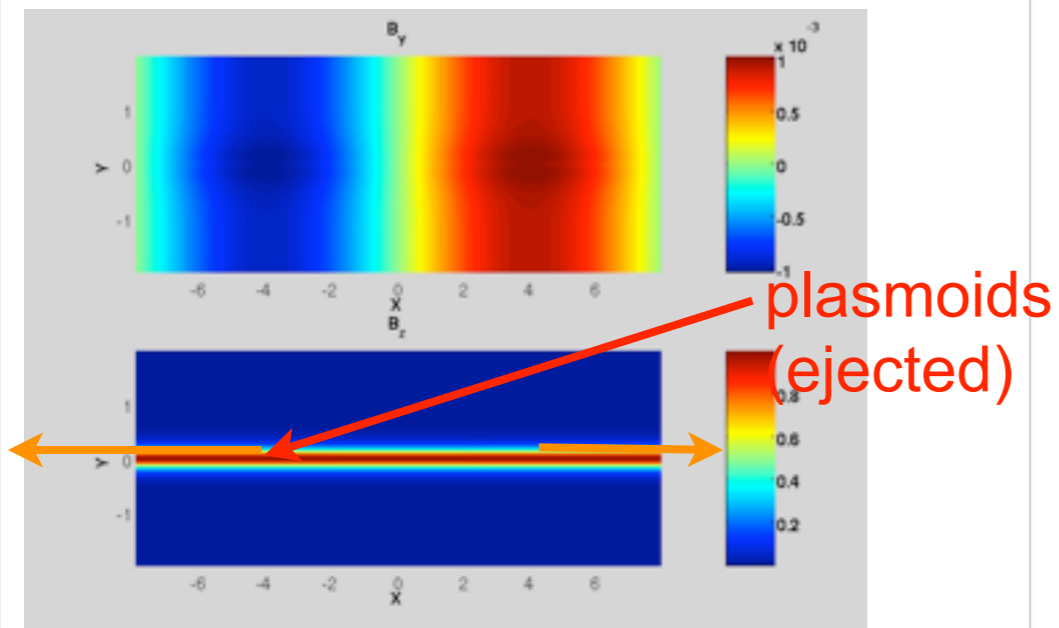
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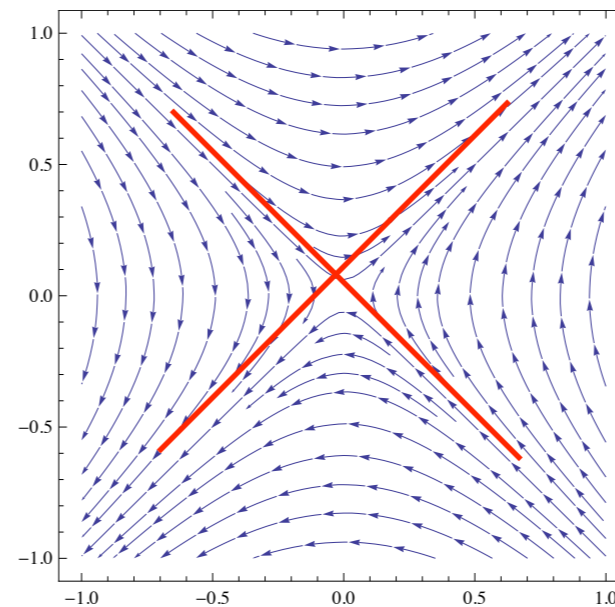
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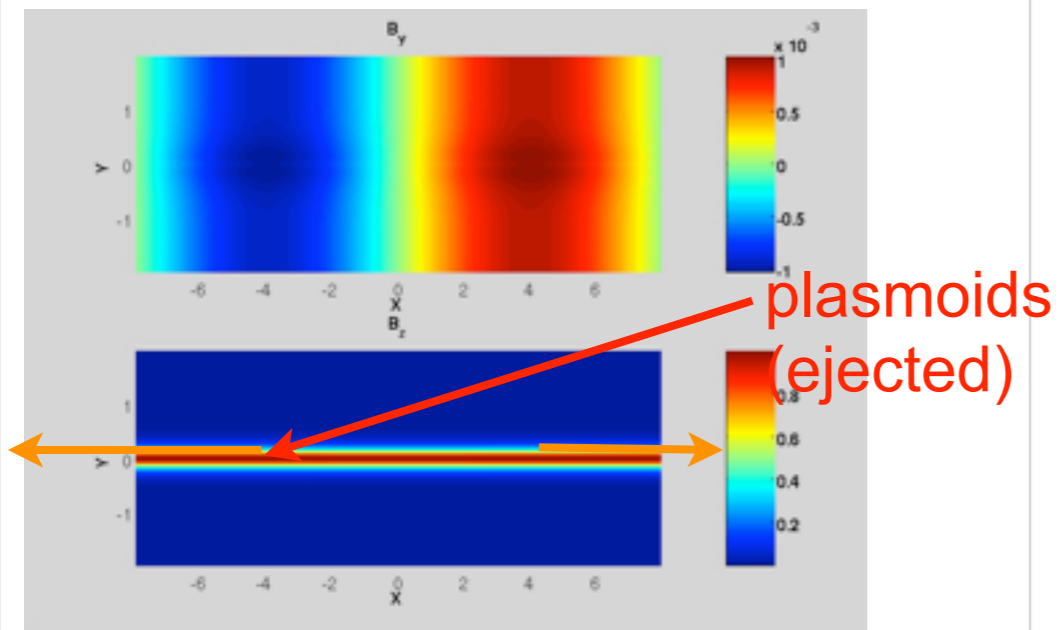
Lyutikov 2003, Komissarov+ 2007

X-point collapse (non-linear tearing?):



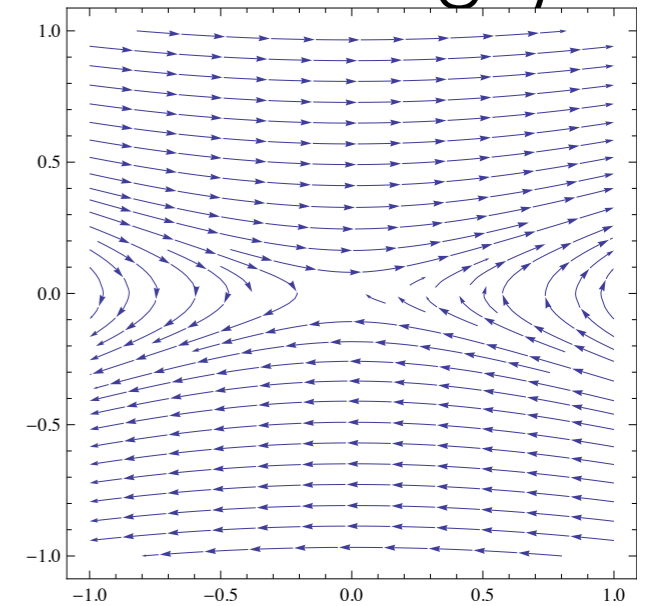
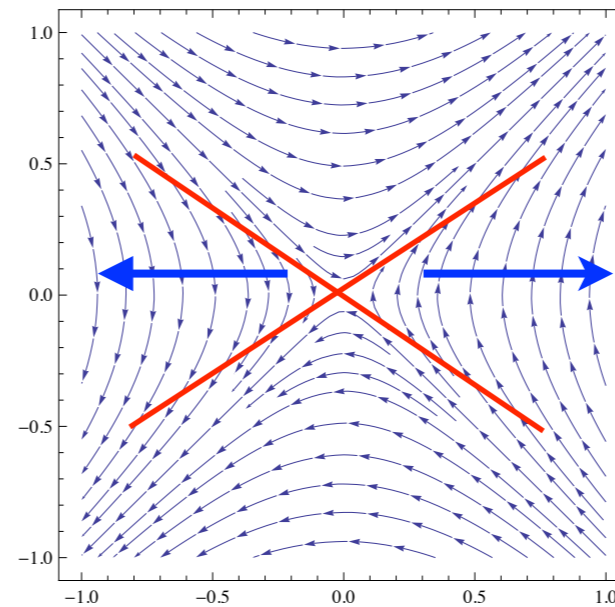
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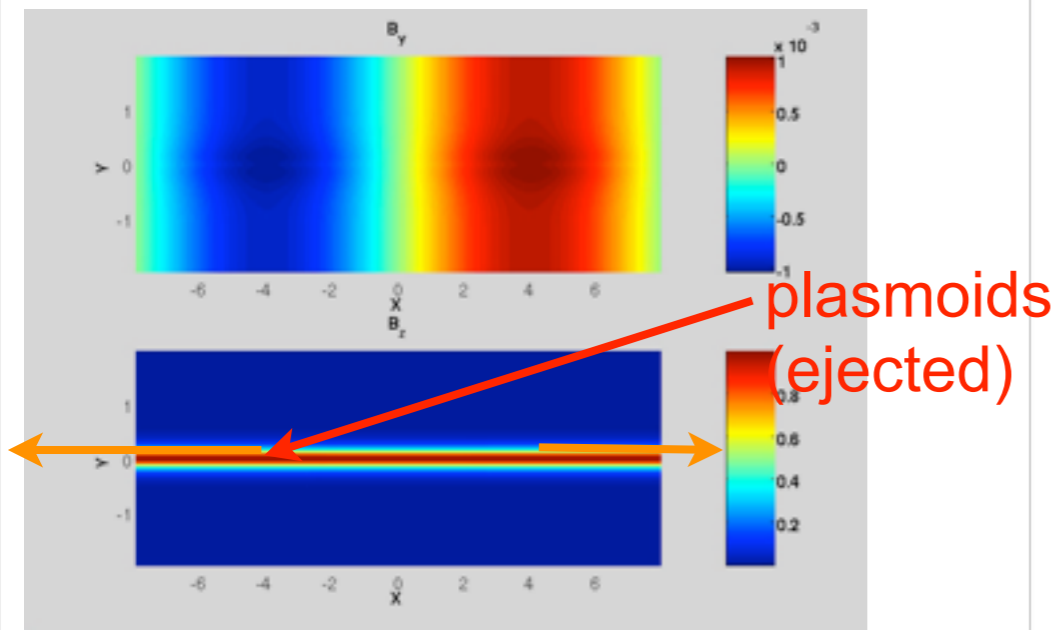
X-point collapse (non-linear tearing?):



- **explosive** dynamics on Alfvén (light) time
- Starting with smooth conditions
- $E \sim B_0$ (field outside), $E > B$ with resistivity

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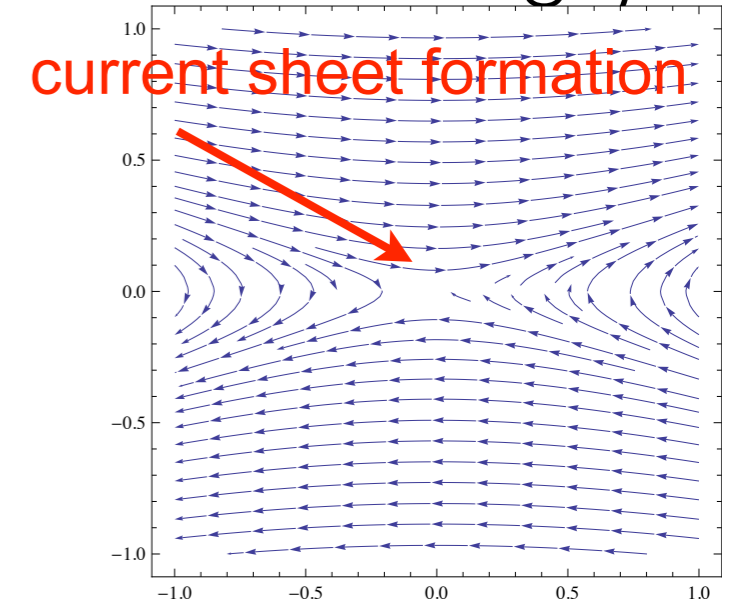
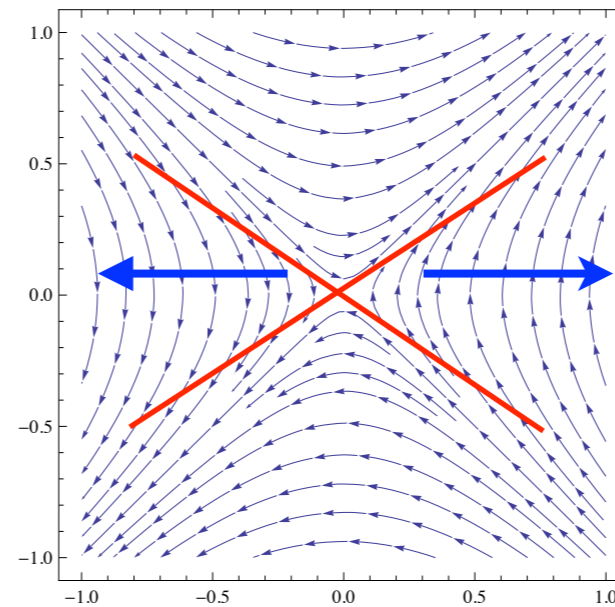
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plasmoids
(ejected)

Lyutikov 2003, Komissarov+ 2007

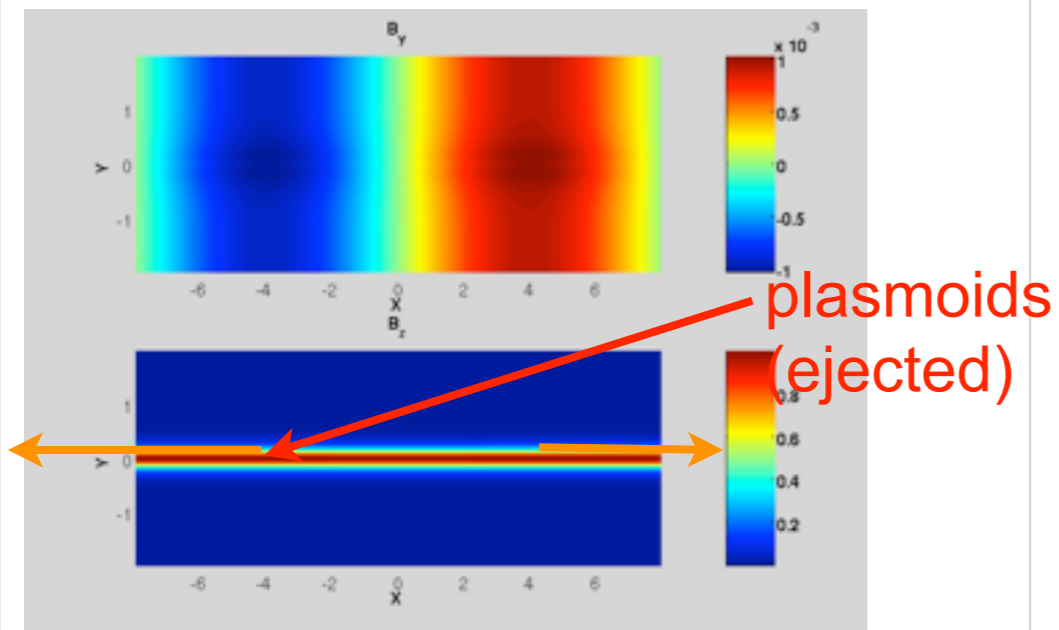
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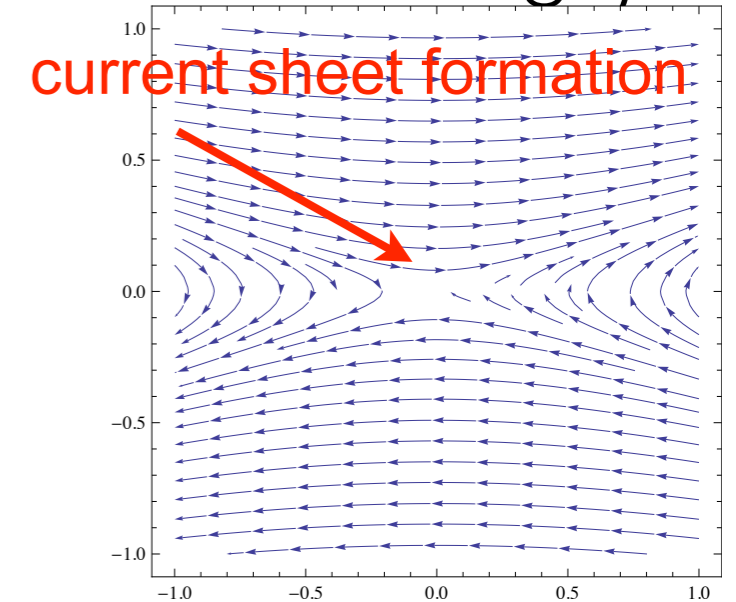
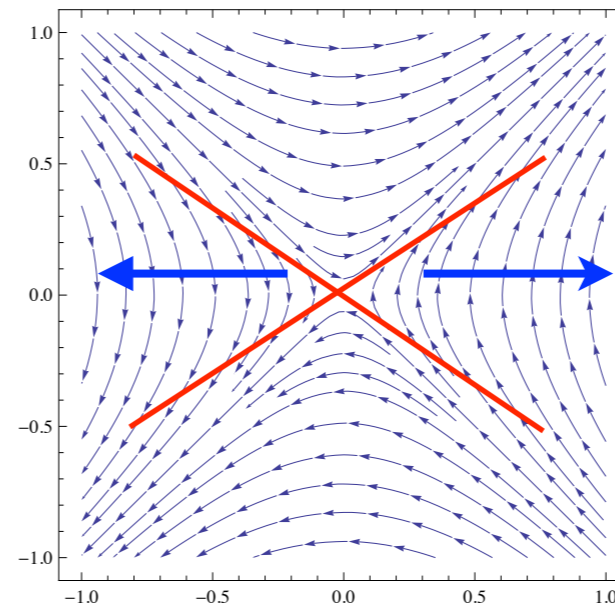
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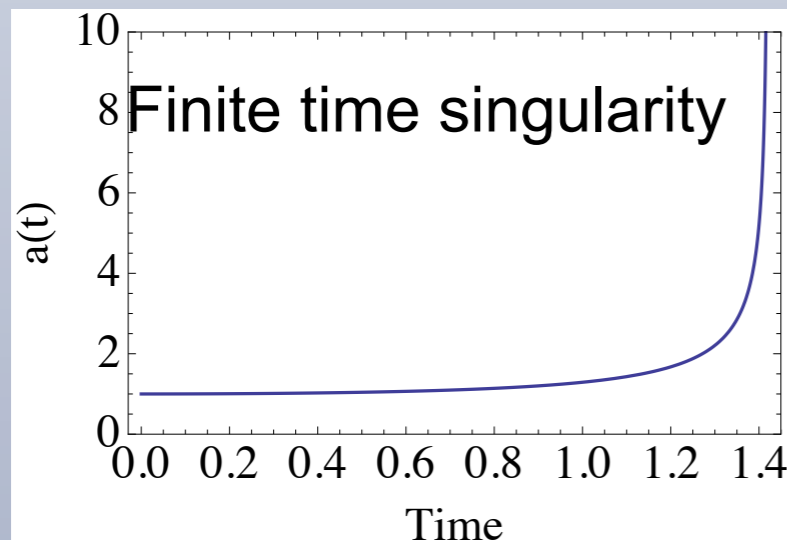


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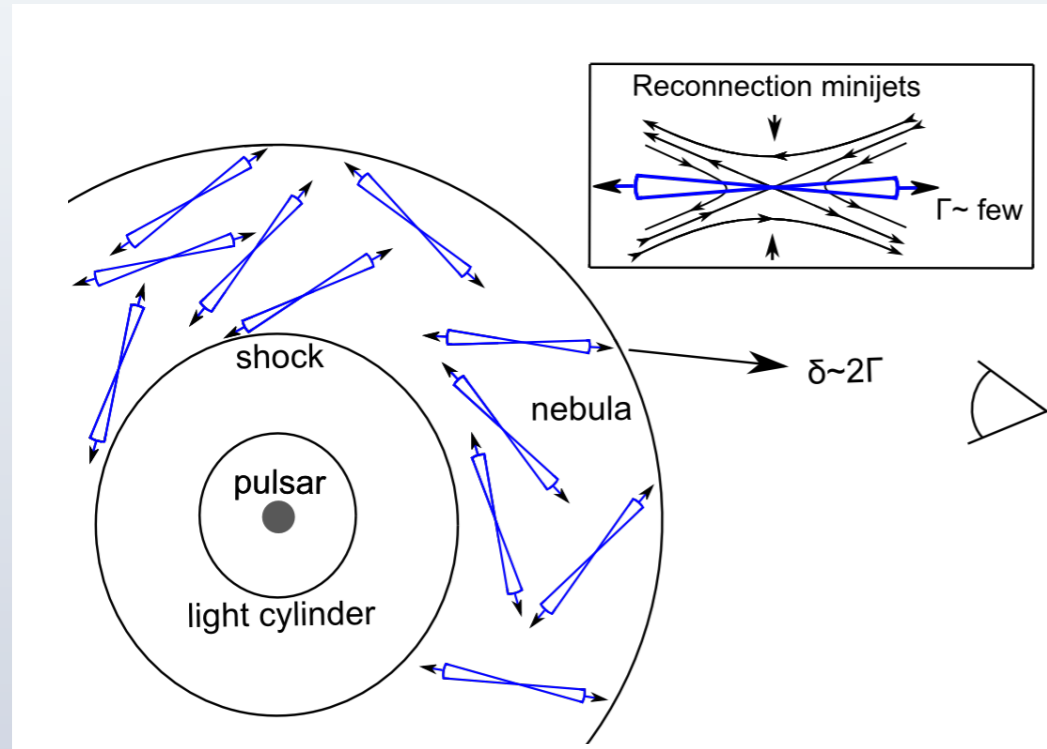


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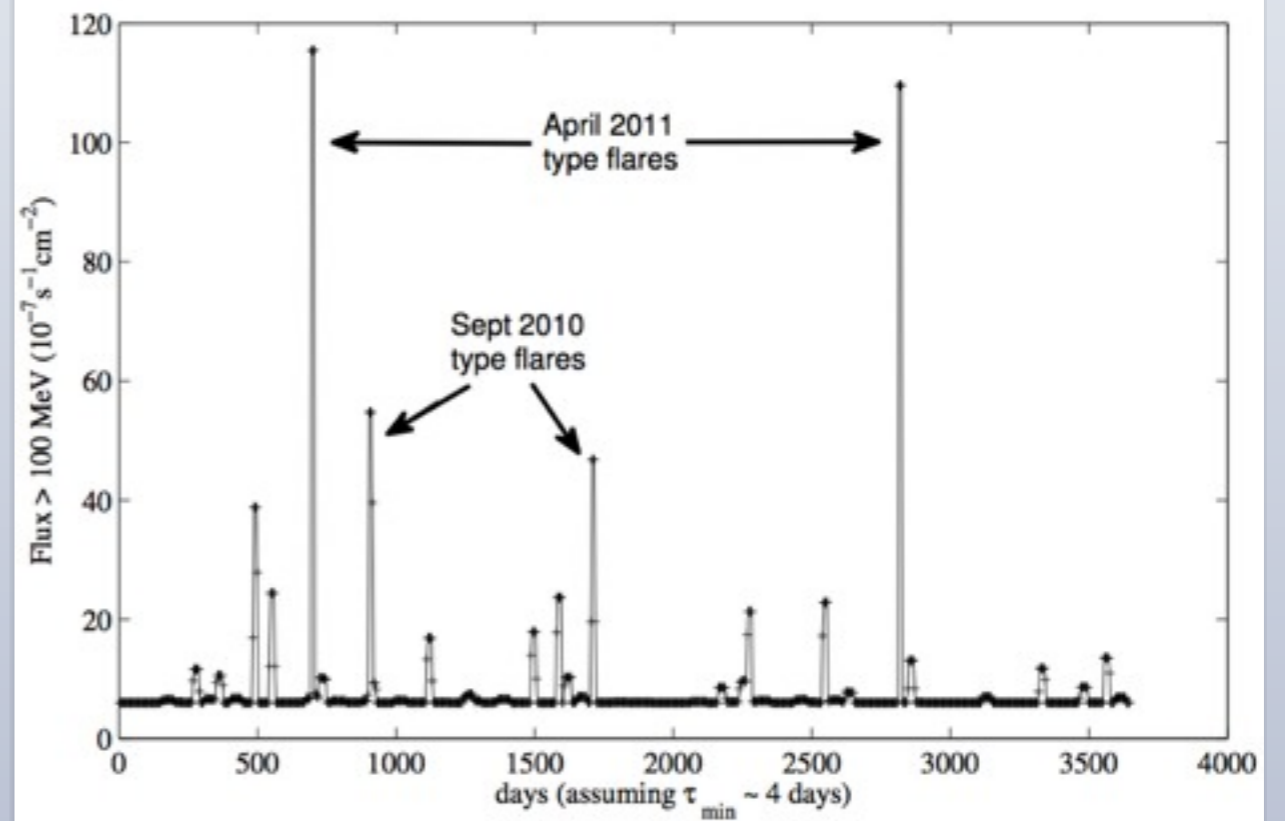


Mini-jets in Crab

Clausen-Brown, Lyutikov 2012



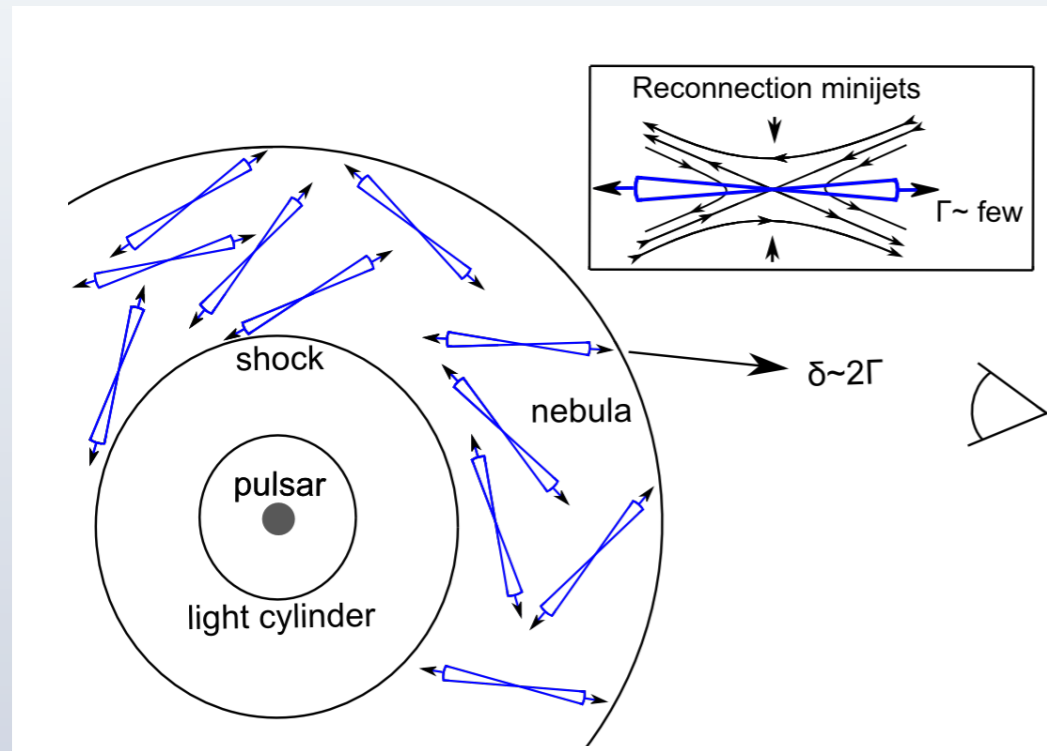
- probability of flare flux $\rho(F)dF \propto F^{-\frac{q+1}{q}} dF \approx \frac{1}{F} dF$
- average flare flux is dominated by bright rare flares.



Power-law from shot noise!

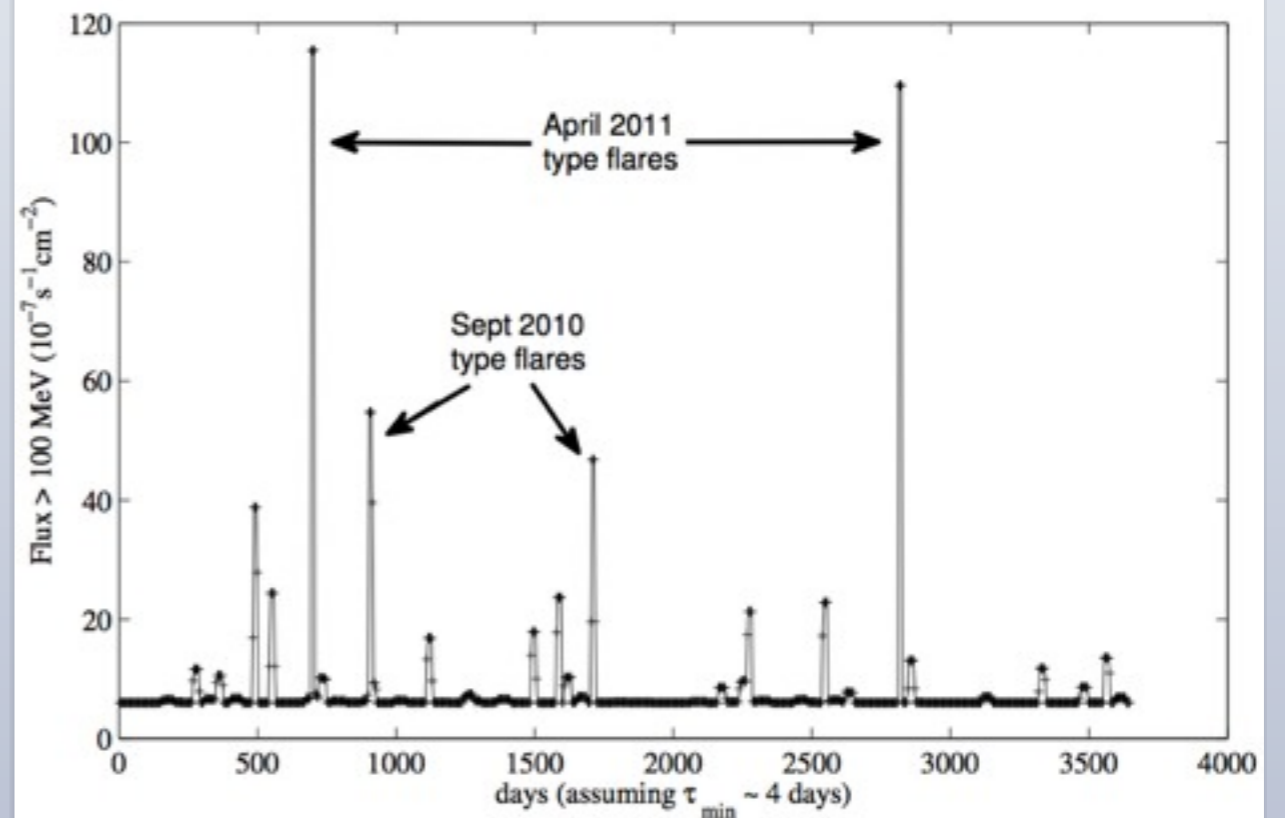
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Clausen-Brown, Lyutikov 2012



- Crab flares are an example how magnetic reconnection
 - can produce bursty radiation
 - can accelerate particles up to the radiation reaction limit, that radiate efficiently (needed number of leptons produced by Crab in only 1 sec)

- probability of flare flux $\rho(F)dF \propto F^{-\frac{q+1}{q}} dF \approx \frac{1}{F} dF$
- average flare flux is dominated by bright rare flares.



Power-law from shot noise!

Fast variability from large radii,

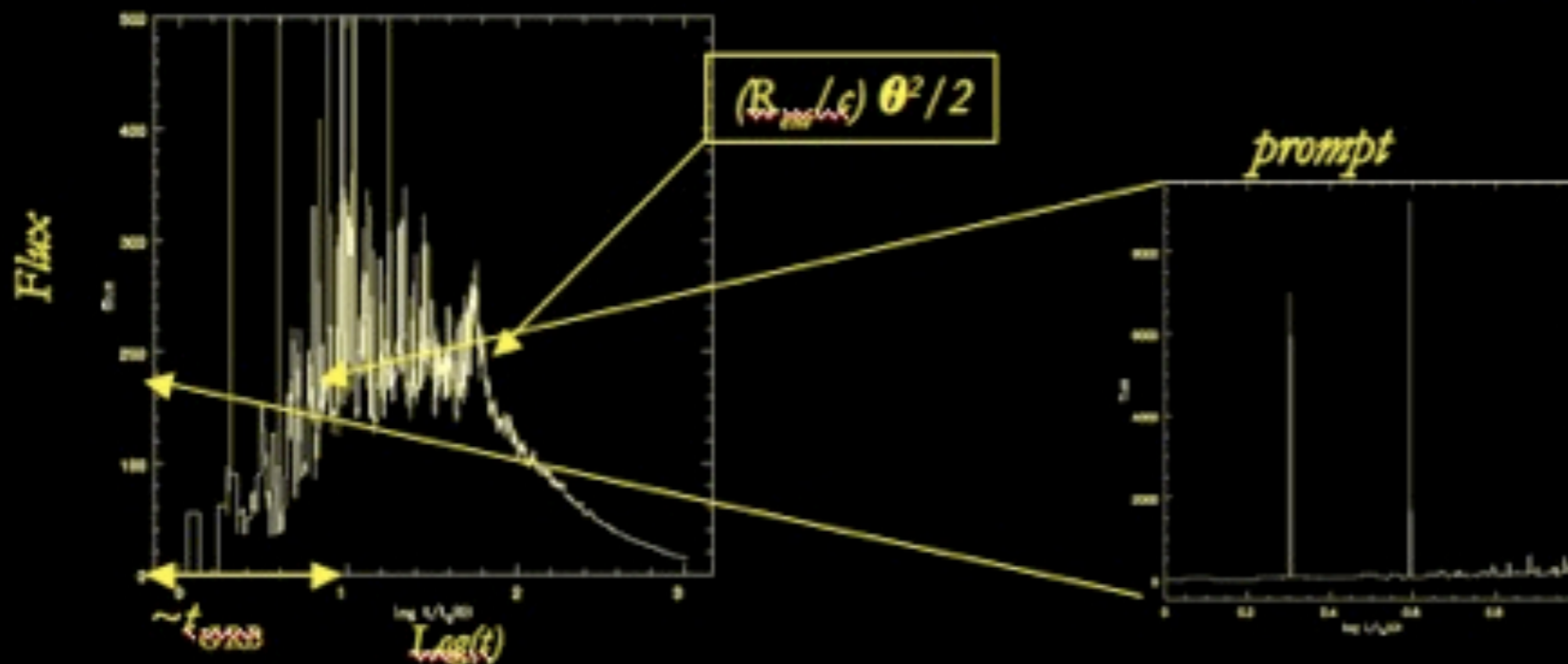
$$R_{em} \sim 10^{15} - 10^{16} \text{ cm}$$

(Lyutikov & Blandford 2003)

- Emission is beamed in outflow frame
 - really beamed $\Delta\theta_{em} \ll 1$
 - random internal motion of emitters, $\Delta\theta_{em} \sim 1/\gamma_{rand}$



- X-flares and breaks are tails of prompt
- fast variability
- no need for long central engine activity
- softening with time, harder spikes
- These are preliminary results: alternatives need to be investigated



$$\Gamma_{bulk} = 50, \gamma_{rand} = 5$$

$$\theta_{ob} = \pi/10,$$

$$\theta_{jet} = 3 \theta_{ob}$$

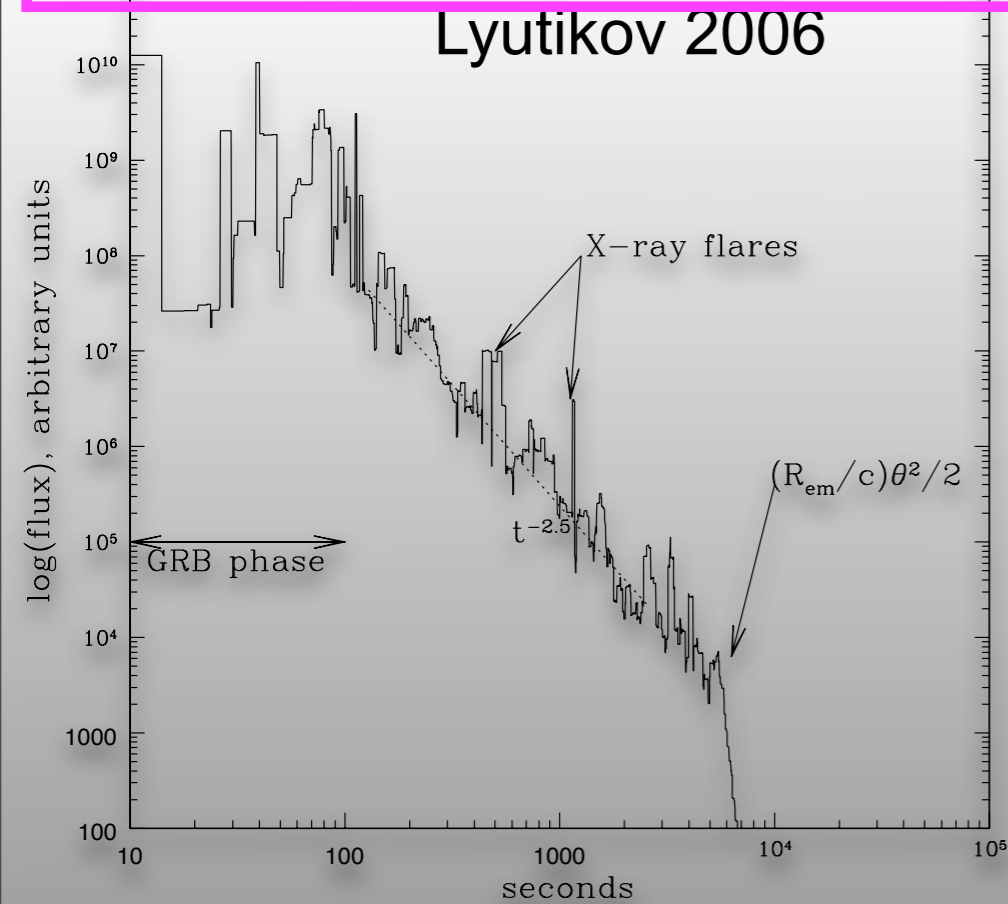
efficiency 10%

(Lyutikov in prog.)

Slide from Venice 2010 GRB meeting

Lyutikov 2006

Also: Ghisellini et al. 2008, Lazar et al. 2009, Giannios et al. 2009, Narayan & Kumar 2009



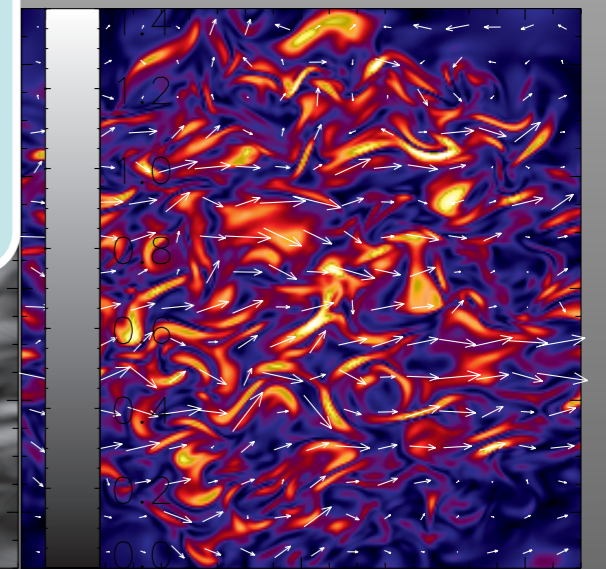
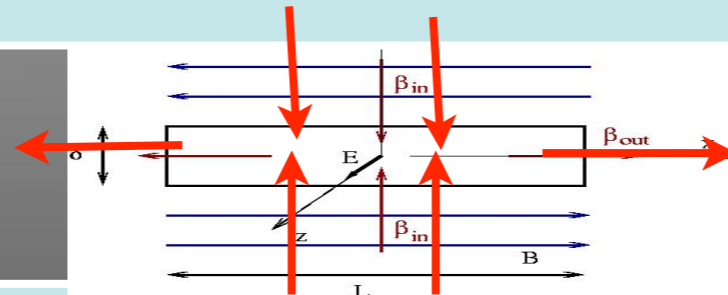
$$\Gamma_{eff} = 2\Gamma\gamma_{rand}$$

$$\Delta t \sim \frac{c}{R} \frac{1}{8\Gamma^2\gamma_{rand}^2}$$

Observed emission can be highly variable and with high efficiency (tapping into most of the proper volume)

- Not fluid "turbulence", $\gamma_{rand} \sim \sqrt{9/8} = 1.06$
- RM & RT instabilities will produce $vT \ll c$ turbulence

- Relativistic reconnection: jets with $\gamma_{out} \sim \sigma \gg 1$ (Lyutikov & Uzdenski 2004)

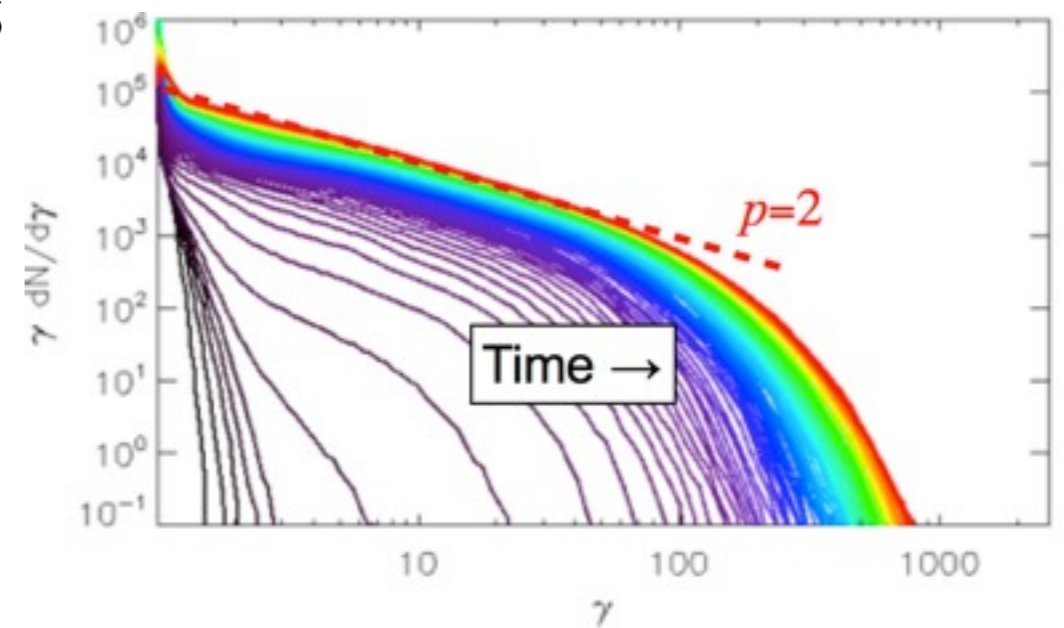


Turbulent reconnection (Lazarian & Vishniac)

- Spectrum is harder during flare (Burrows et al 2005)
- Are flares becoming longer and softer as function of flare time?
- Can some "Shorts" be "one spike Long"? (failed SN-type)
- Can explain optical -gamma correlations in 080319B? E.g. emitting "blobs" expand, killing both

Particle acceleration in reconnection

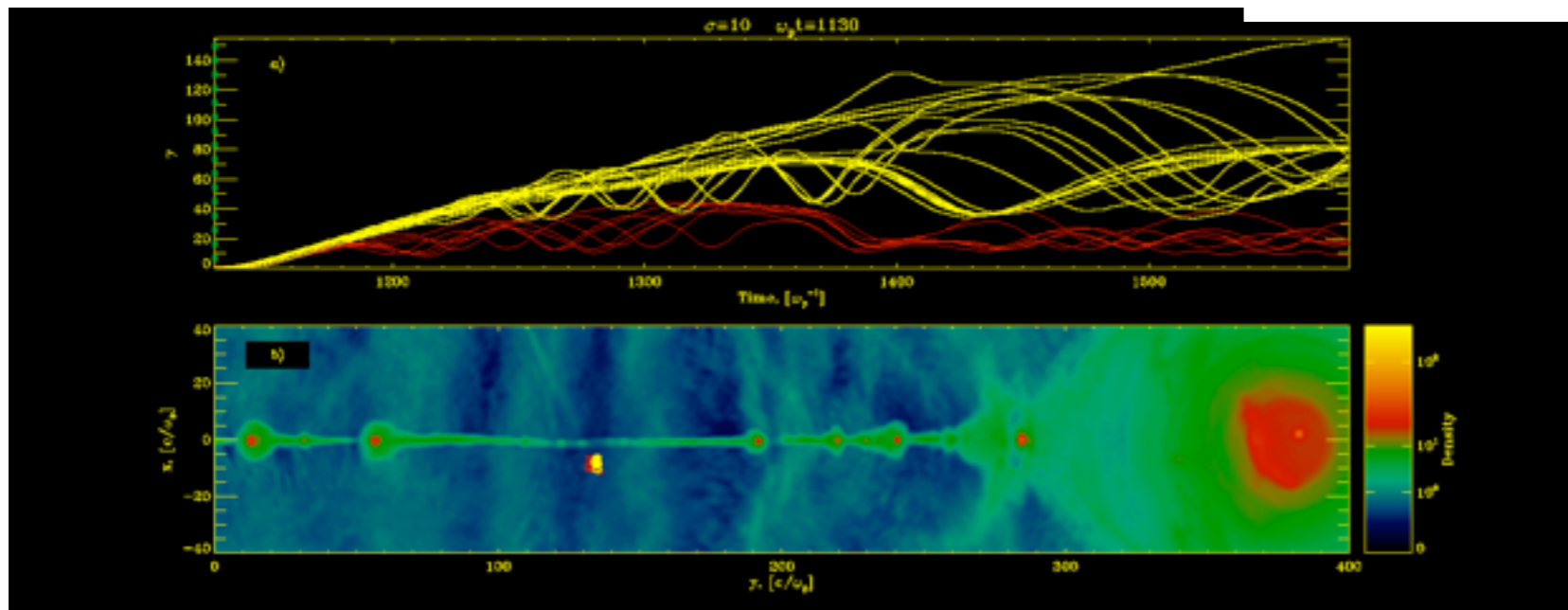
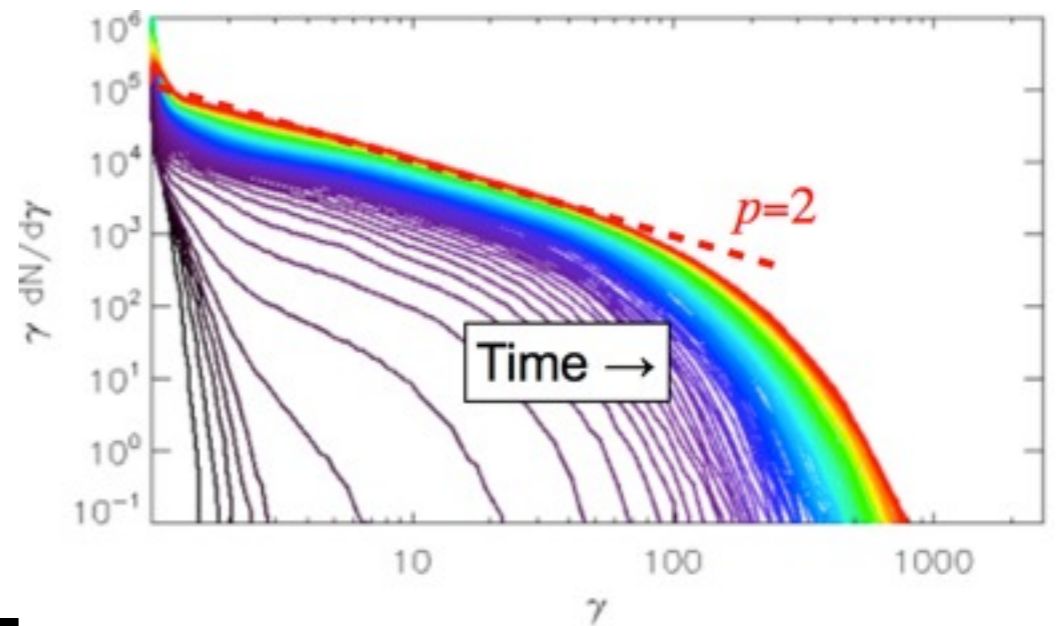
- Shock acceleration: correct kinetic spectrum of particles follows from macroscopic jump conditions
- Reconnection - no simple scaling...
- But same result!



Sironi

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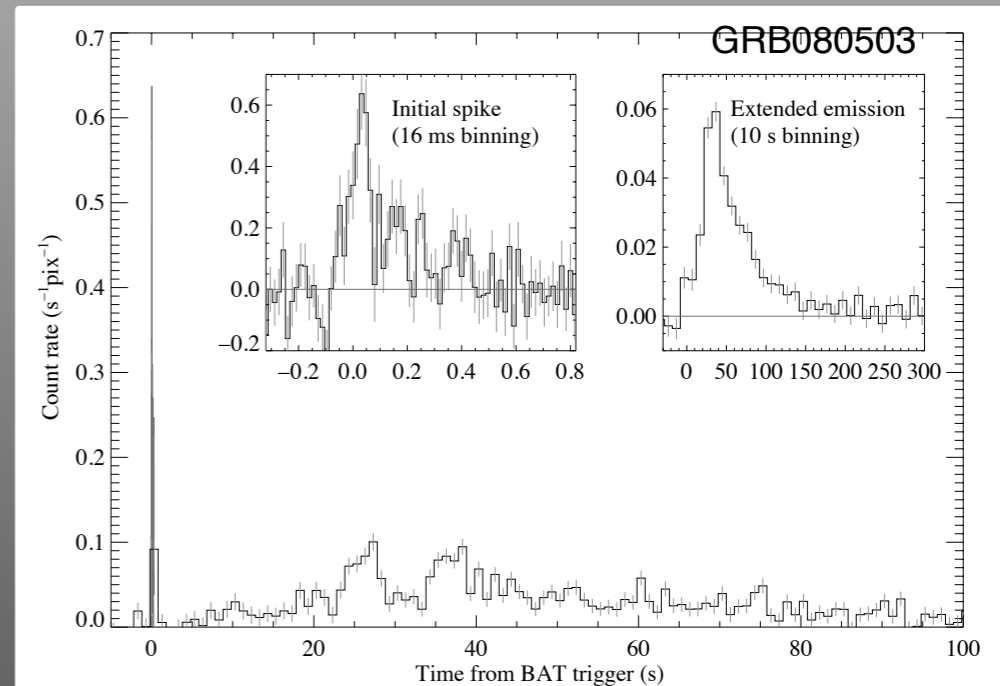
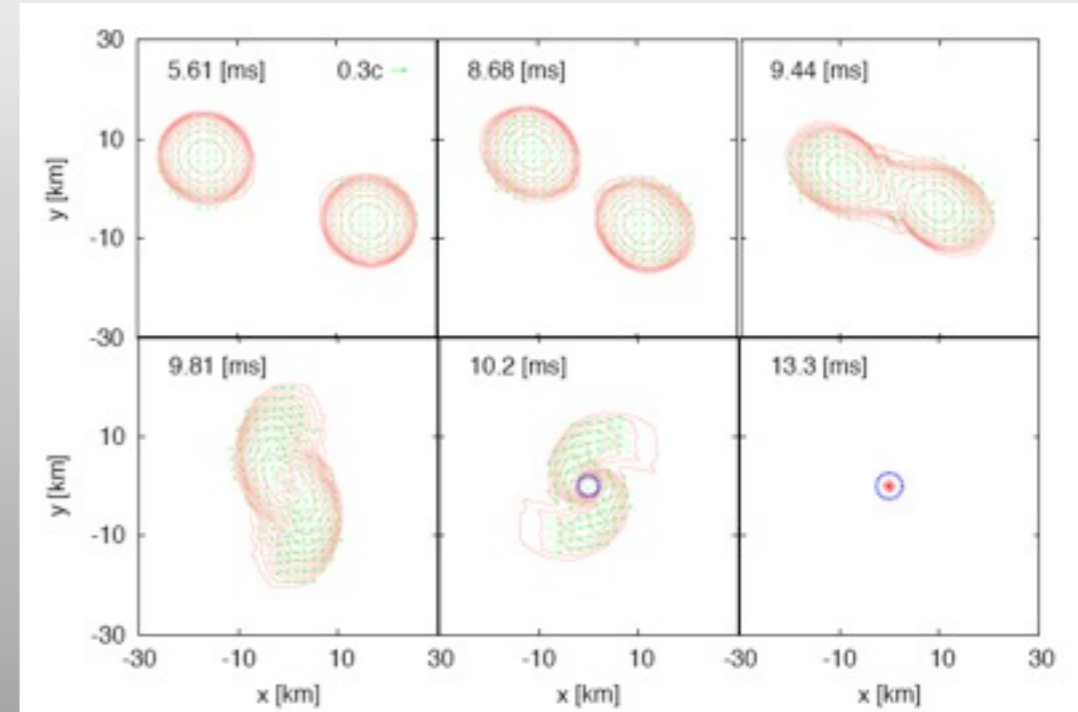


Sironi

Long-lived engines in short GRBs

NS-NS merger as paradigm for Short GRBs

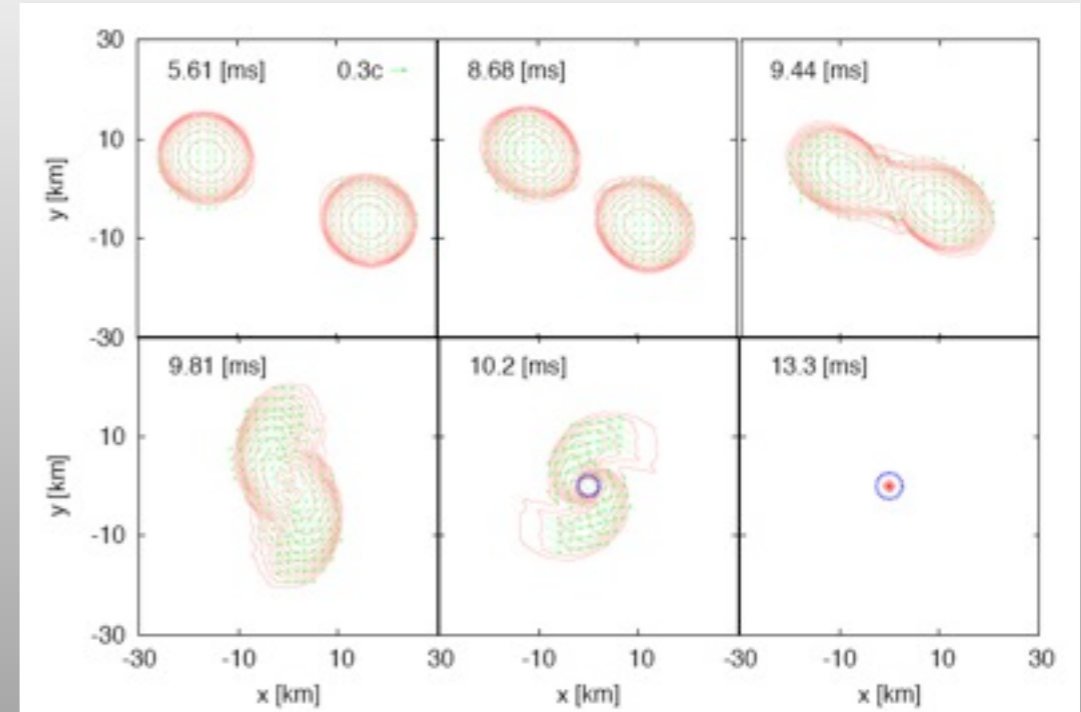
- Active stage of NS-NS merger takes 10-100 msec, then **collapse into BH**
- Transient NS - 100 msec, (**NOT** 100 sec!)
- Very little mass is ejected, drains out quickly
- Many short GRBs have long 100 sec tails, energetically comparable to the prompt spike.
- Many GRBs have late time flares, 10^5 sec



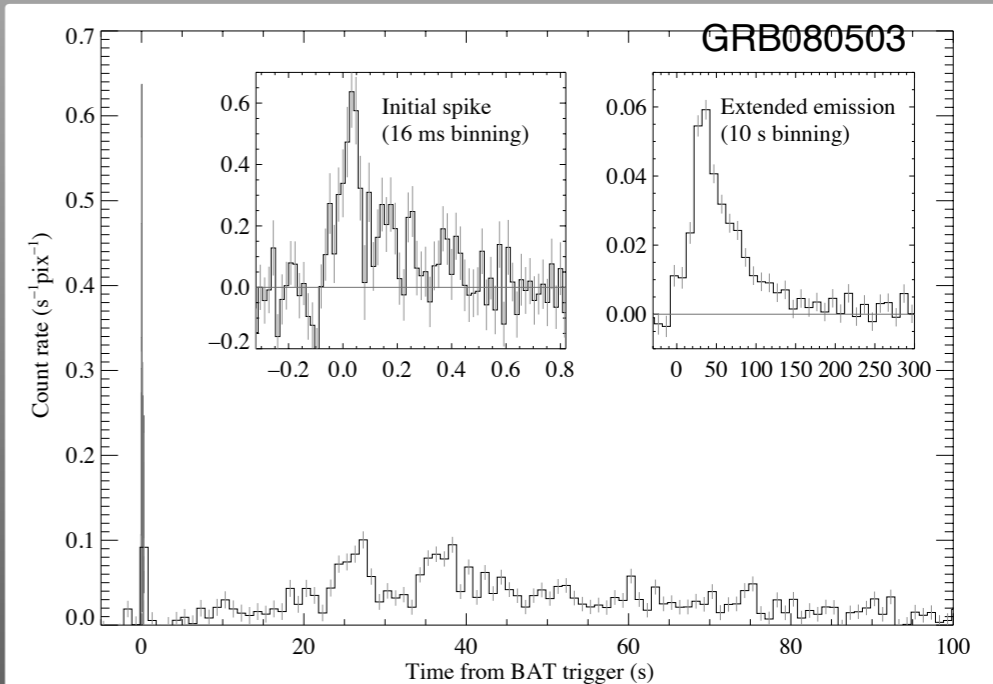
100 sec tail has ~ 30 times more energy than the prompt spike

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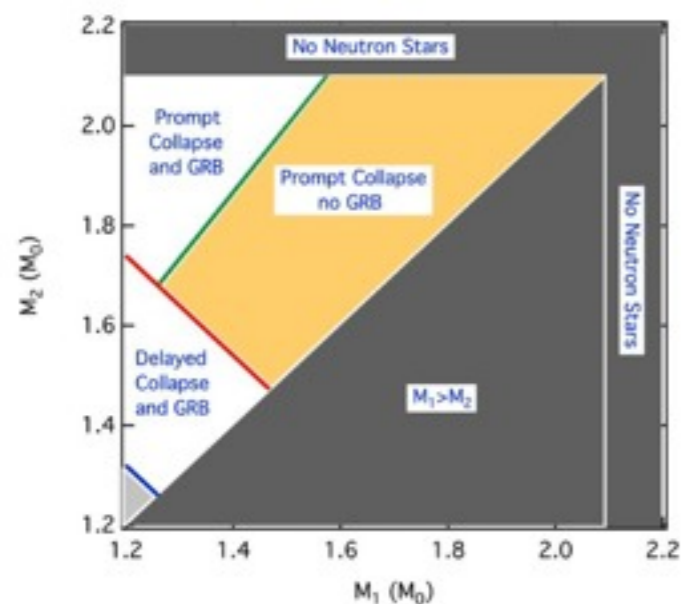
How to explain energetically dominant activity on ~ 100 sec, while the engine lives 10-100 msec?



100 sec tail has ~ 30 times more energy than the prompt spike

Can merger of two NSs leave a NS? (millisecond magnetar)

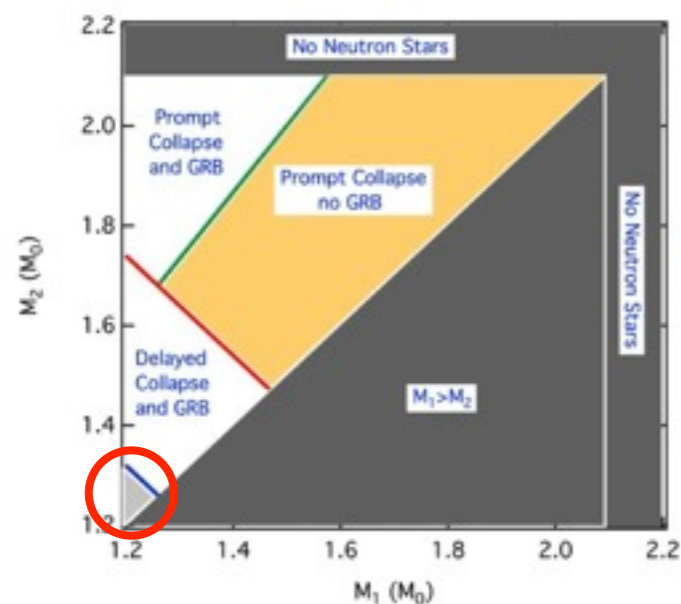
- There is a $2M_{\text{Sun}}$ NS
 - Need **both** $M_{\text{NS}} < 1.2 M_{\text{Sun}}$
 - And throw out $\sim 0.3 M_{\text{Sun}}$
 - And very stiff EoS



Ozel + 2010

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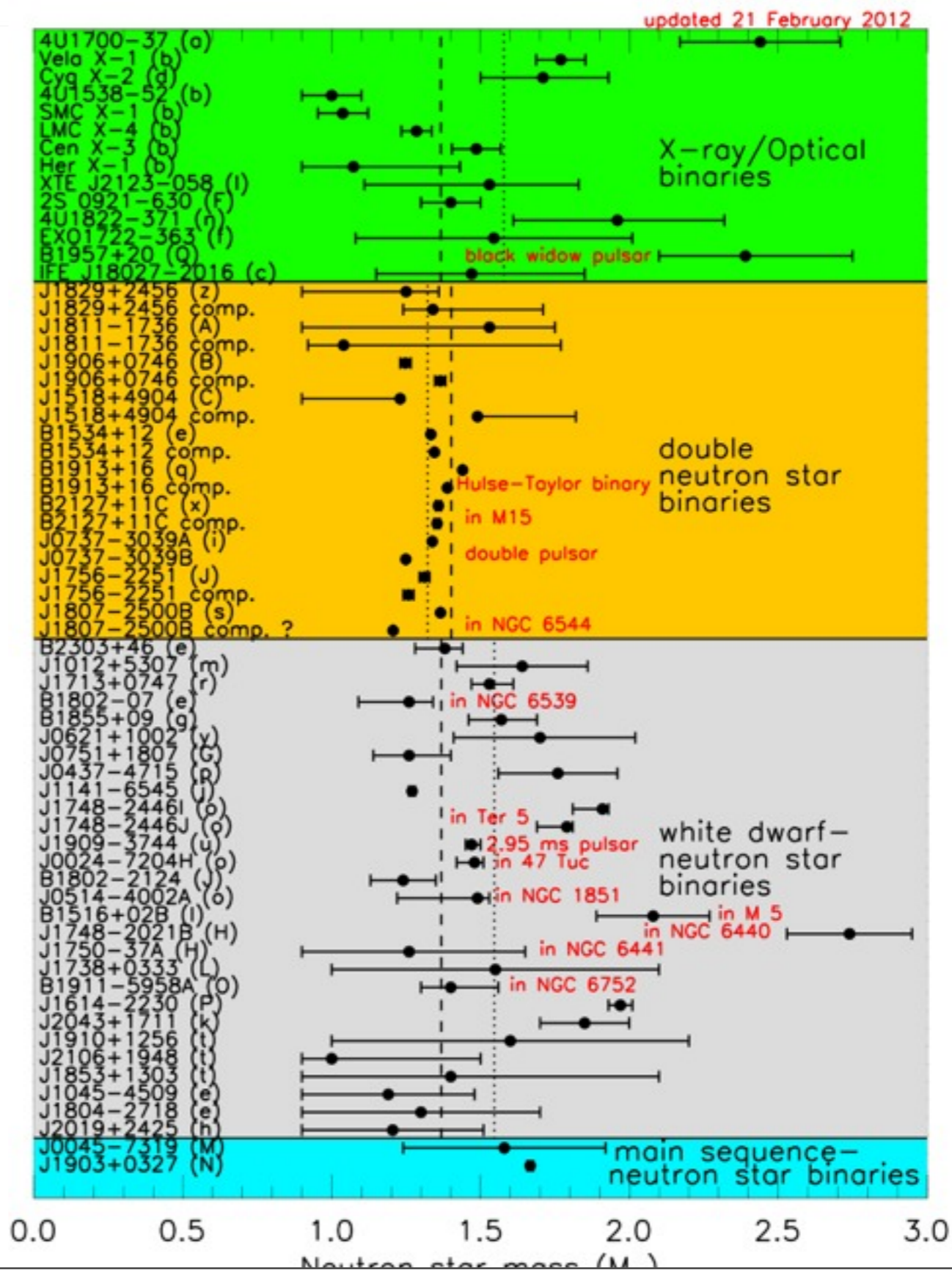
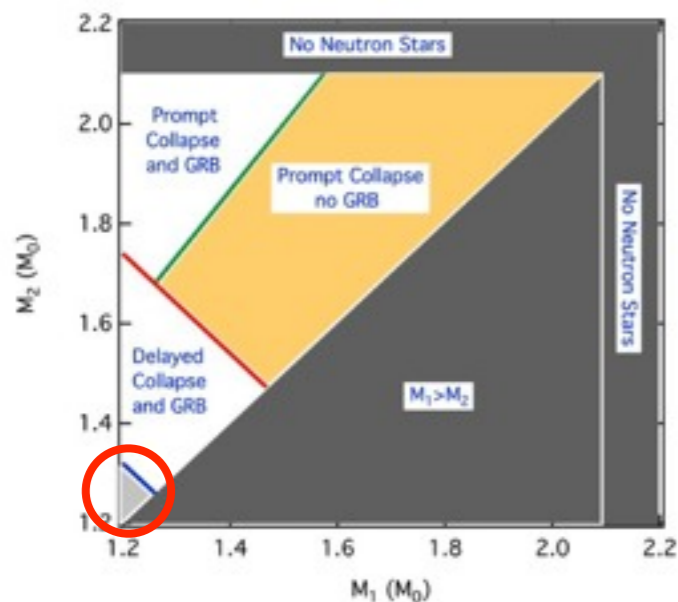


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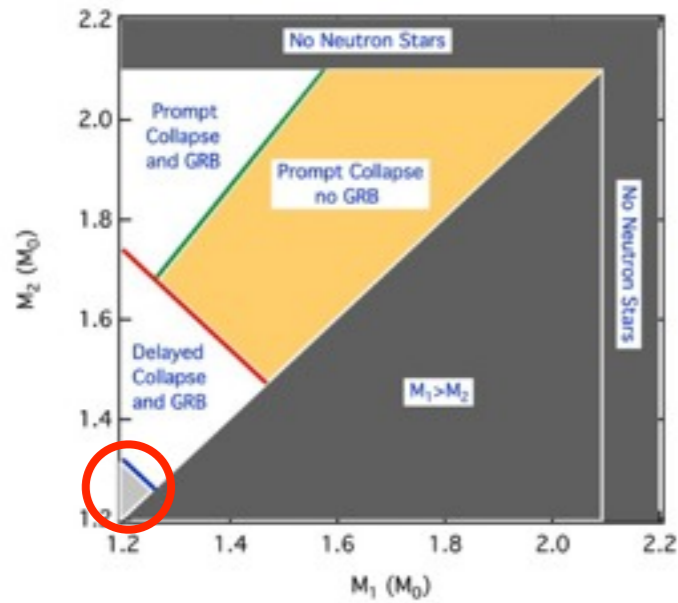
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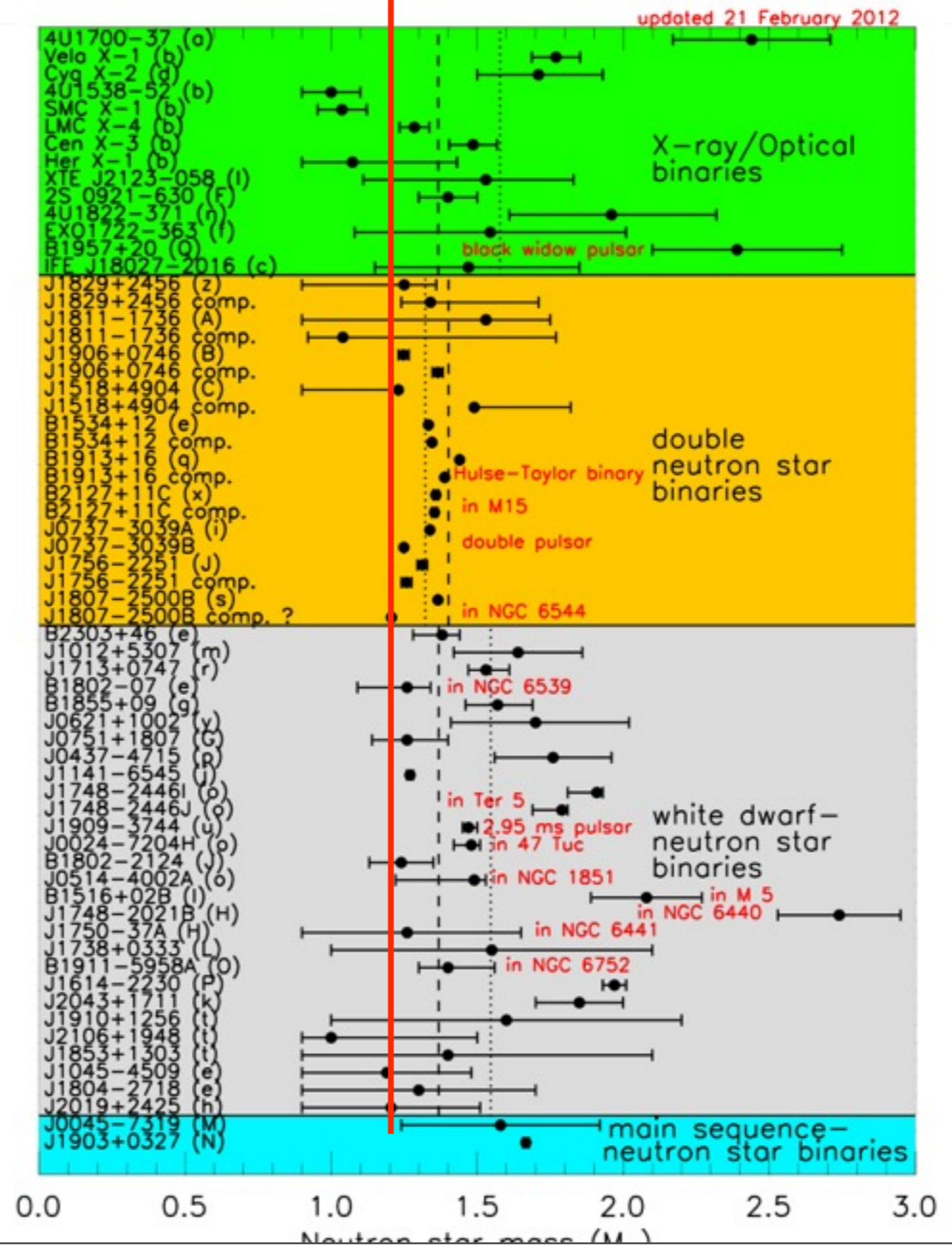


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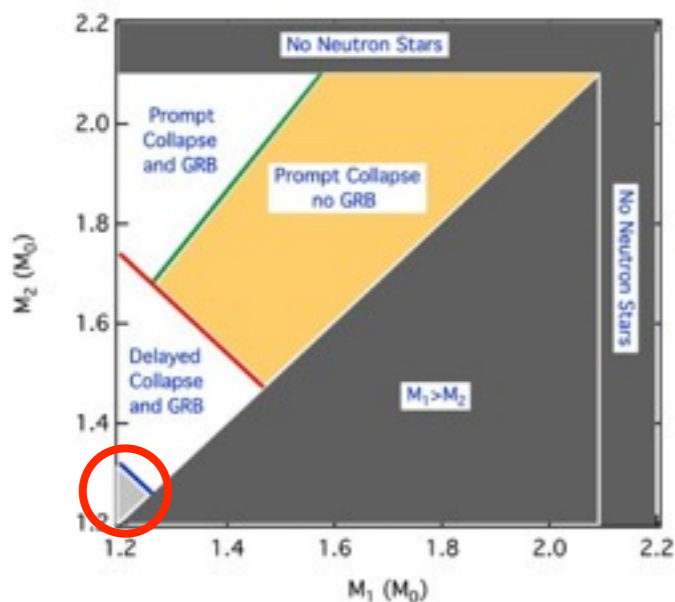


Ozel + 2010



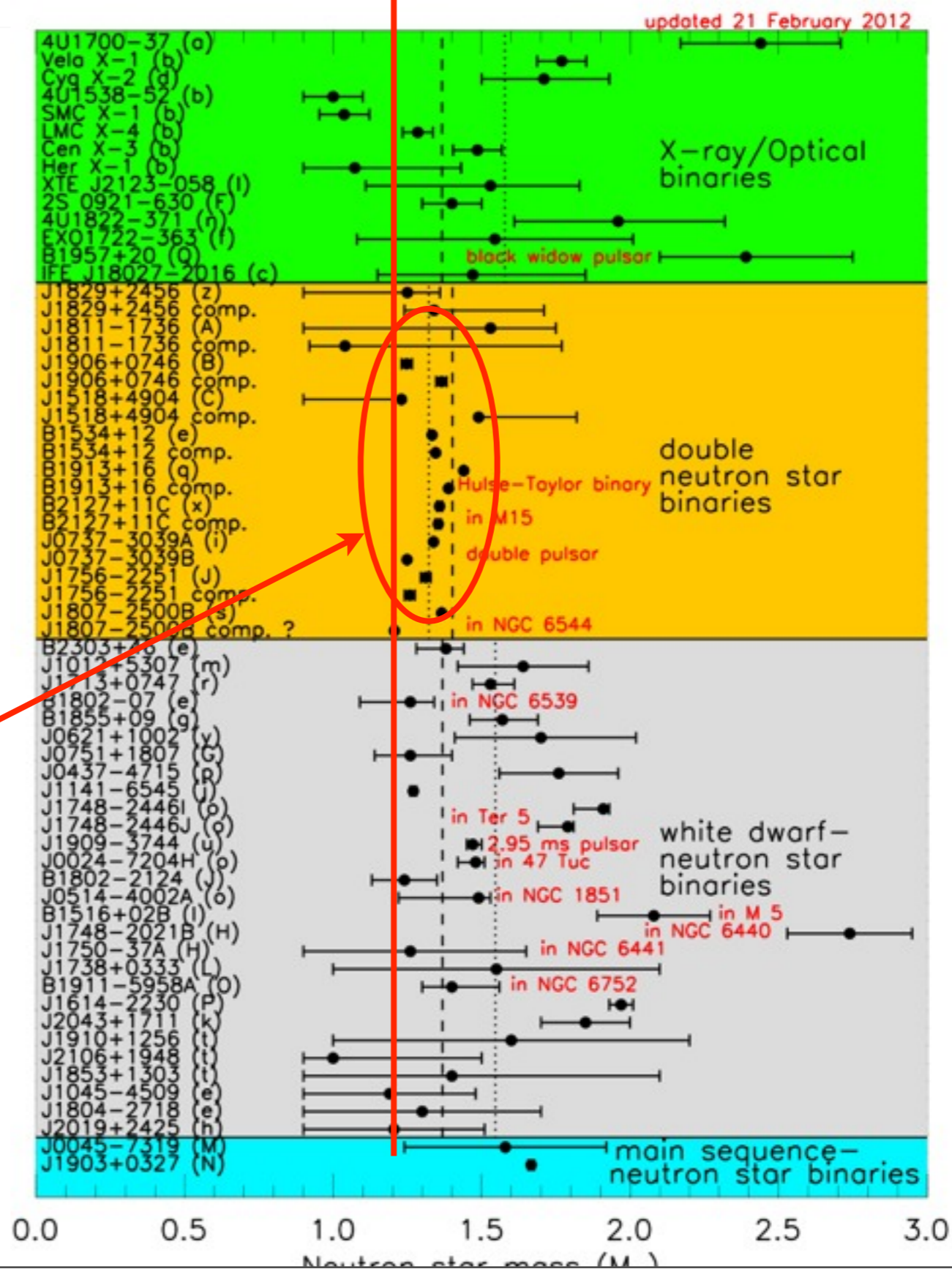
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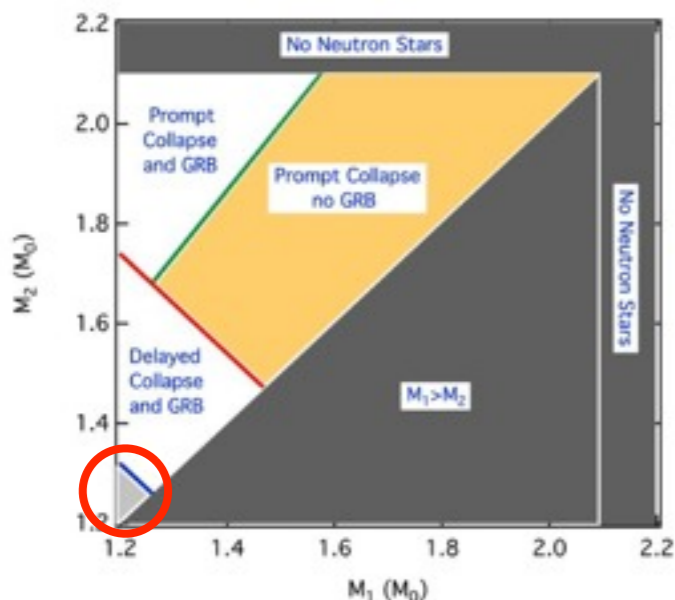
Ozel + 2010

The best determined masses (down to $10^{-4} M_{\text{Sun}}$) are in NS-NS binaries,
 $M_{\text{min}} = 1.25 M_{\text{Sun}}$



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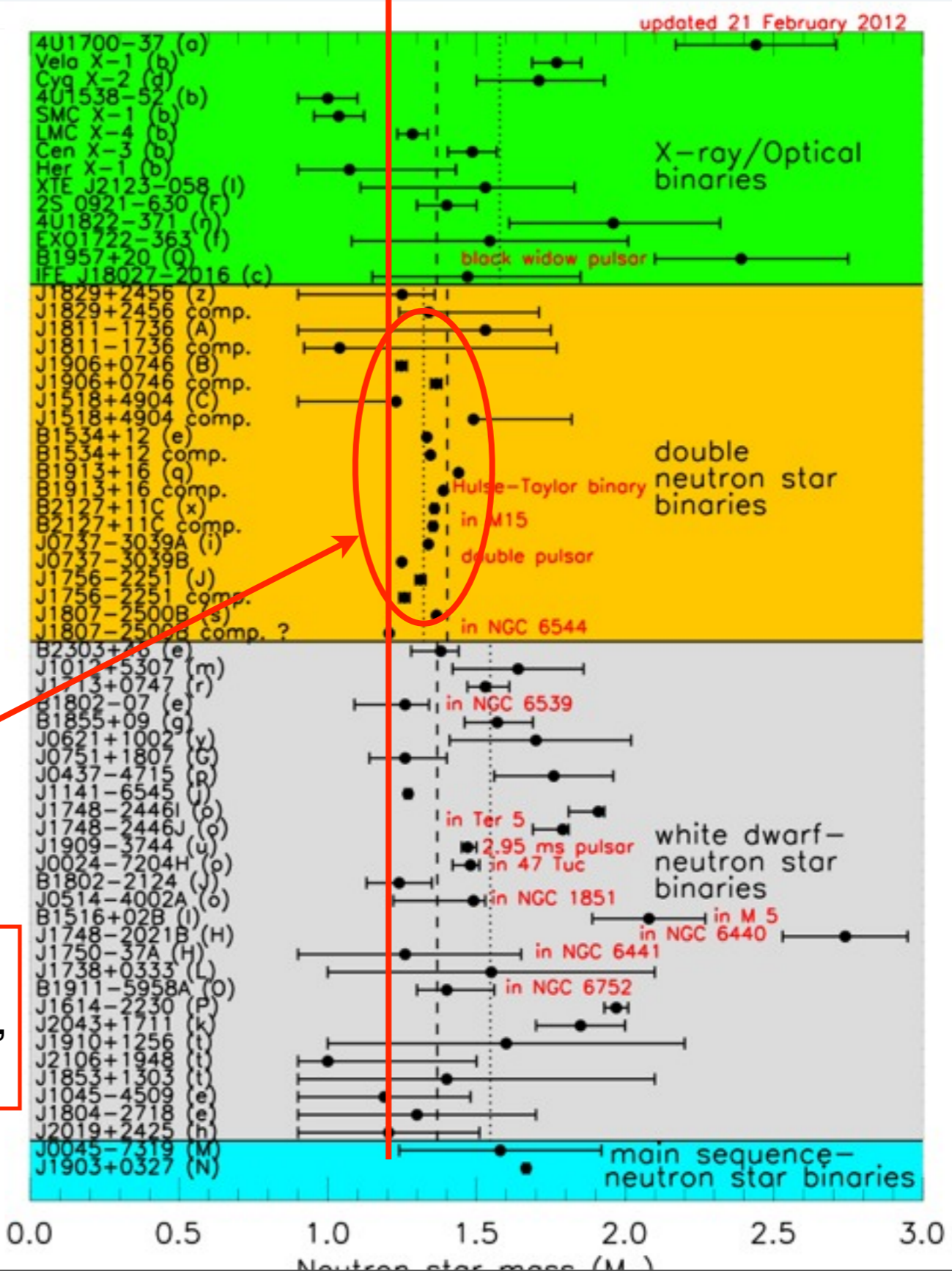


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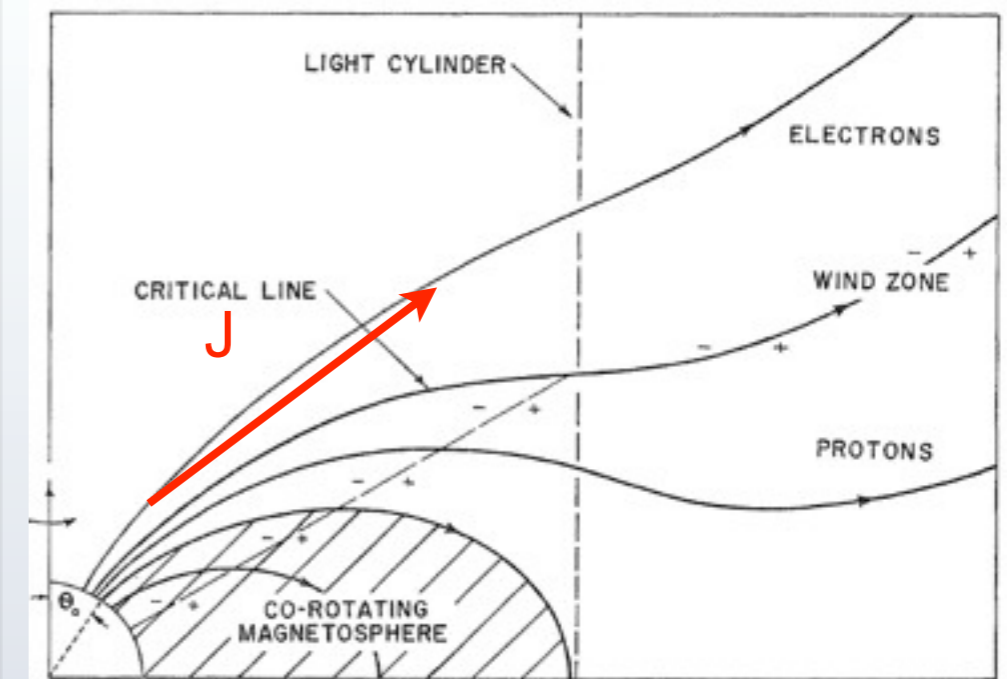
This cannot be the dominant channel of NS-NS mergers and, thus, of short GRBs.



BHs part-timing as magnetars

Newly formed **isolated spinning** astrophysical black holes can keep magnetic fields for times much longer than predicted by the “No hair” theorem, working as \sim millisecond magnetars

- **Rotating NS** - unipolar inductor
 - generate plasma out of vacuum
 - open B-field lines to infinity
- Blandford & Znajek: BHs do the same
- Outside plasma: $\mathbf{E} \cdot \mathbf{B} = 0$ - frozen-in B-field
- If a BH keeps producing plasma, like a NS, B-field cannot slide off: **field lines that connected NS surface to infinity, has to connect horizon to infinity**



Goldreich & Julian, 1969

- **The “no hair” theorem is not applicable to collapse of rotating NSs: high plasma conductivity introduces topological constraint (frozen-in B-field).**

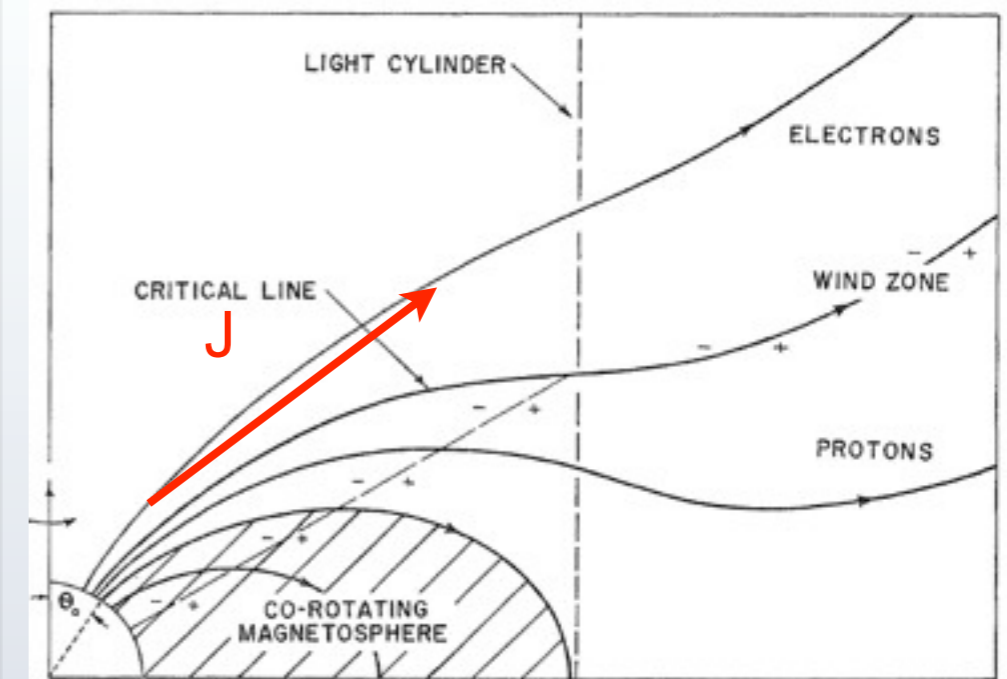
Conserved number: open magnetic flux:

$$N_B = e\Phi_\infty / (\pi c \hbar)$$

$$\Phi_\infty \approx 2\pi^2 B_{NS} R_{NS}^3 / (P_{NS} c)$$

Can be measured at infinity: BH hair

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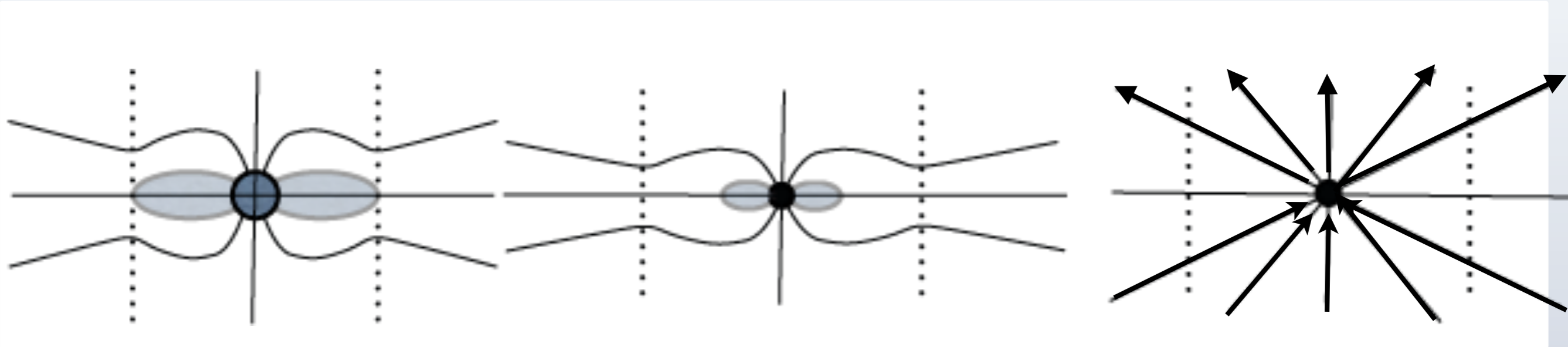
Countable BH hair!

$$N_B = e\Phi_\infty / (\pi c \hbar)$$

$$\Phi_\infty \approx 2\pi^2 B_{NS} R_{NS}^3 / (P_{NS} c)$$

Can be measured at infinity: BH hair

BH can only have open field lines: split monopole magnetosphere

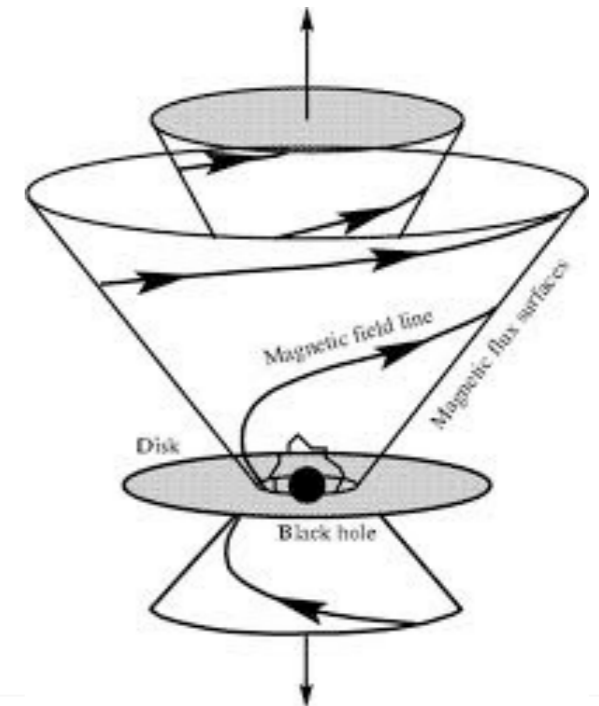


- Analytics: time-dependent force-free B-field in Schwarzschild geom.

$$B_\phi = -\frac{R_s^2 \Omega \sin \theta}{\alpha r} B_s, \quad B_r = \left(\frac{R_s}{r}\right)^2 B_s,$$

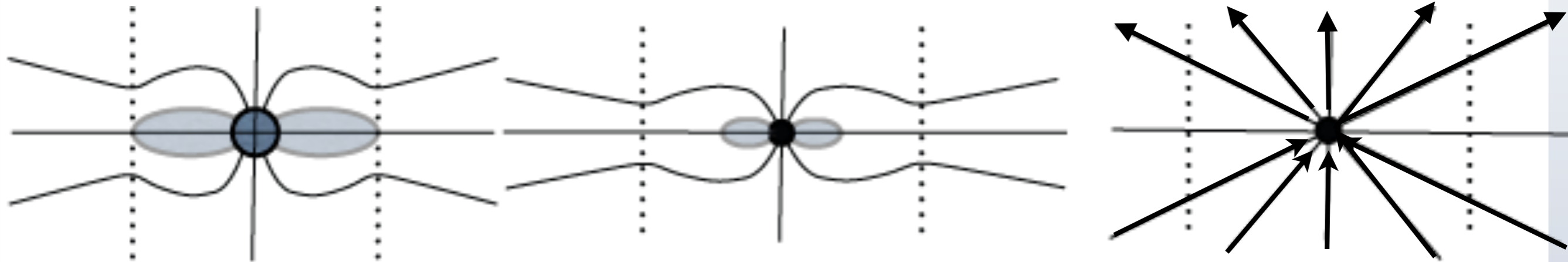
$$E_\theta = B_\phi, \quad j_r = -2 \left(\frac{R_s}{r}\right)^2 \frac{\cos \theta \Omega B_s}{\alpha}$$

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BH can only have open field lines: split monopole magnetosphere

pulsar

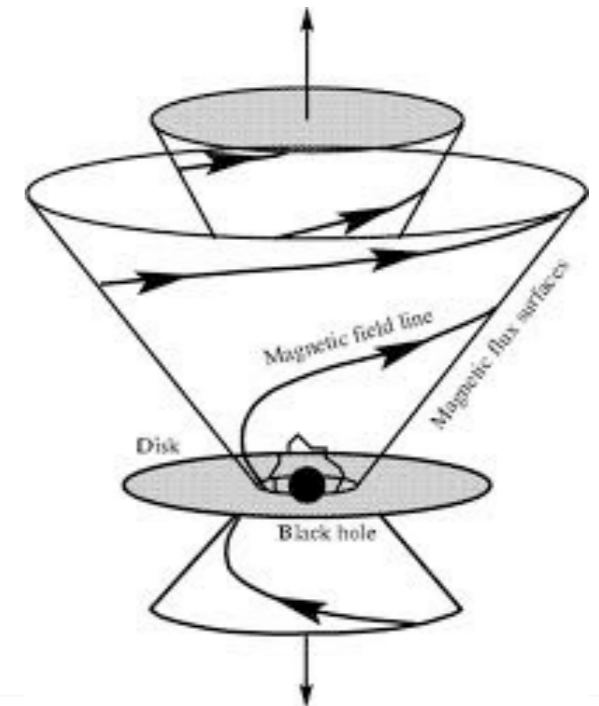


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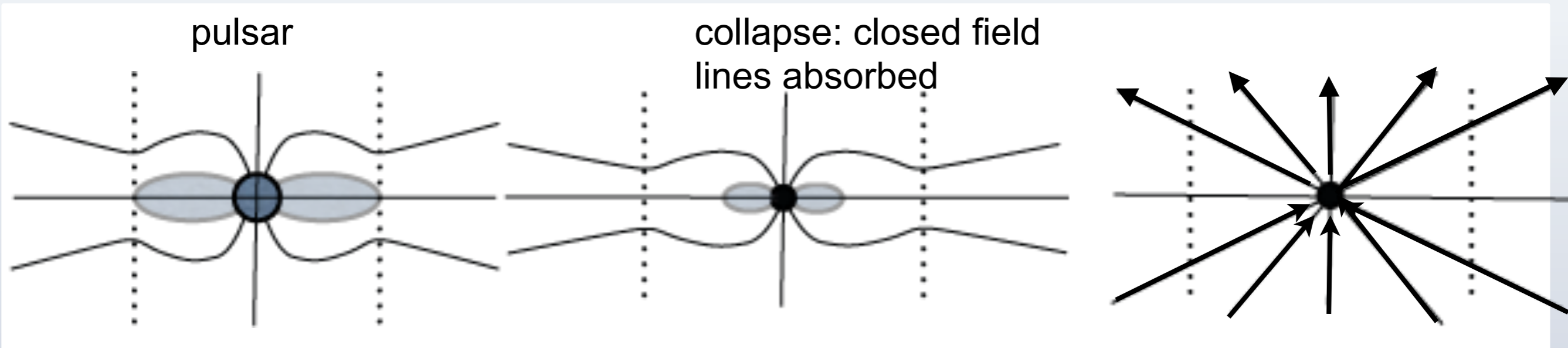
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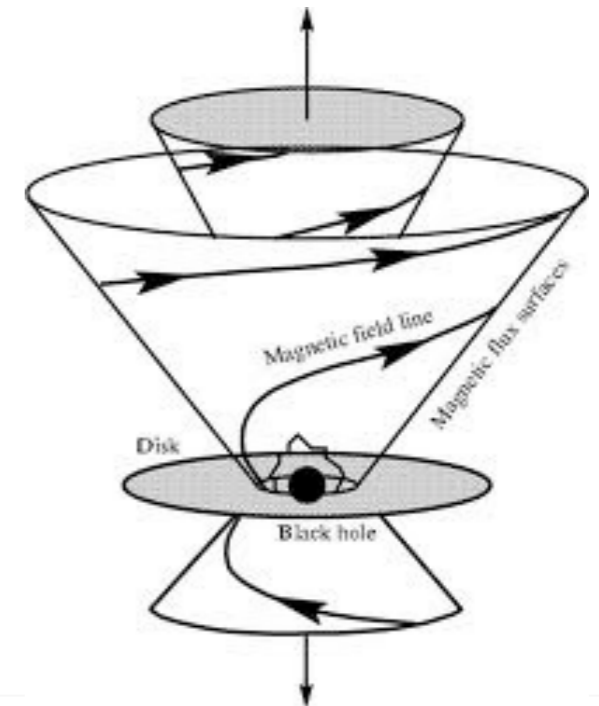


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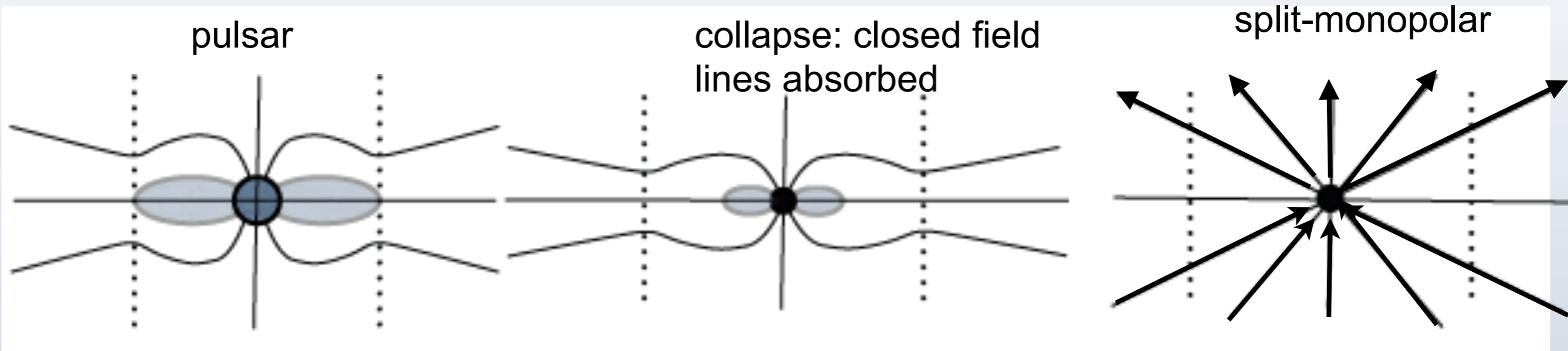
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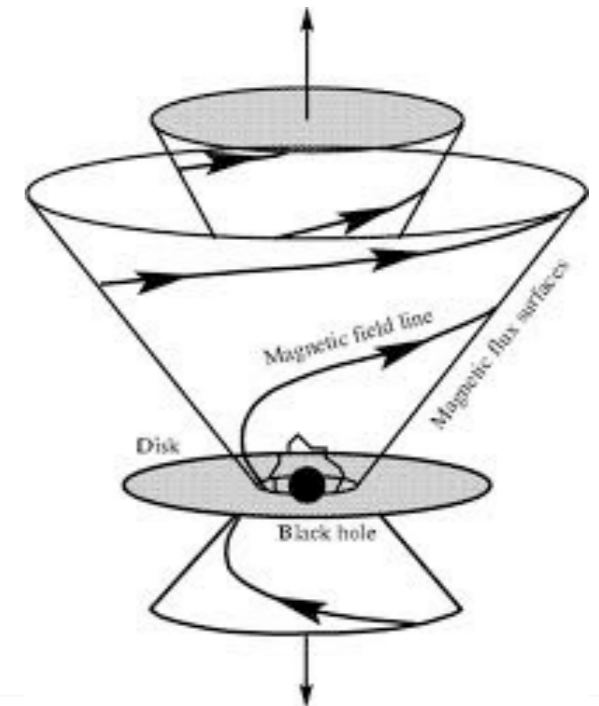


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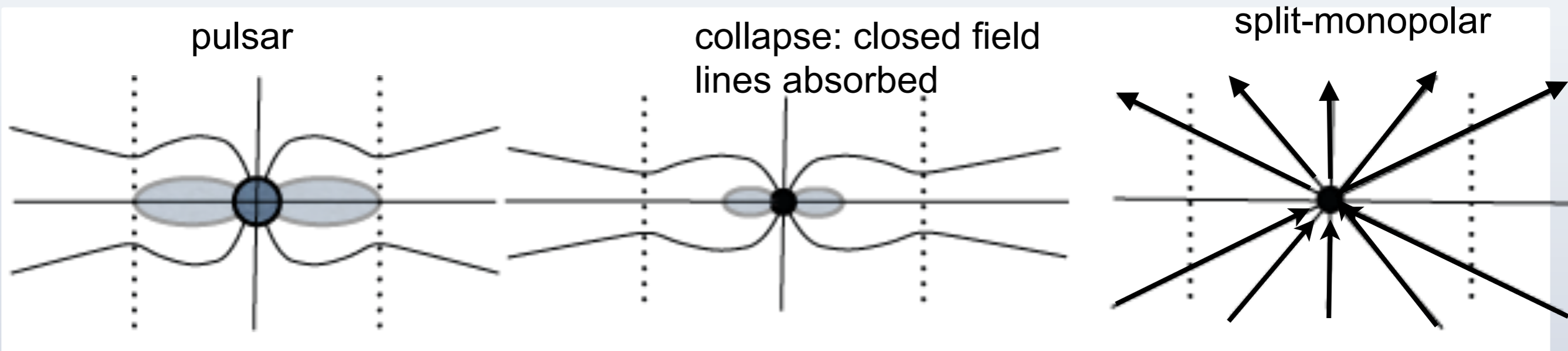
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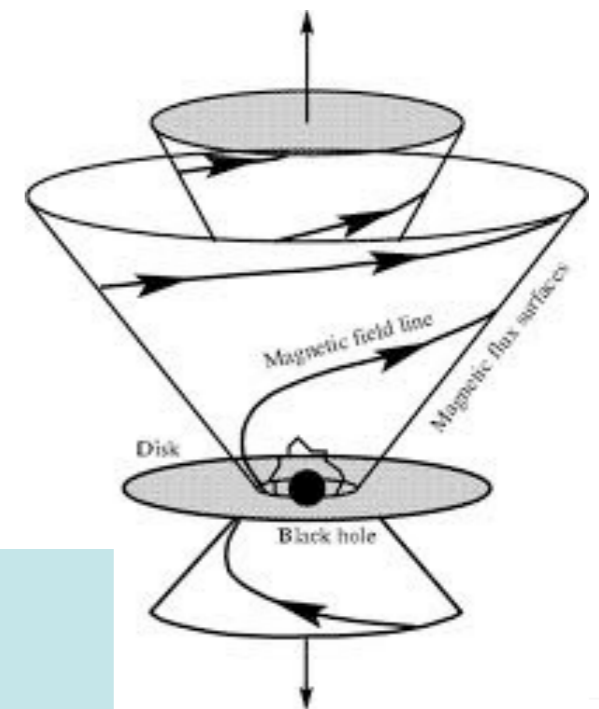
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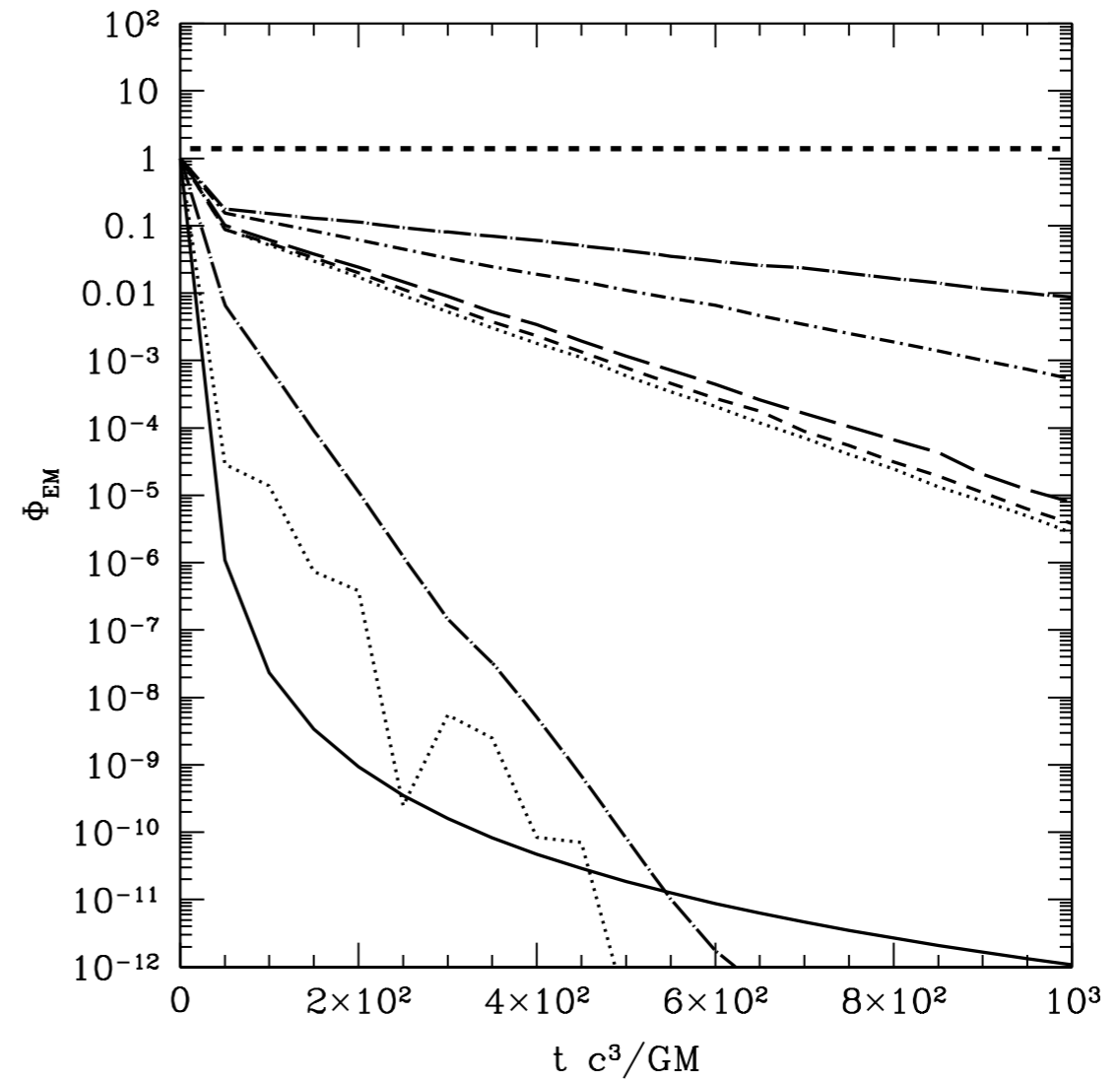
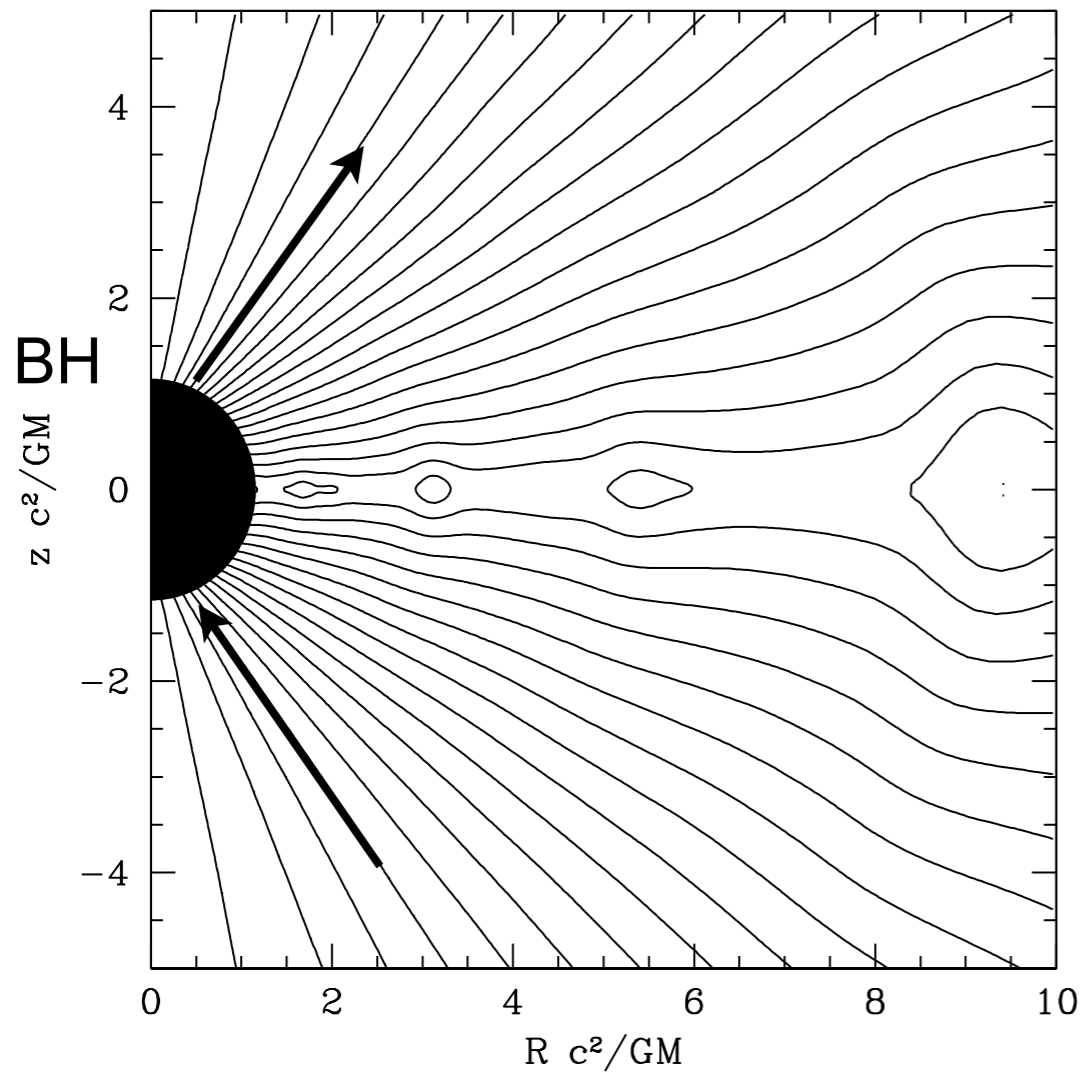
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Take a relativistic object with monopolar B-field, rotate it arbitrarily (slowly, $a \ll 1$). The field will remain monopolar



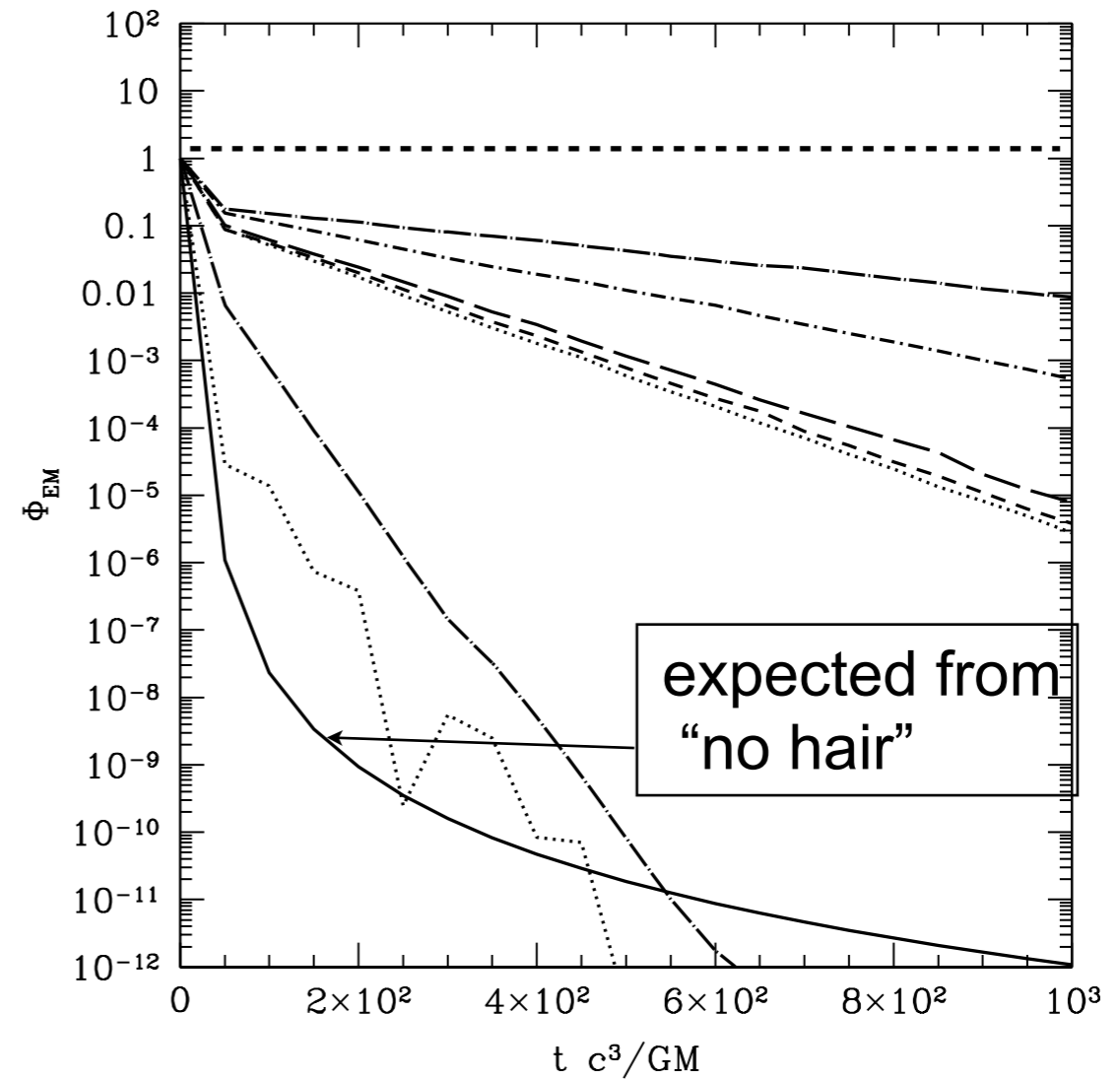
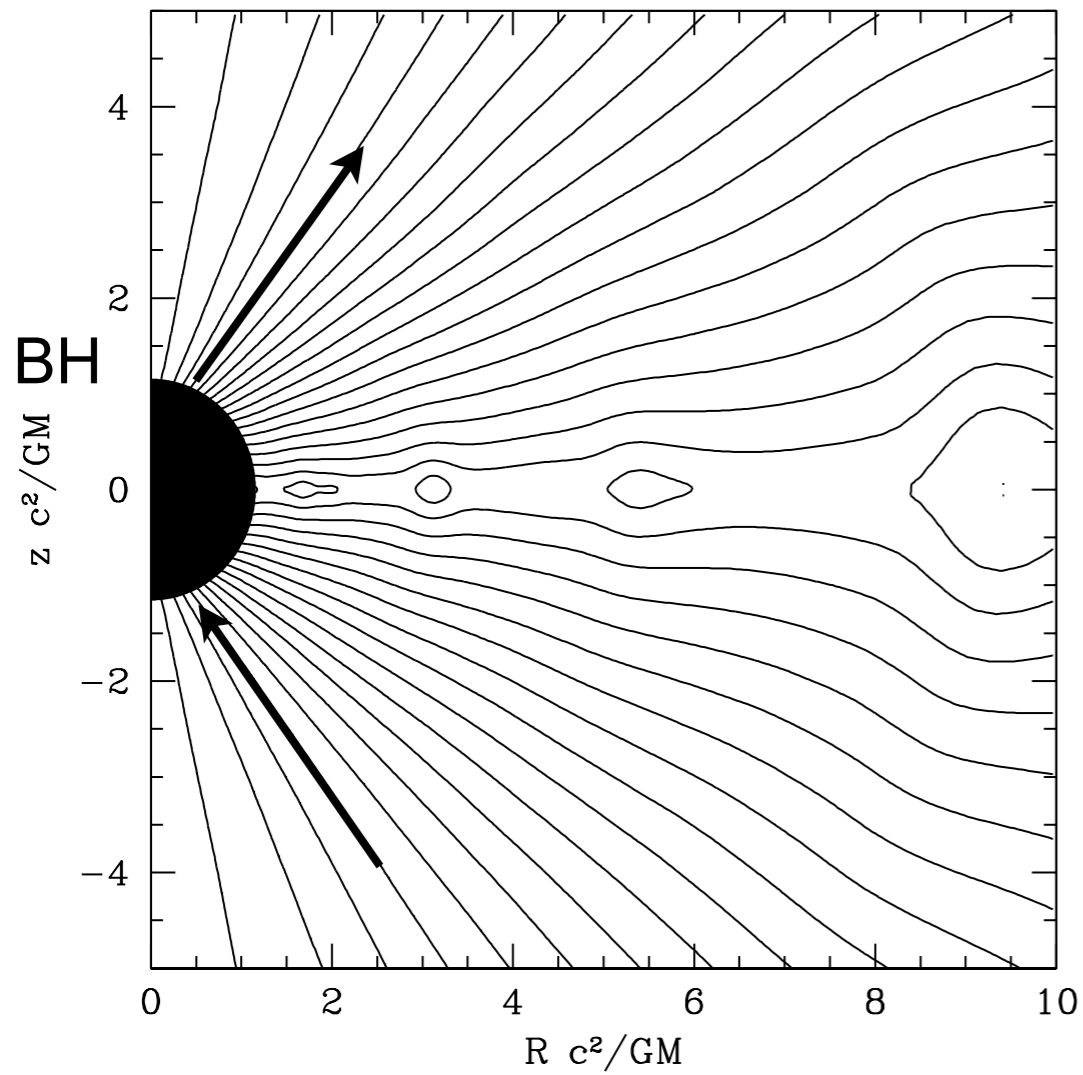
Simulations *(Lyutikov & McKinney, 2011)*

- Split-monopole
magnetosphere
- Slow balding



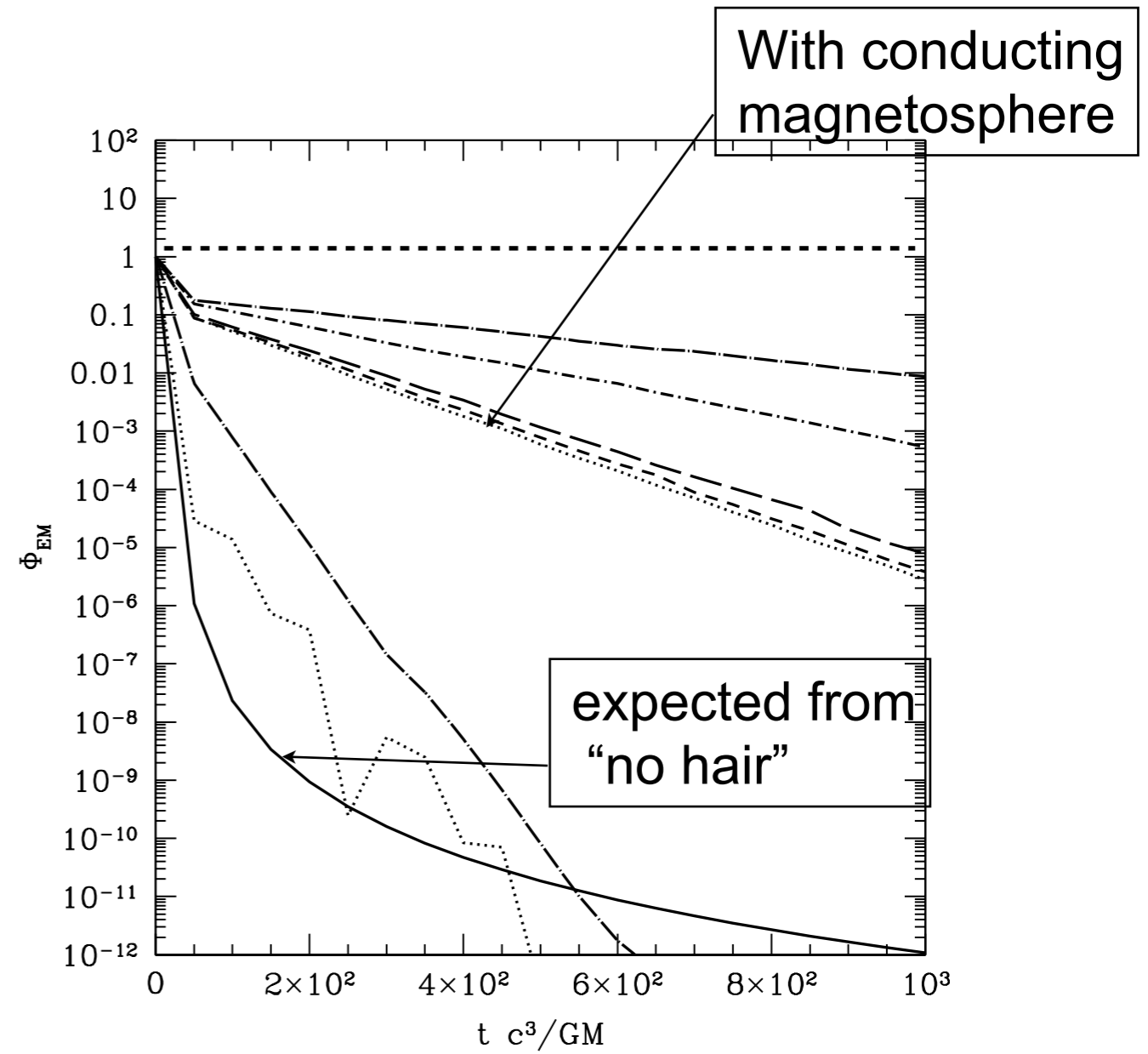
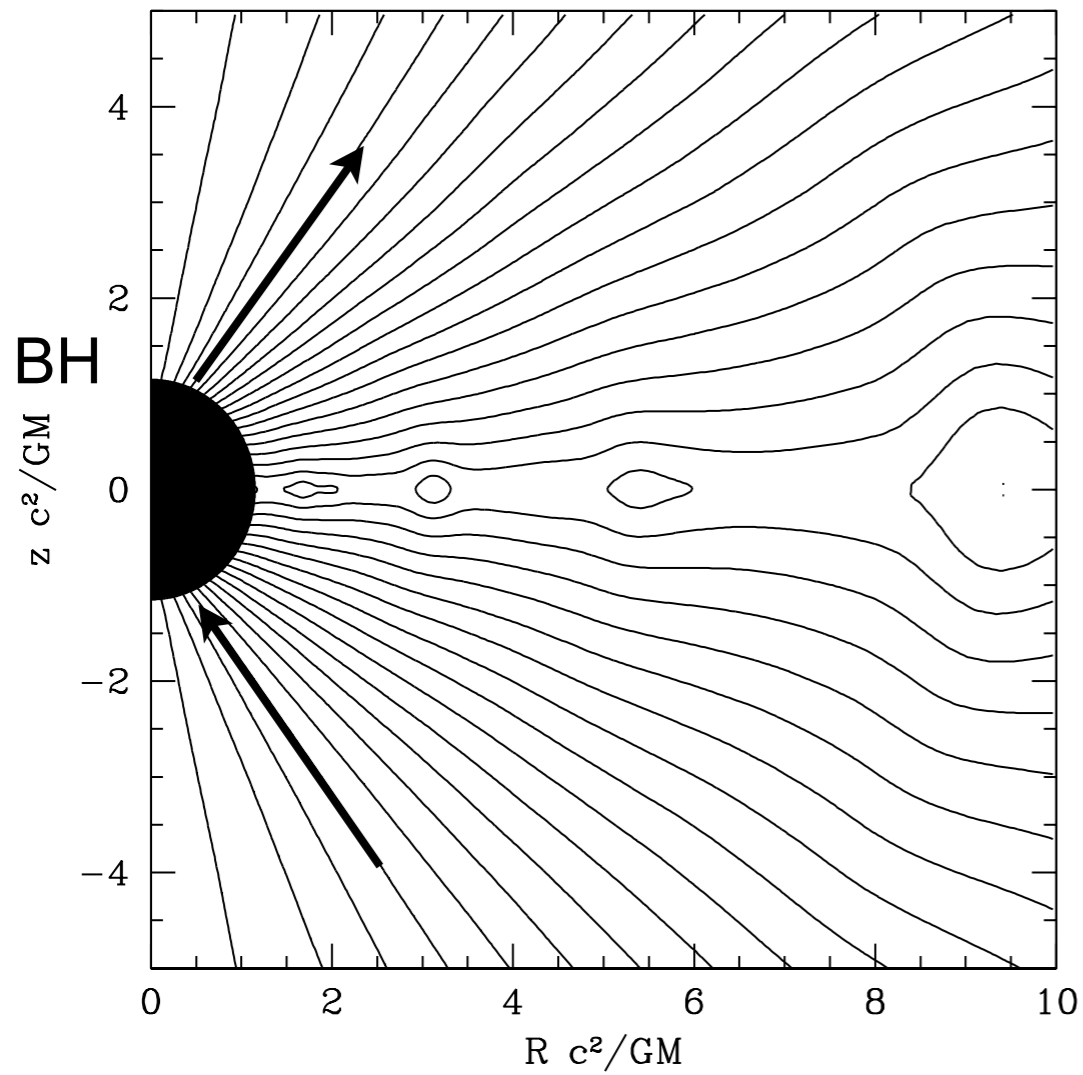
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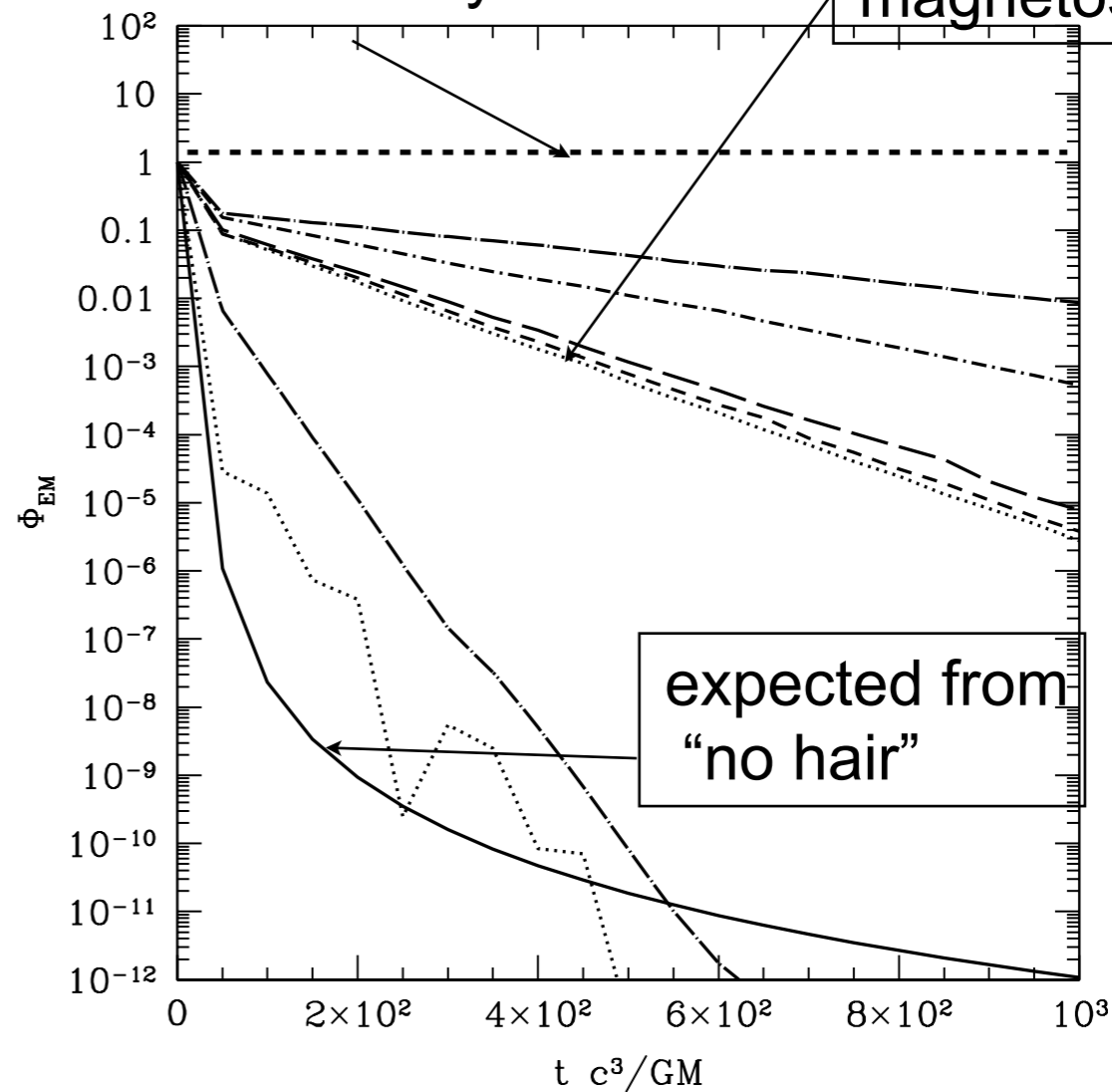
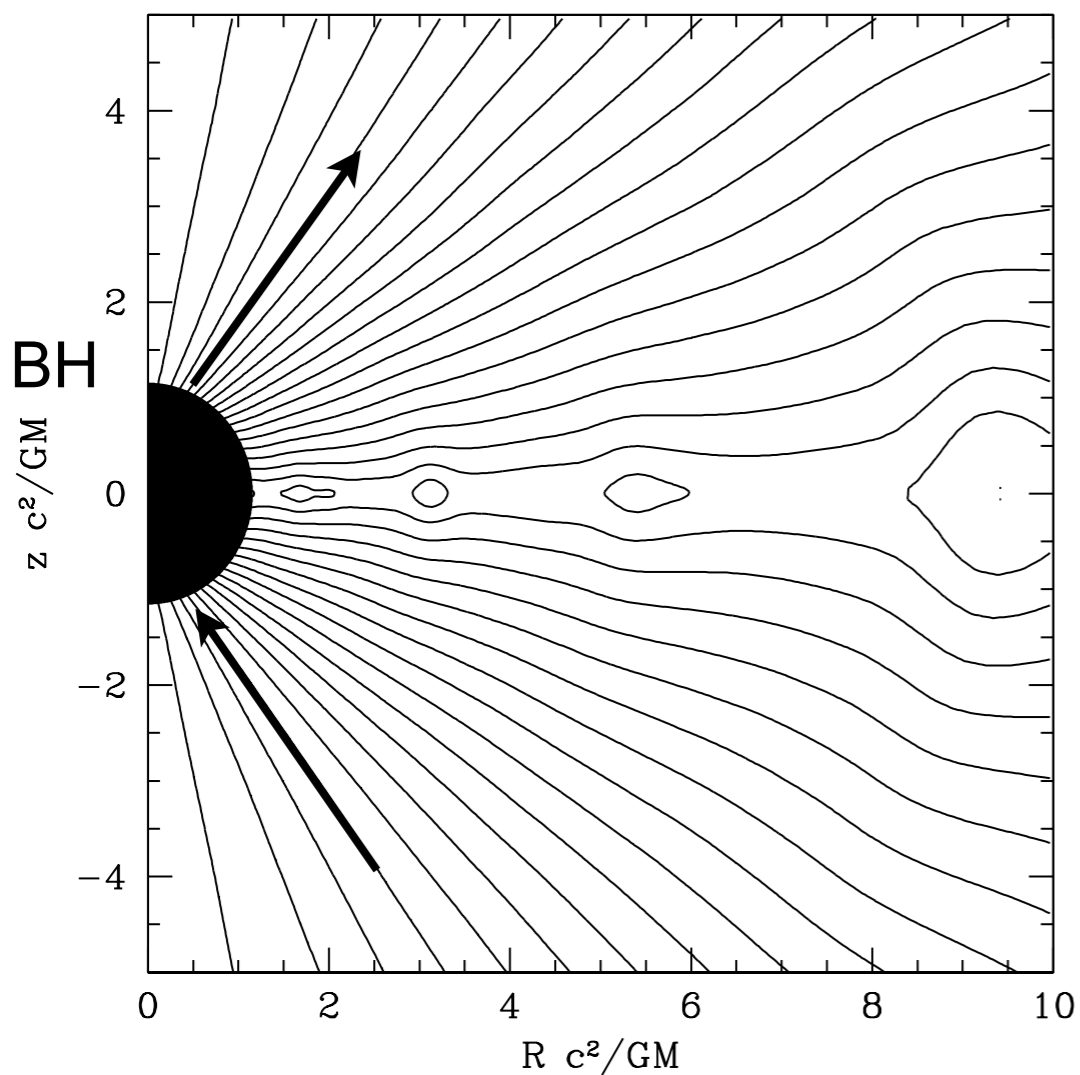


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- Split-monopole magnetosphere
- Slow balding

Expected for no numerical resistivity

With conducting magnetosphere

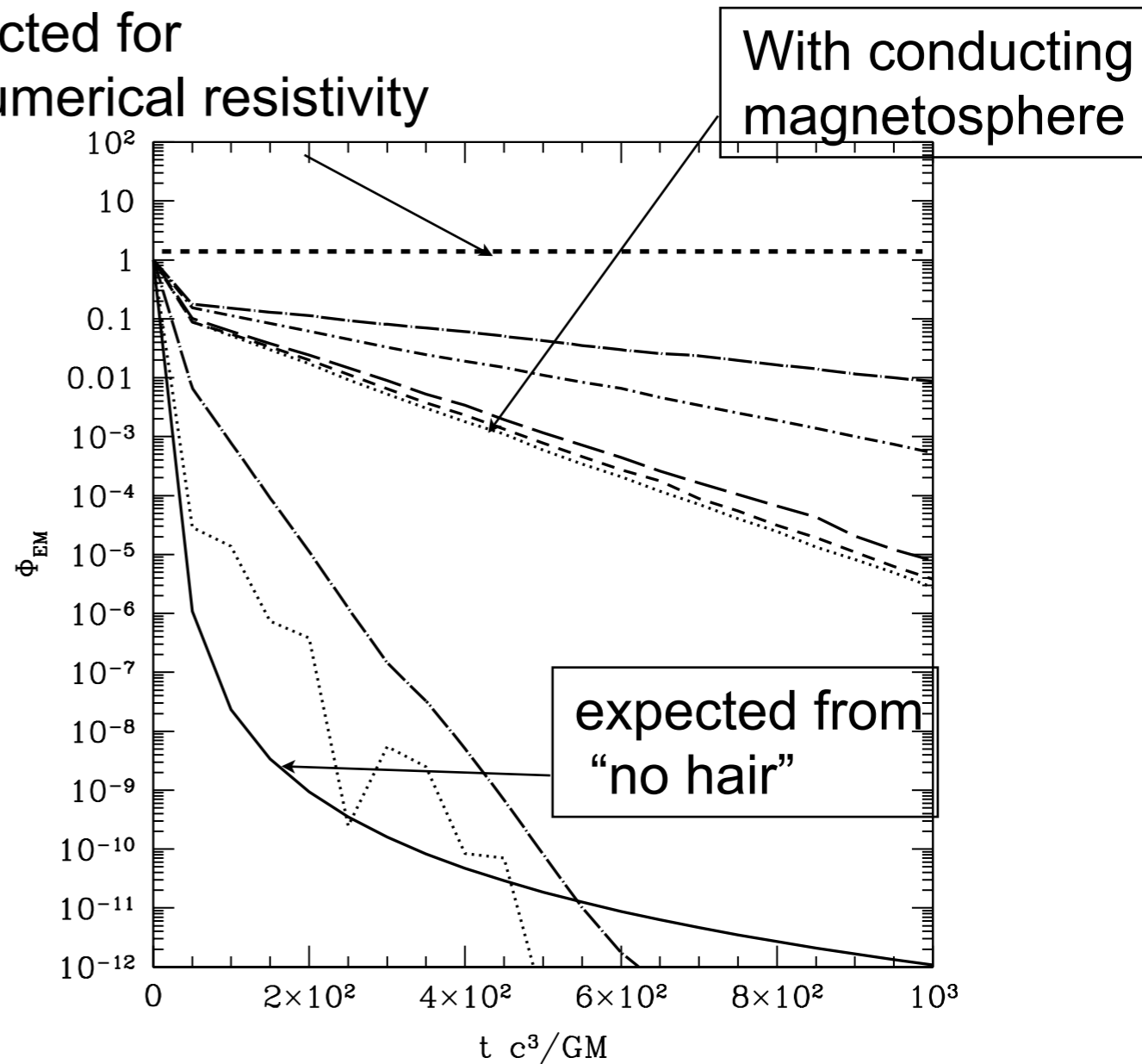
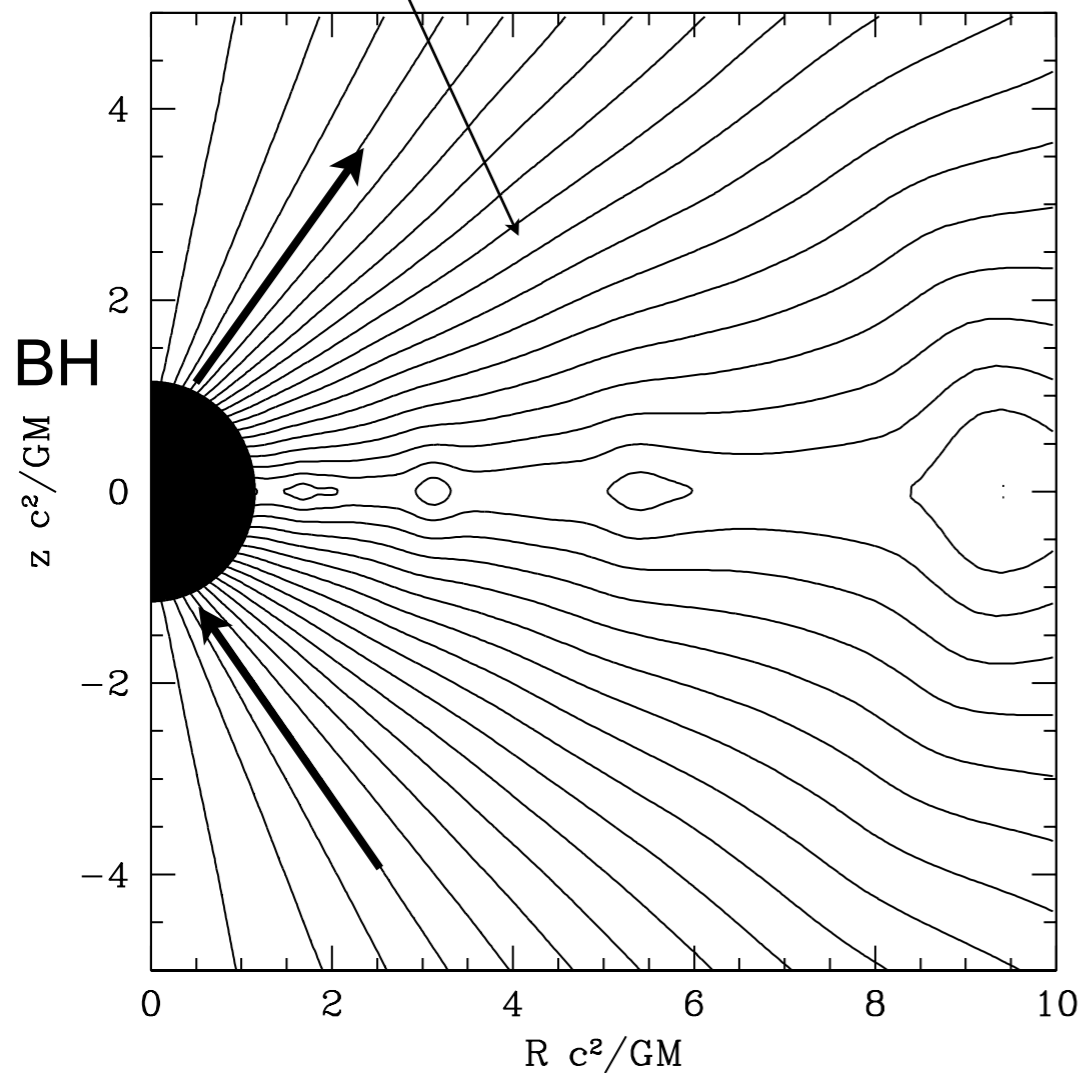


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Fields are NOT anchored in heavy crust

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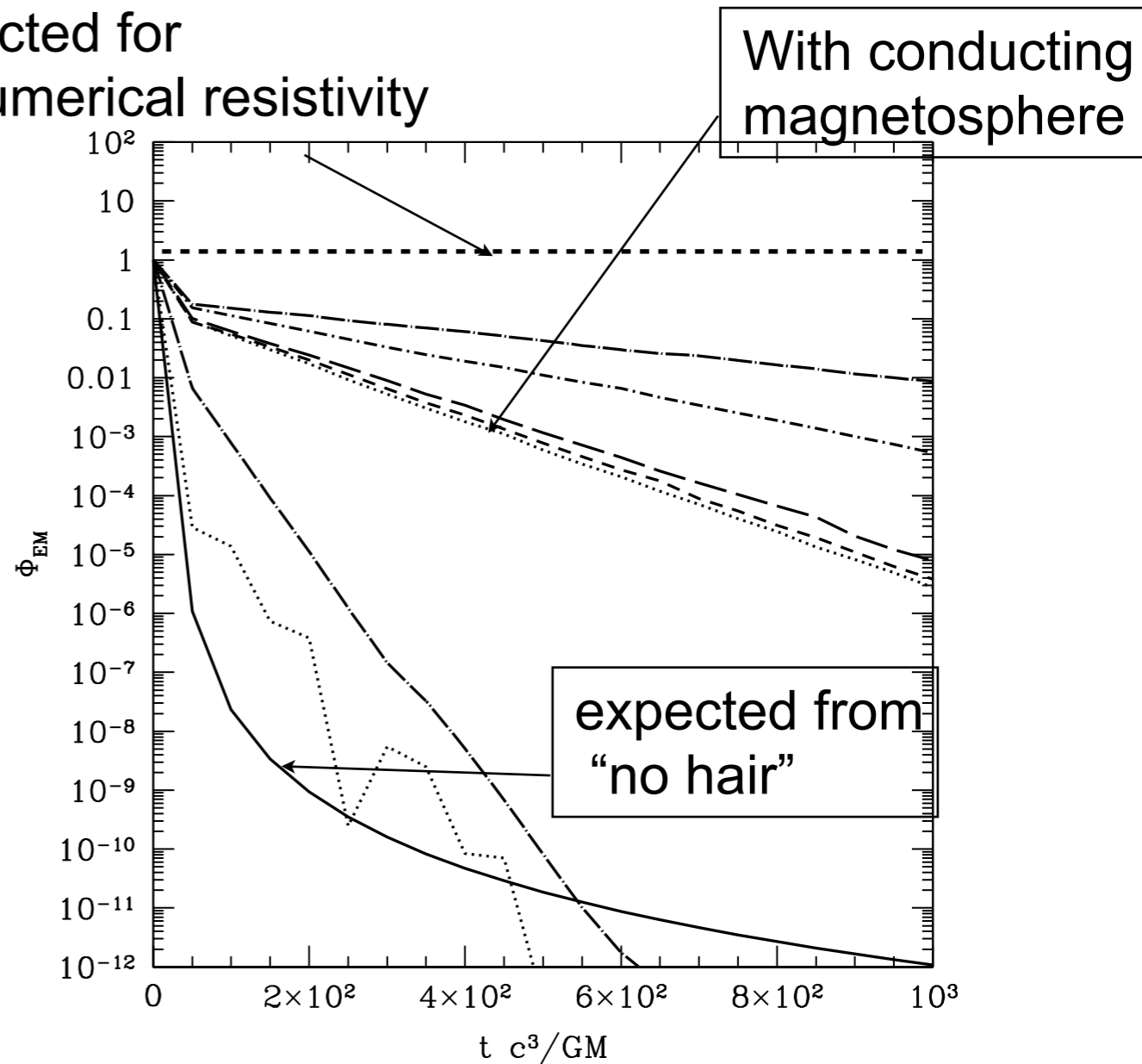
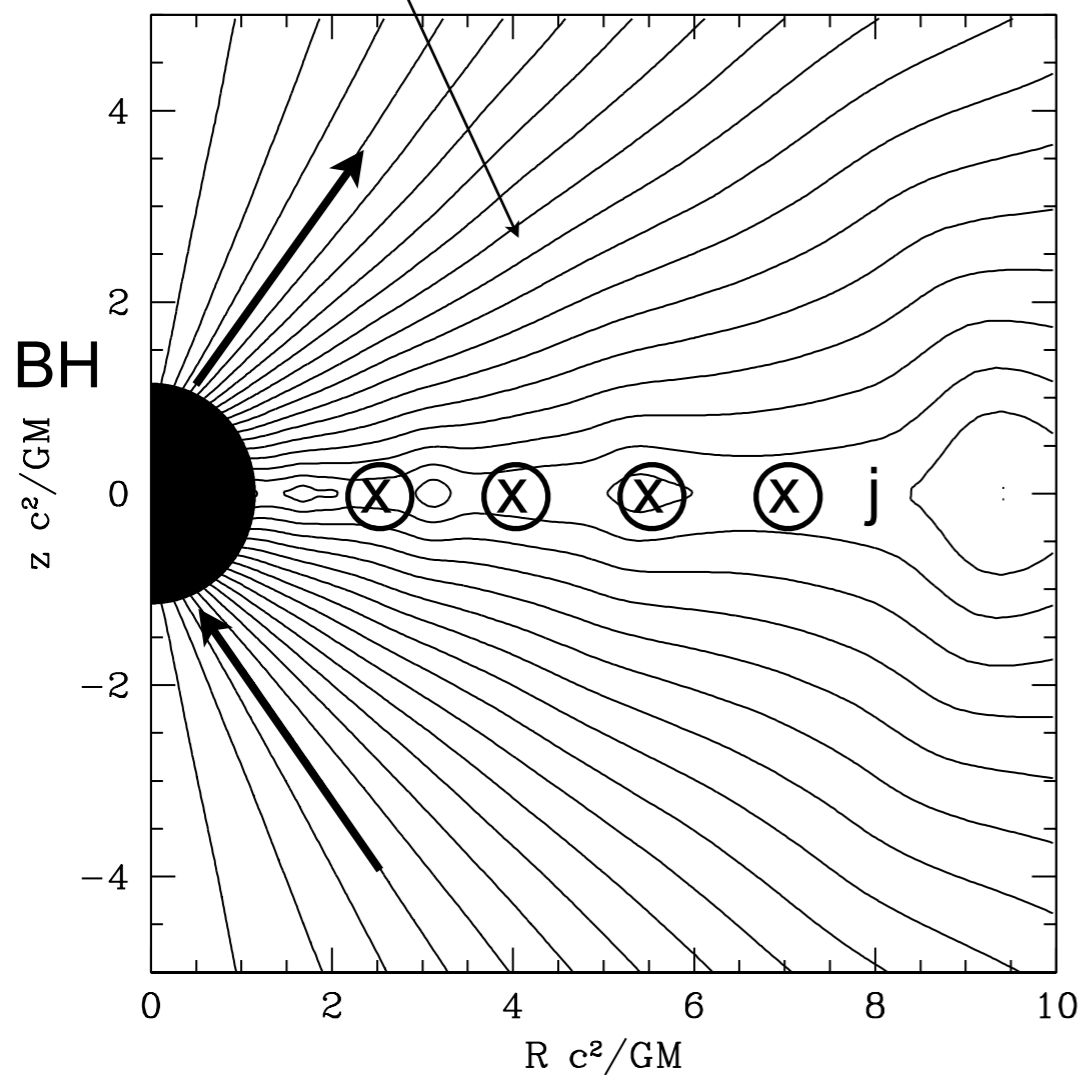


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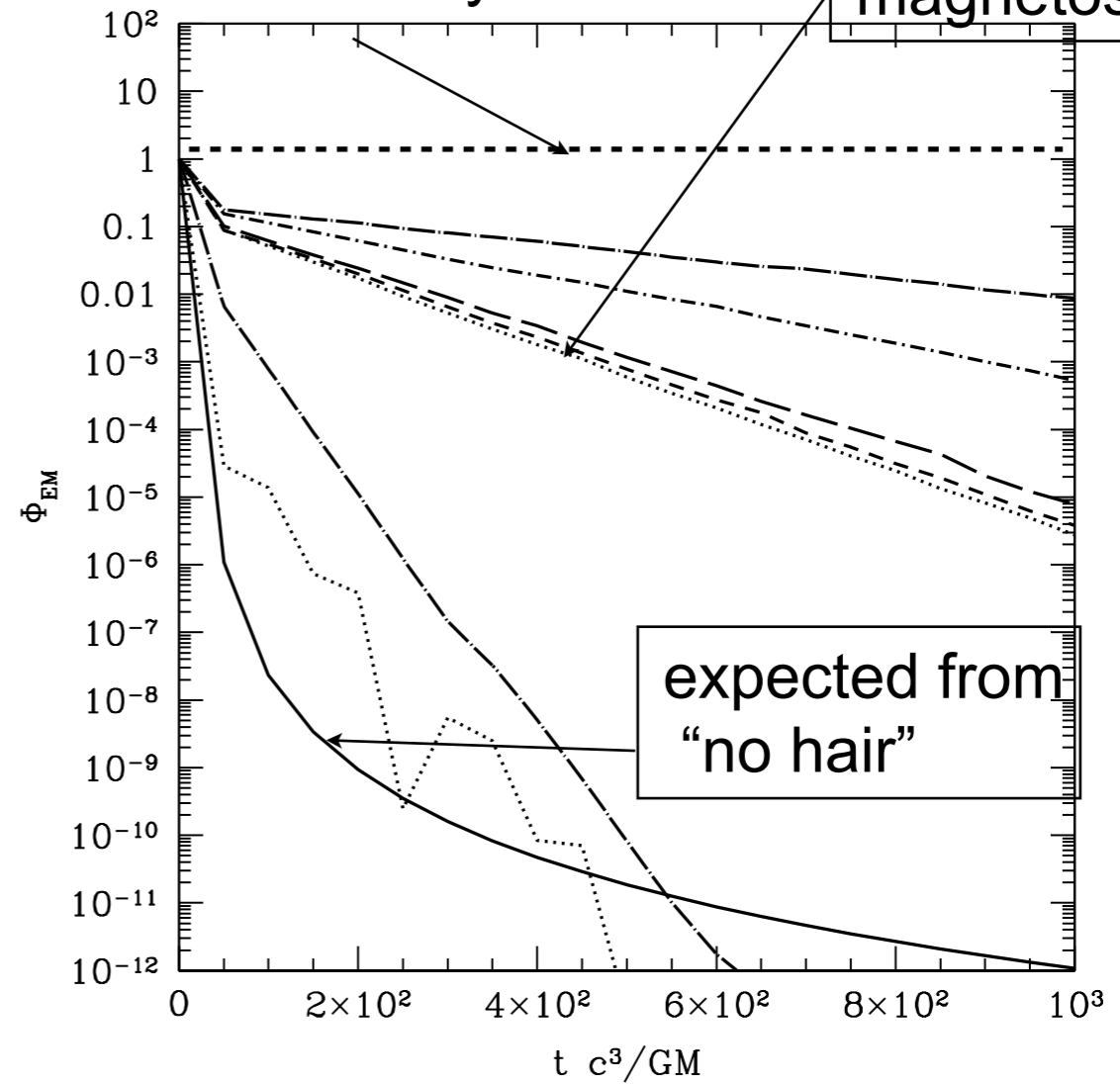
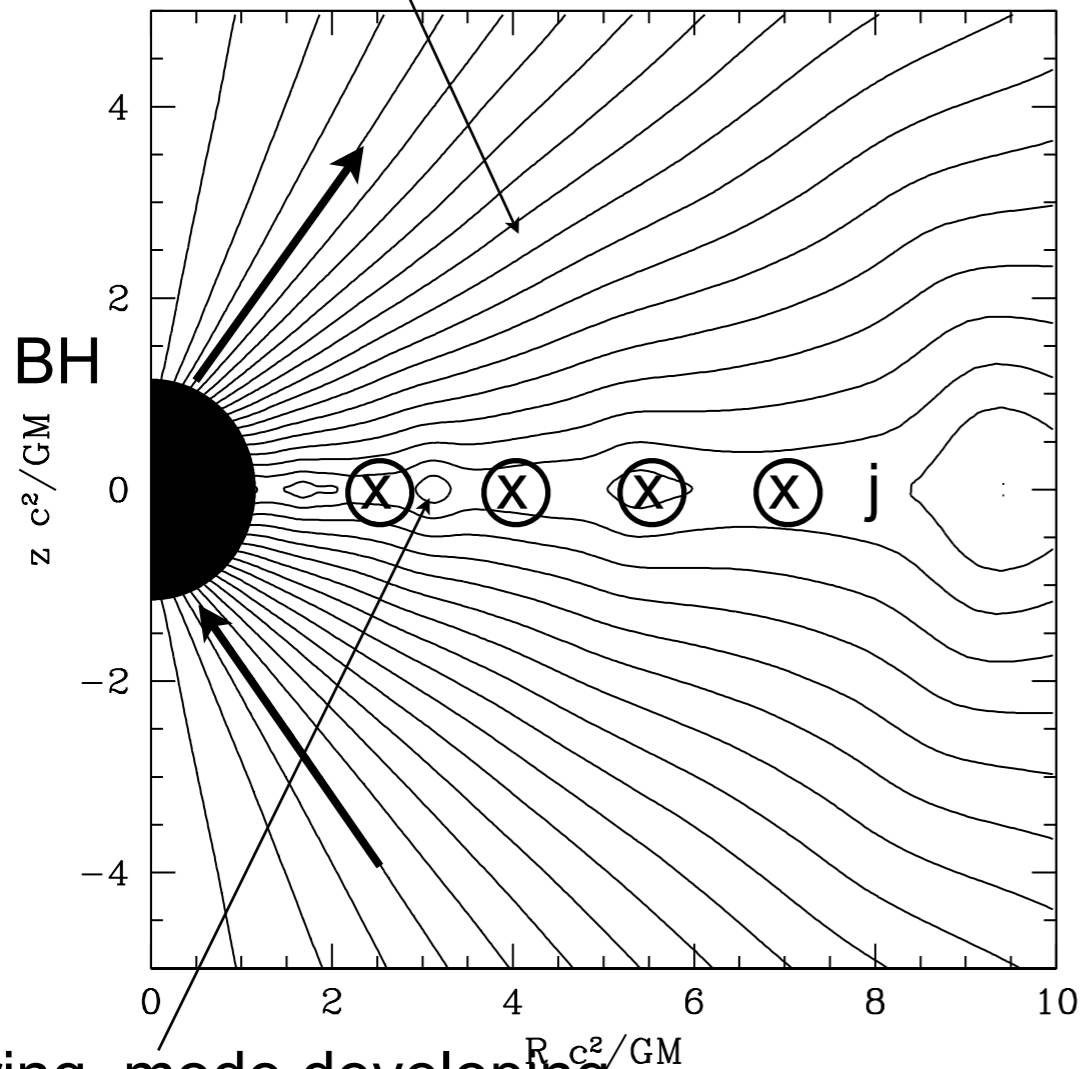
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Tearing mode developing

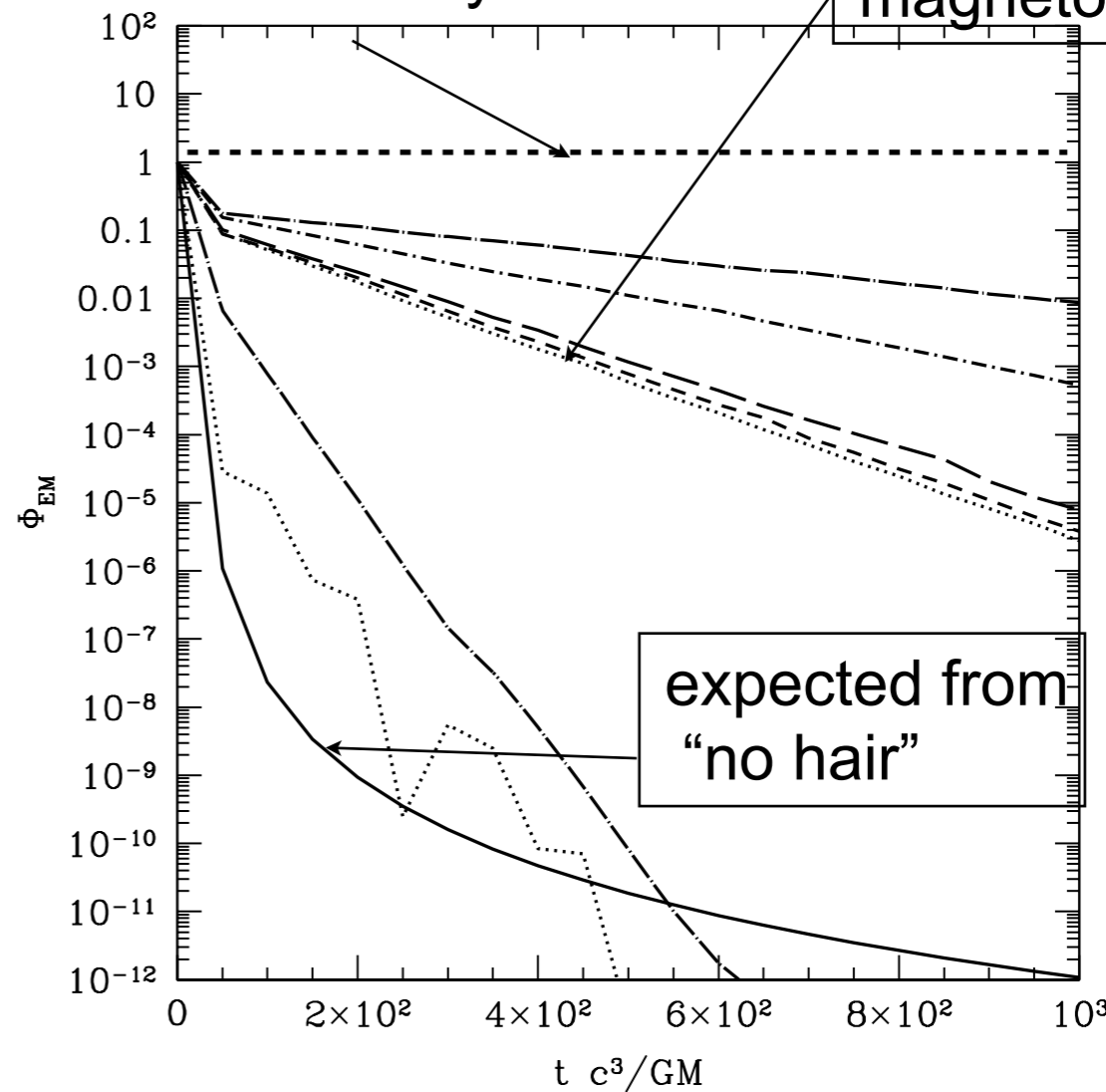
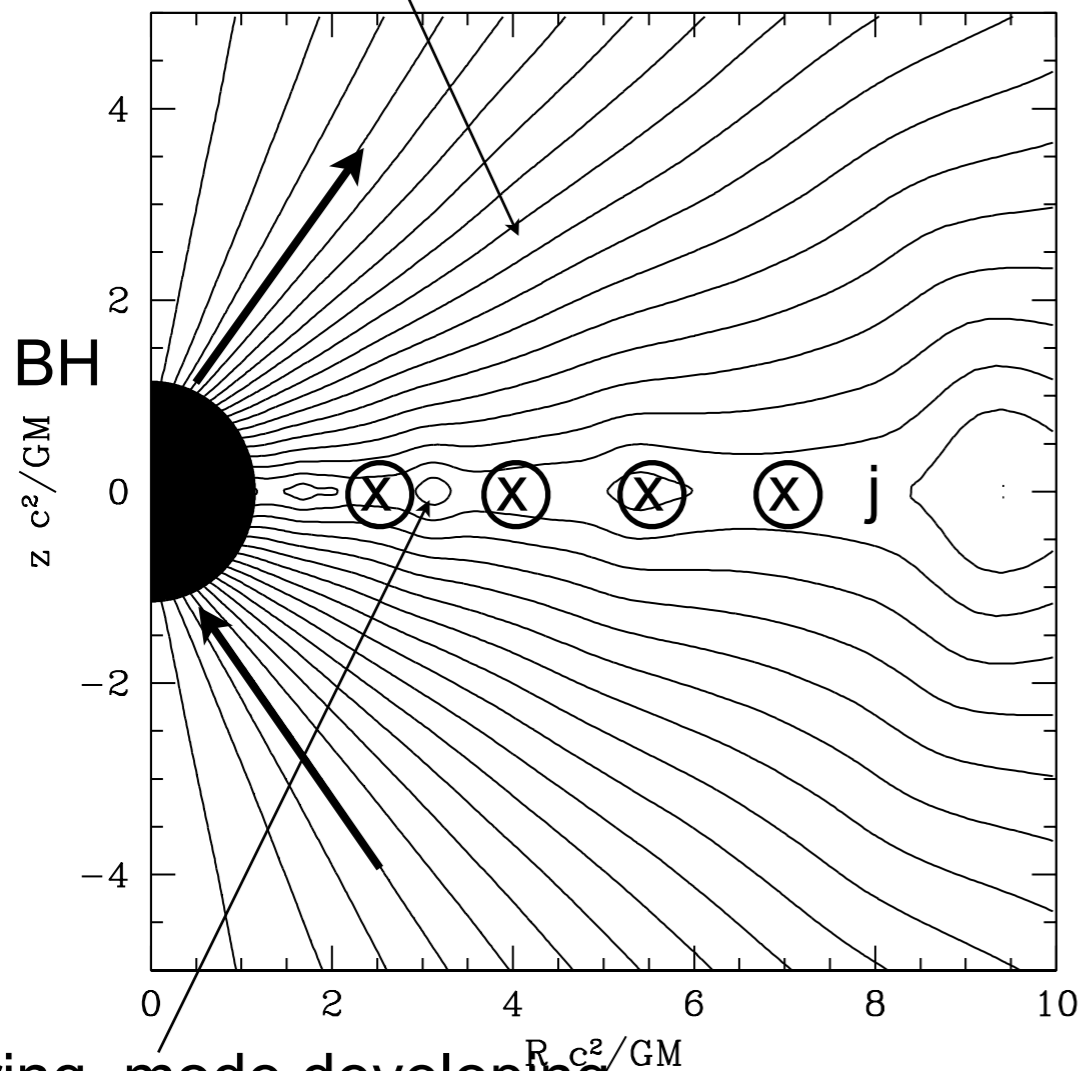
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Fields are contained by the equatorial current, just like in BZ, but this current is self-produced

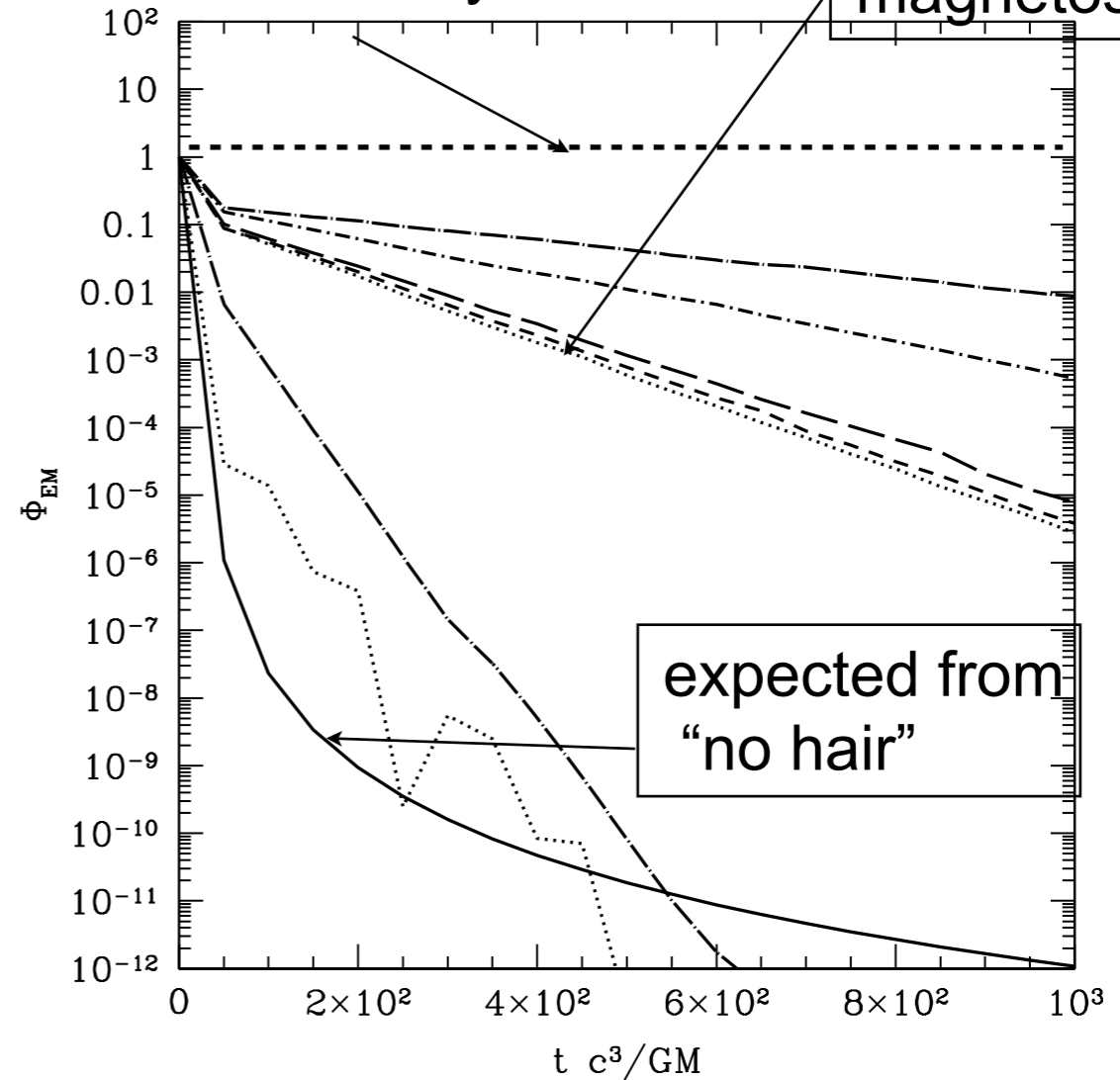
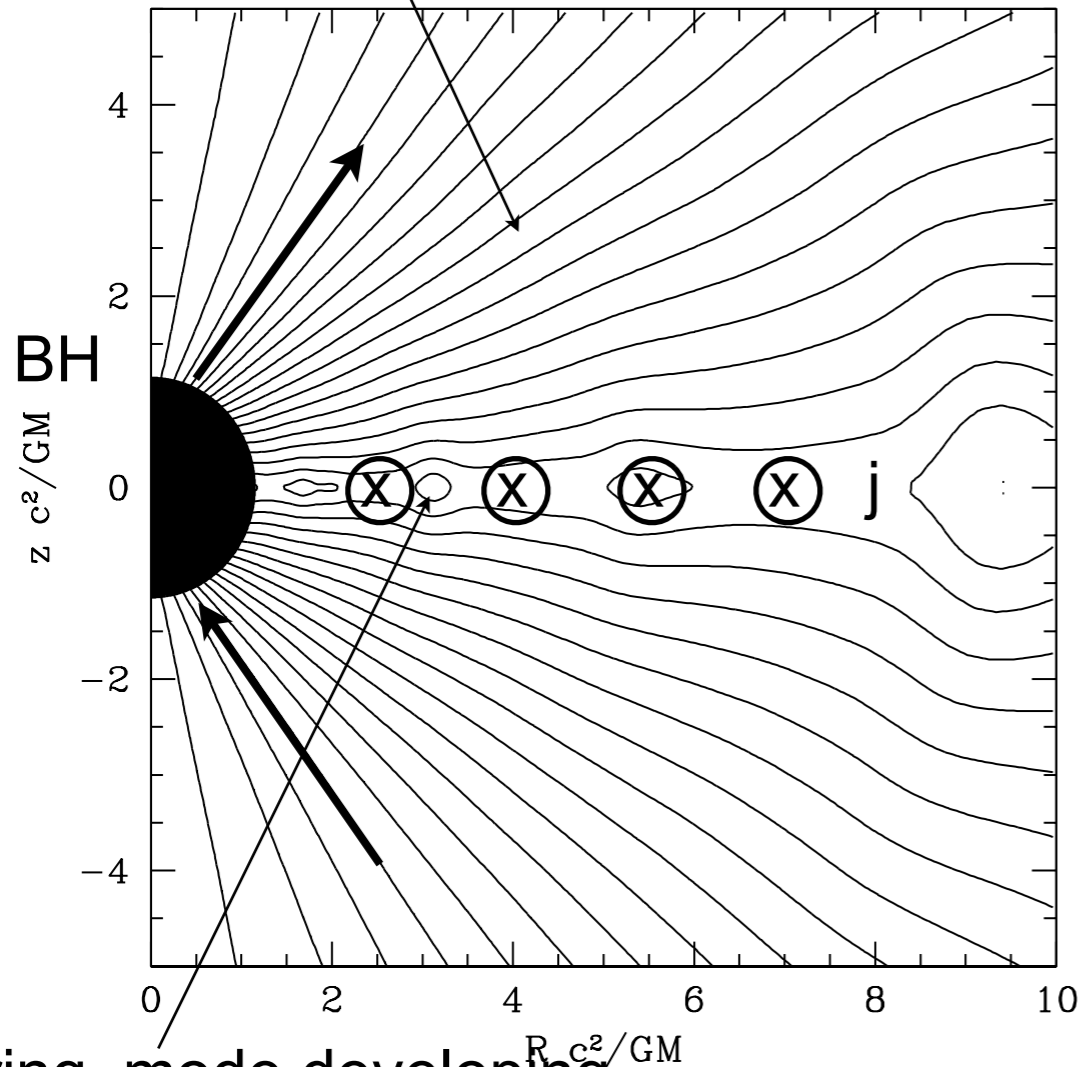
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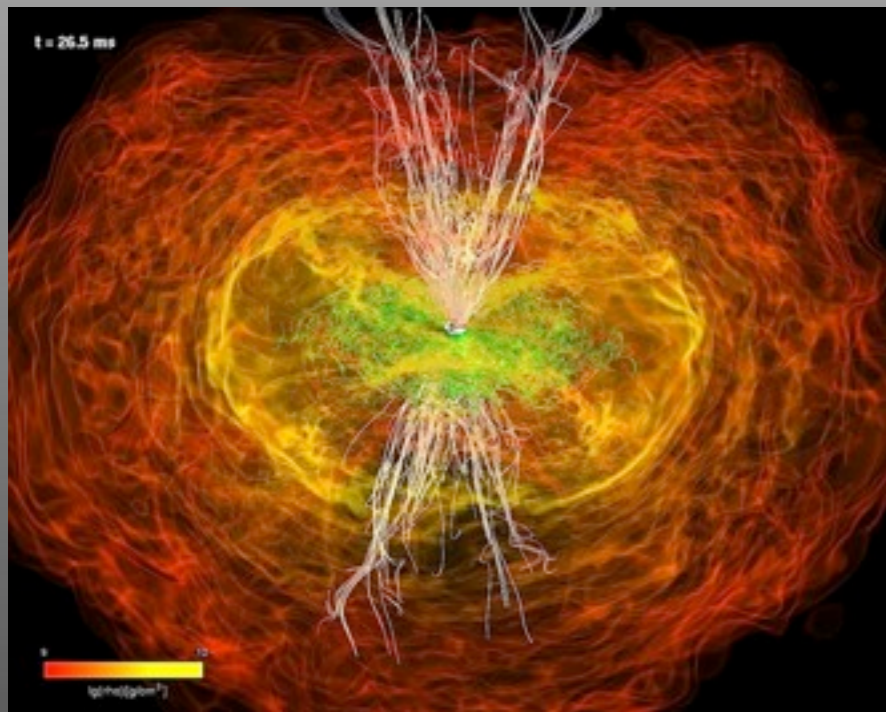
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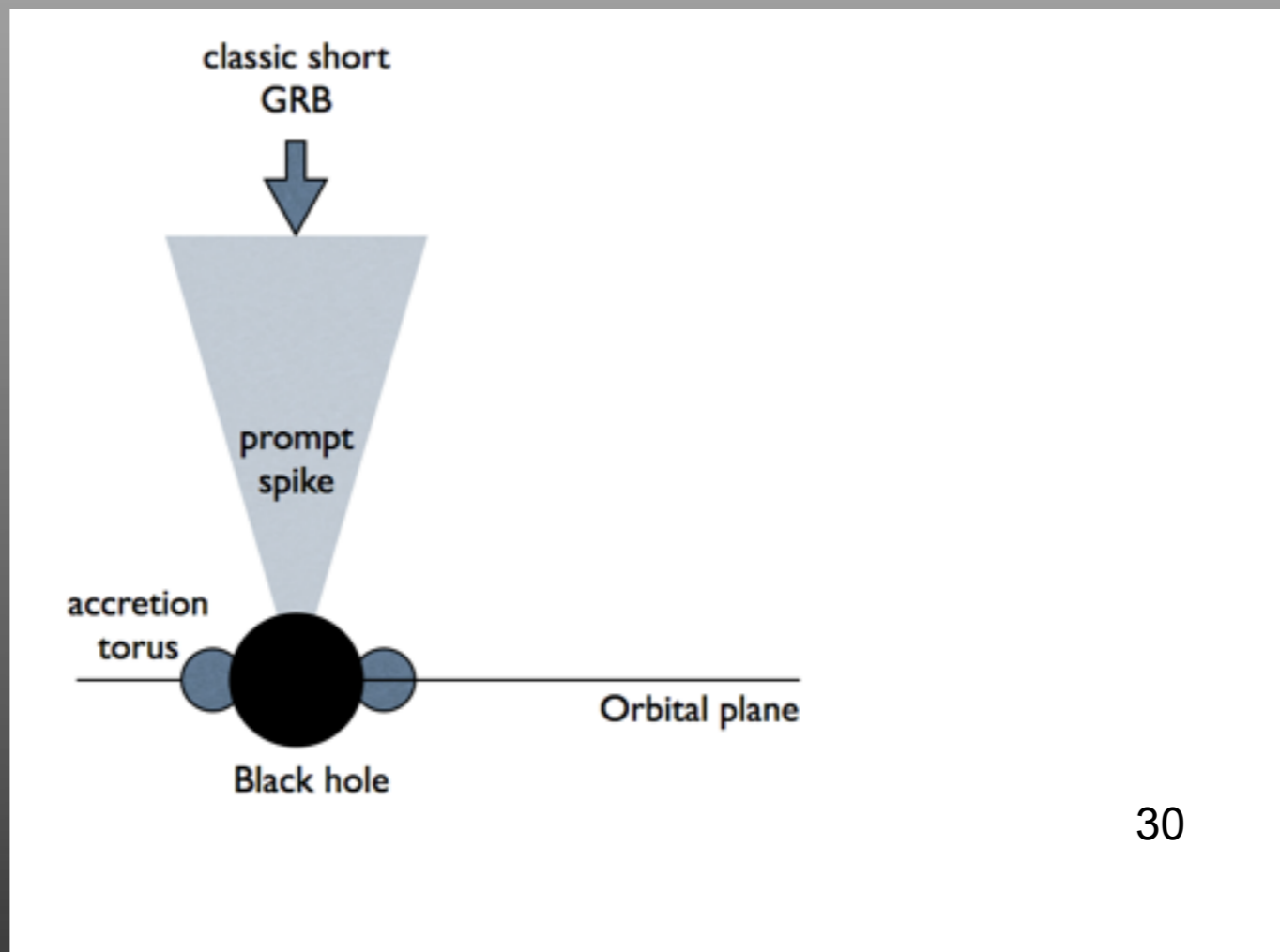
Biggest problem: hard to predict resistive time

The electromagnetic model of short GRBs

- NS-NS merger generates $B \sim 10^{15}$ G in the torus around BH (Rezzolla et al.)
- BH-torus launches a jet along the axis: prompt spike
- After ~ 100 msec torus collapse, isolated BH spins down electromagnetically, produces **equatorially-collimated** flow, $L \propto \sin^2 \theta$: prompt tail
- **Tail is more energetic**, but de-boosted for axial observer
- Very late re-brightening of the remnant

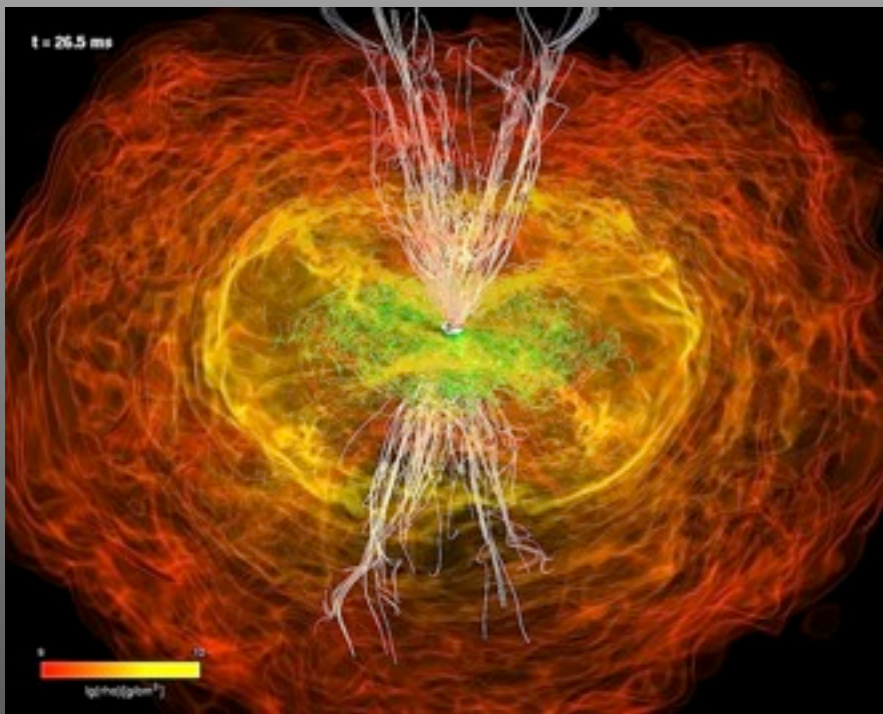


Rezzolla et al

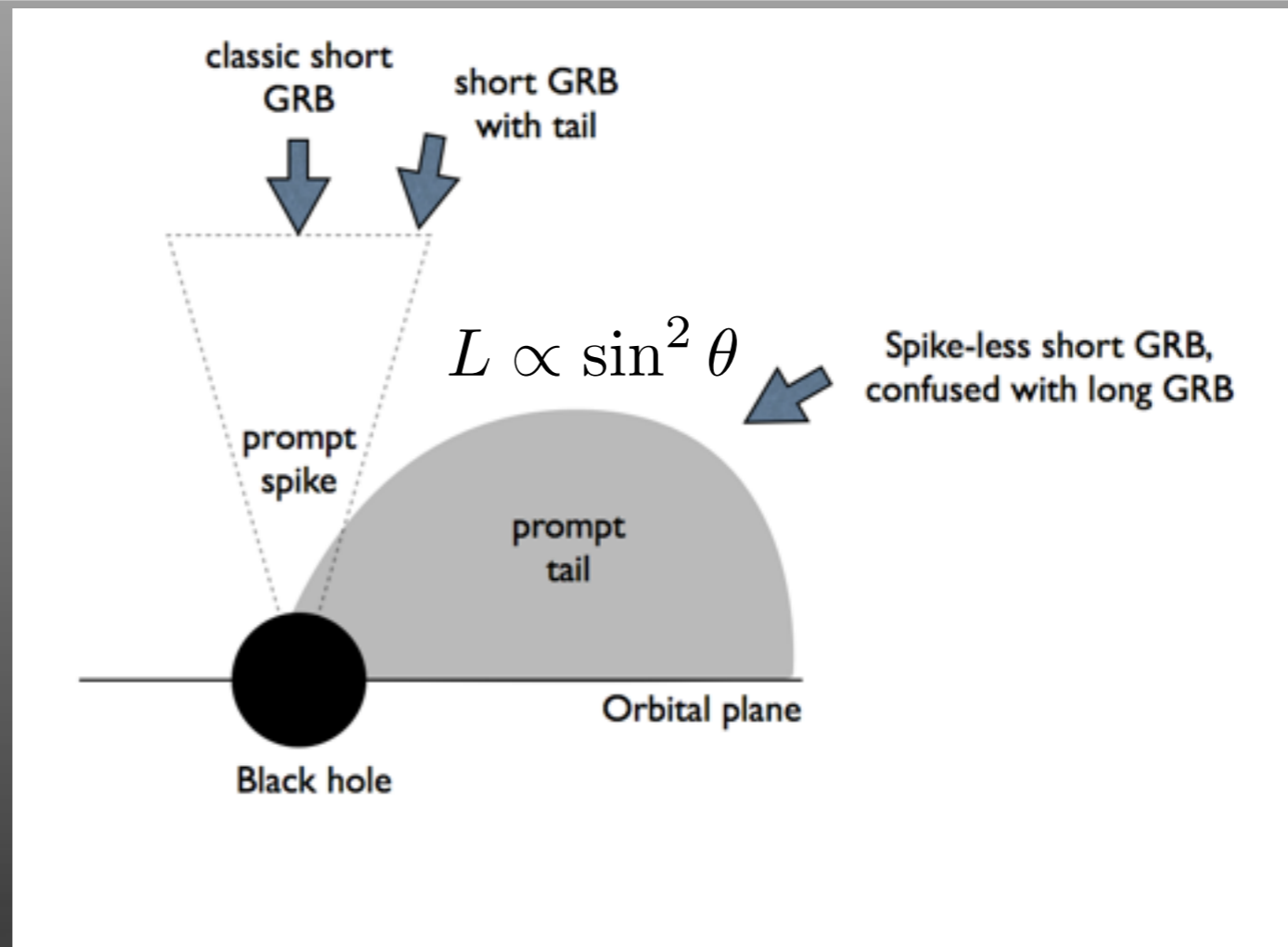


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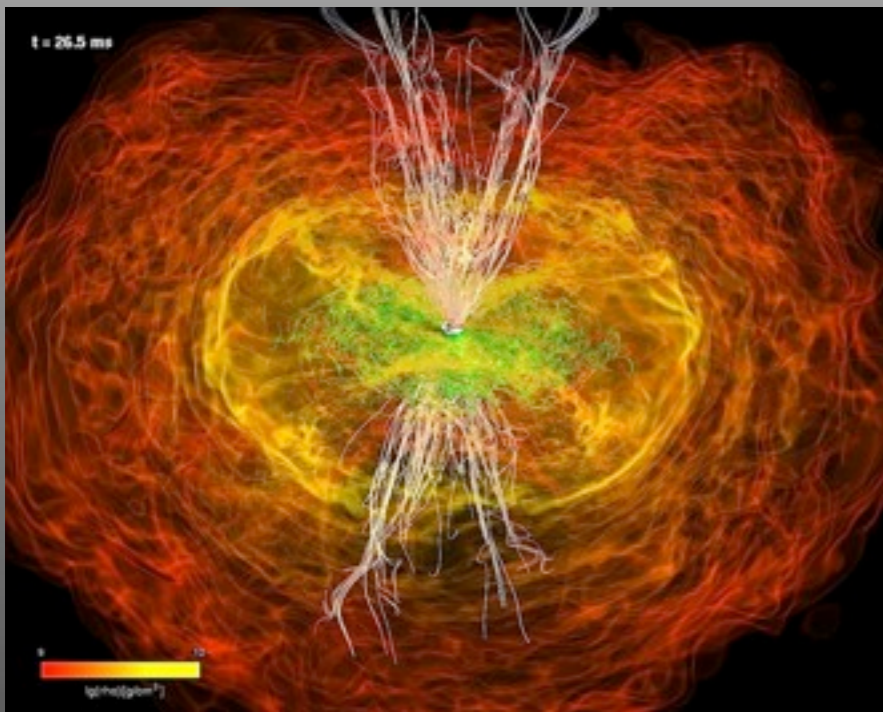


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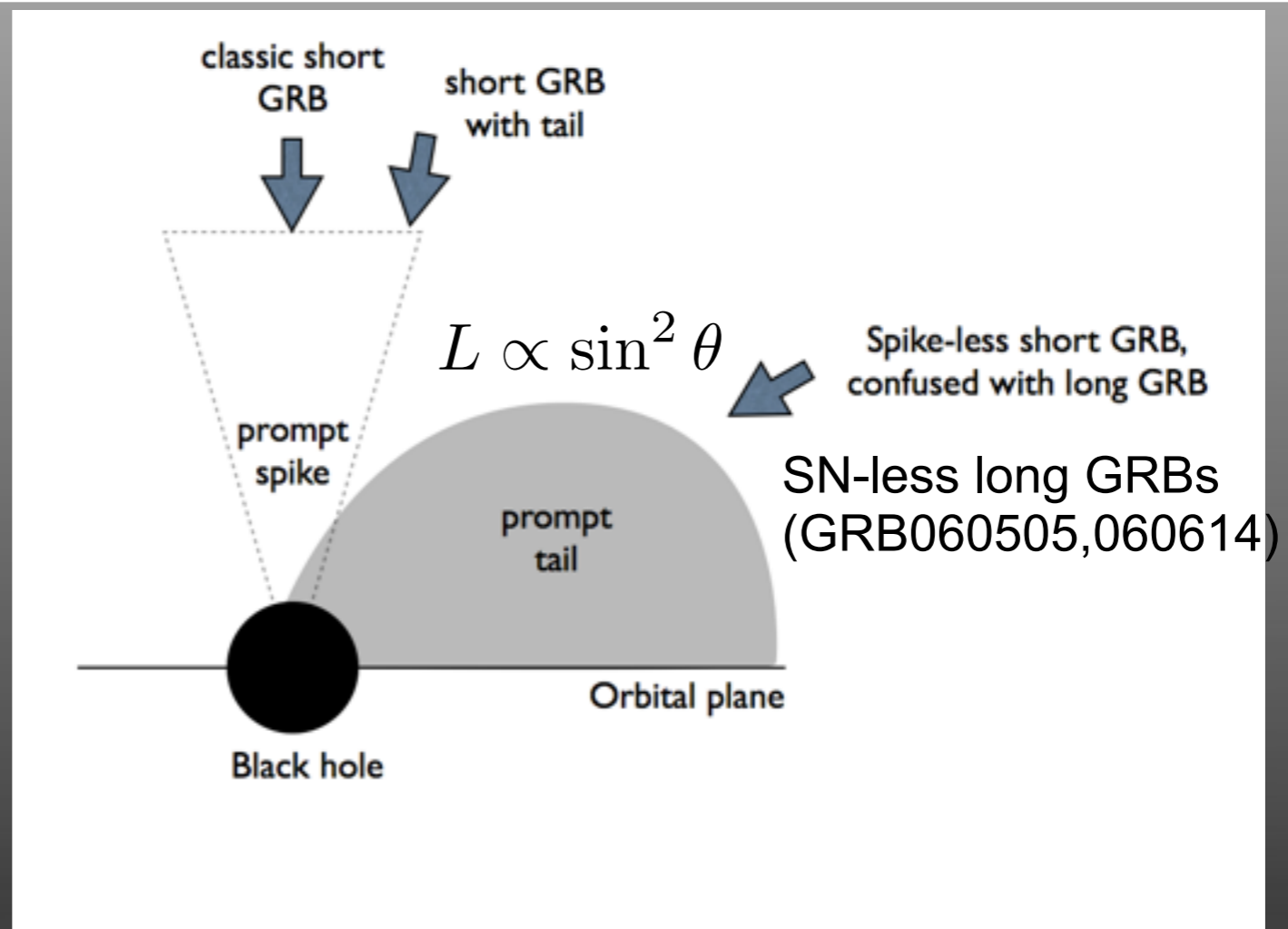


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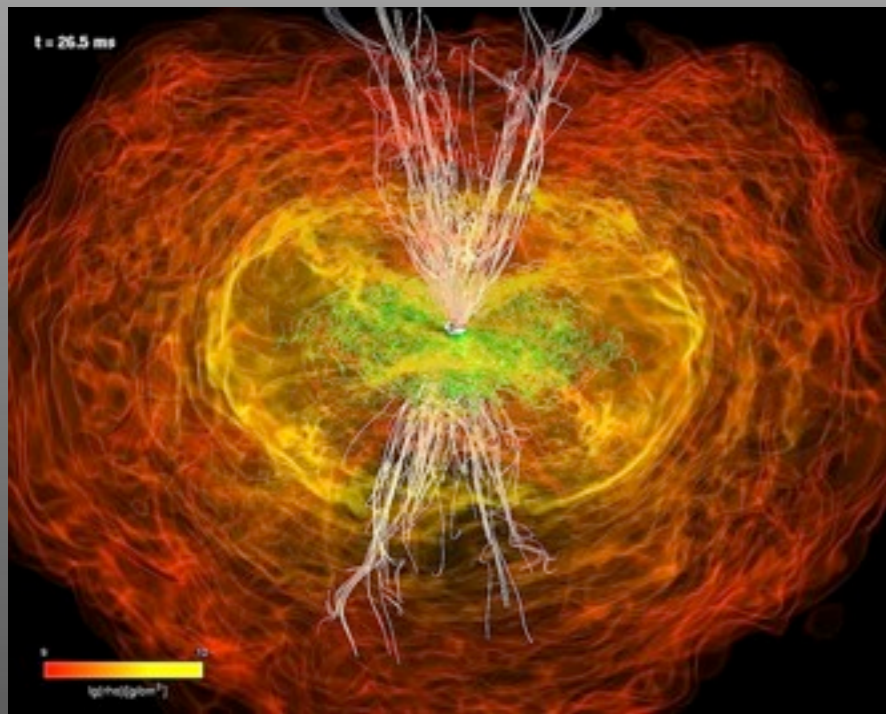


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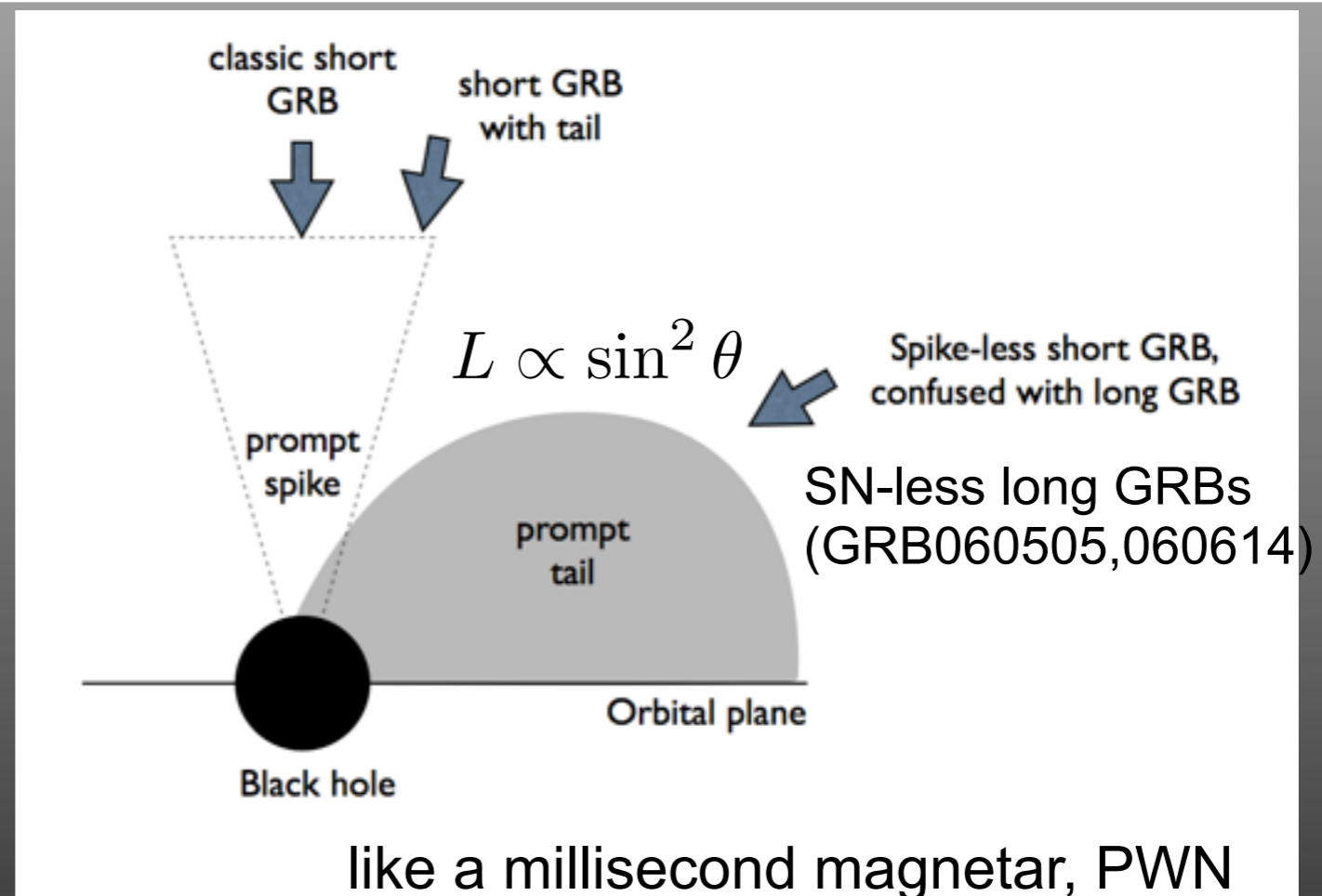


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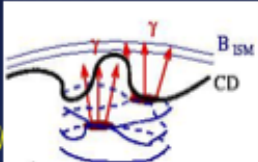
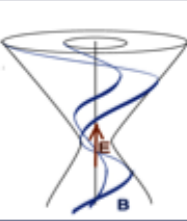


Millisecond magnetar as GRB central engine

- Millisecond magnetars is a promising central source
 - can produce clean (after few seconds), highly relativistic outflows
 - can operate on long time scales without external feeding
- Magnetic dissipation/particle acceleration *a la* Crab flares can be important (dominant?) in GRBs
 - Bursty, short time scales from large radii
 - fast efficient acceleration
 - non-thermal tail
- Newly born BHs may work as millisecond magnetars - prompt tails in short GRBs

Particle acceleration

- I. *Instability of polar currents*
(~Z-pinch) – narrow emission
- II. *Instability of shell currents* –
 4π emission
 - Electro-magnetic turbulence particle acceleration
 - Breakdown of MHD (Lytikov&Blackman00)
 - Inertial acceleration (Usov&Smolsky00)
 - Synchrotron emission



Lytikov 2003

