

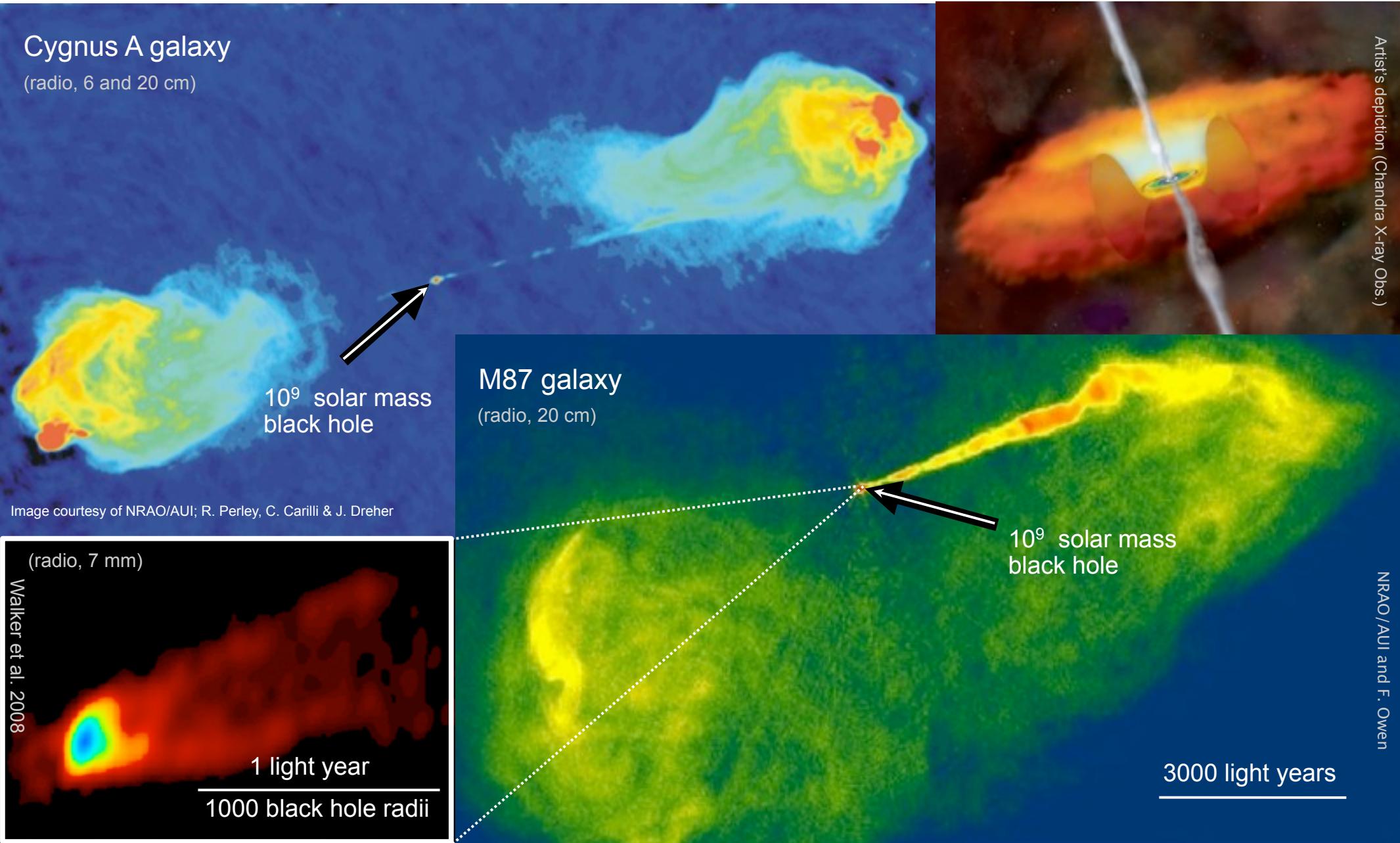
General Relativistic MHD Simulations of Relativistic Jets in AGN

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Jon McKinney, Stanford-Maryland

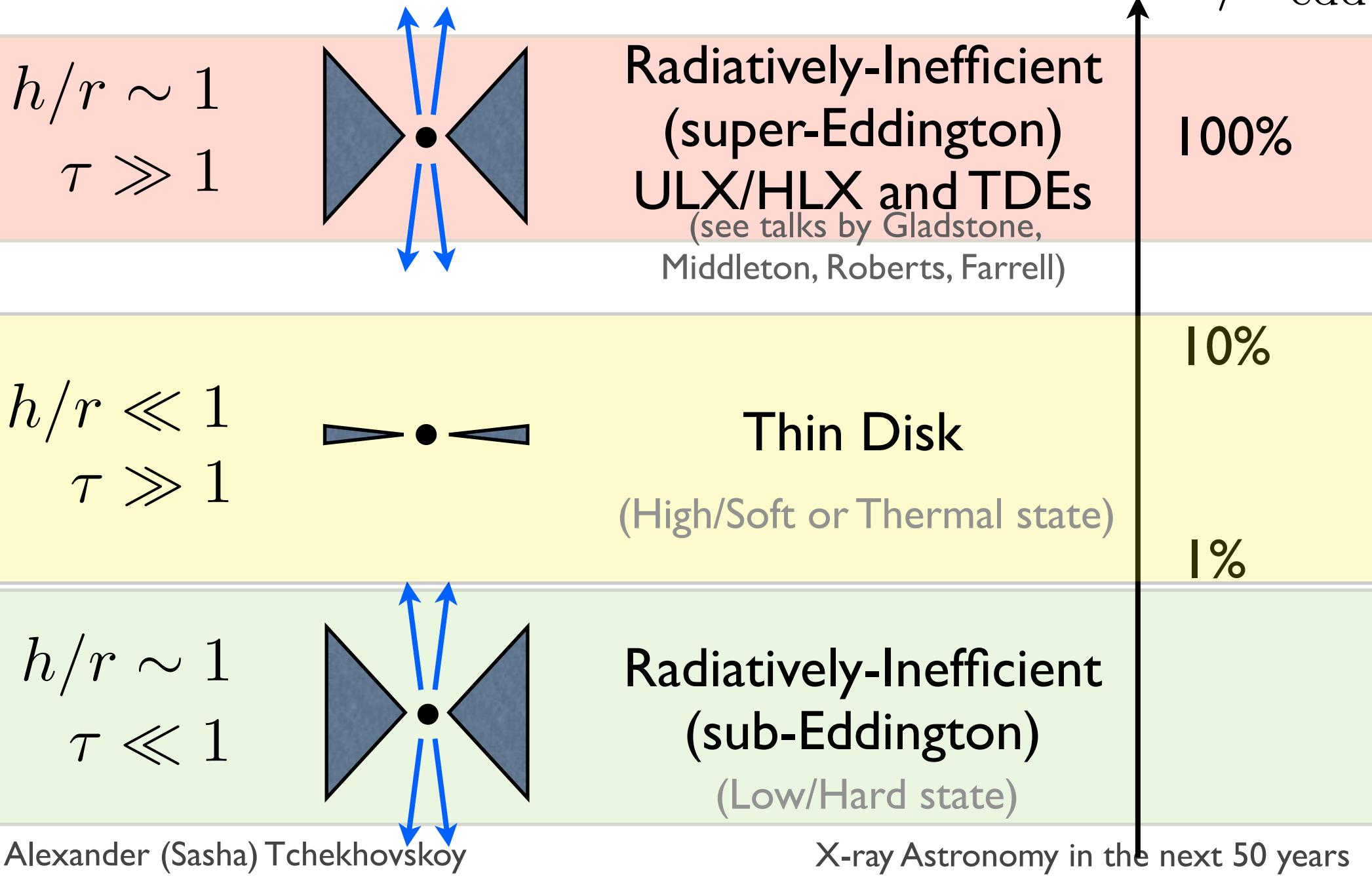
Jets: Beautiful & Challenging



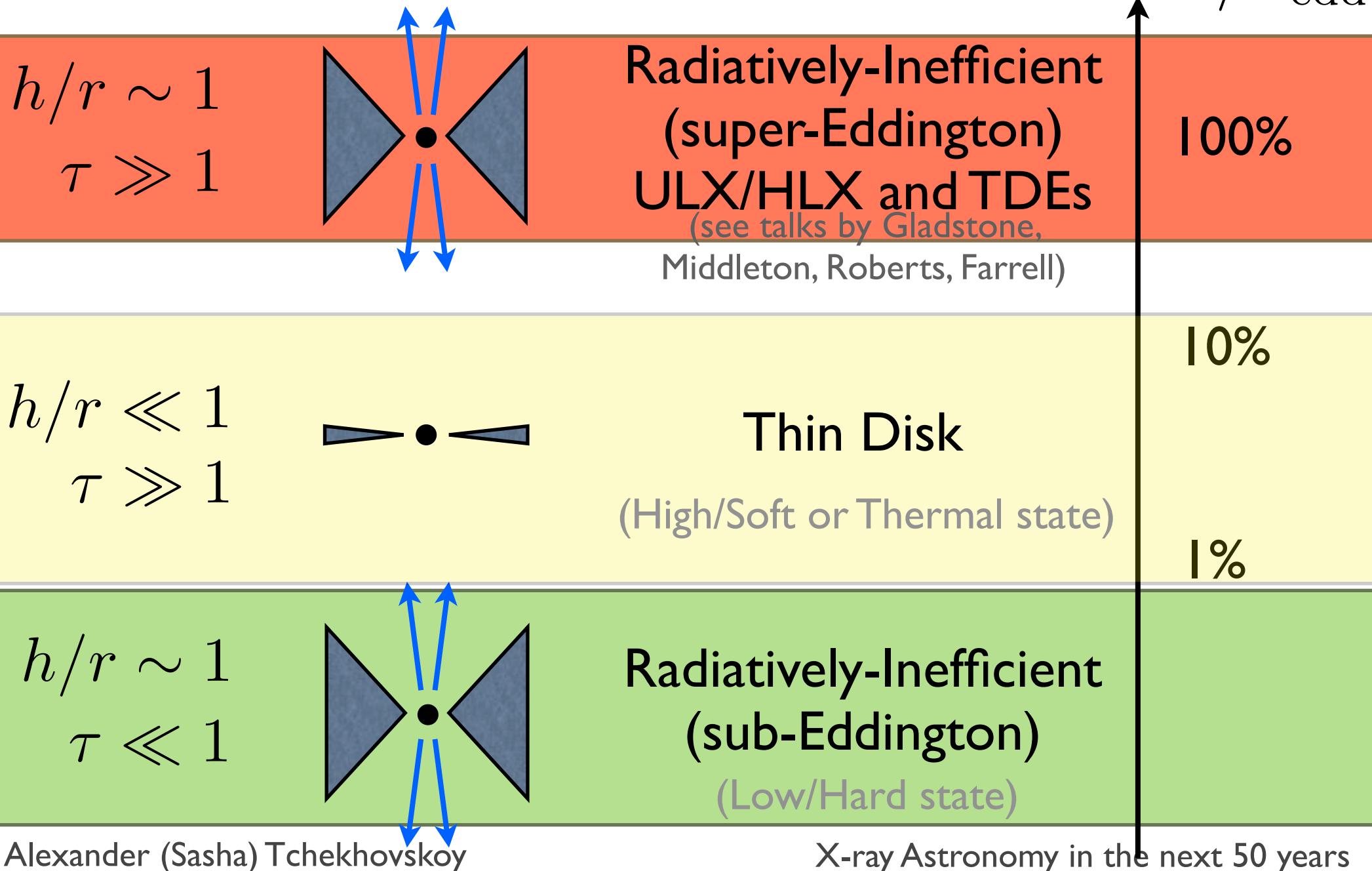
Fundamental Questions

- What sets the maximum power of jets?
- Are jets powered by black holes or inner regions of accretion disks?
- How does jet power depend on BH spin?
- Does accretion always spin up BHs to high spins?

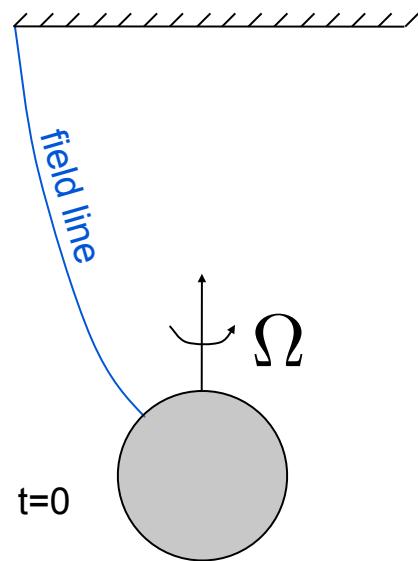
When Do We Observe Jets?



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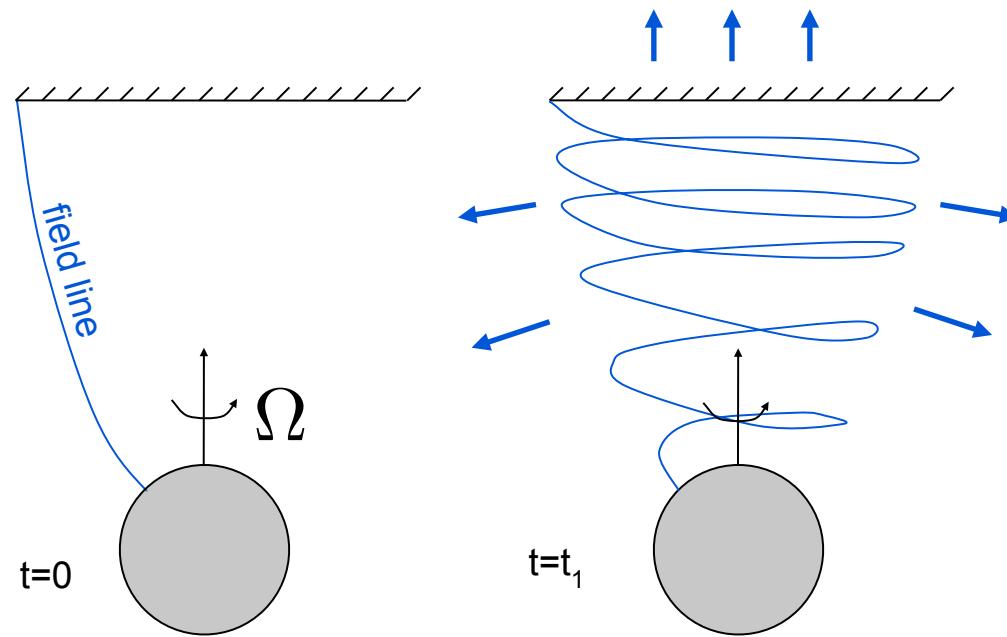


How Do Jets Work?

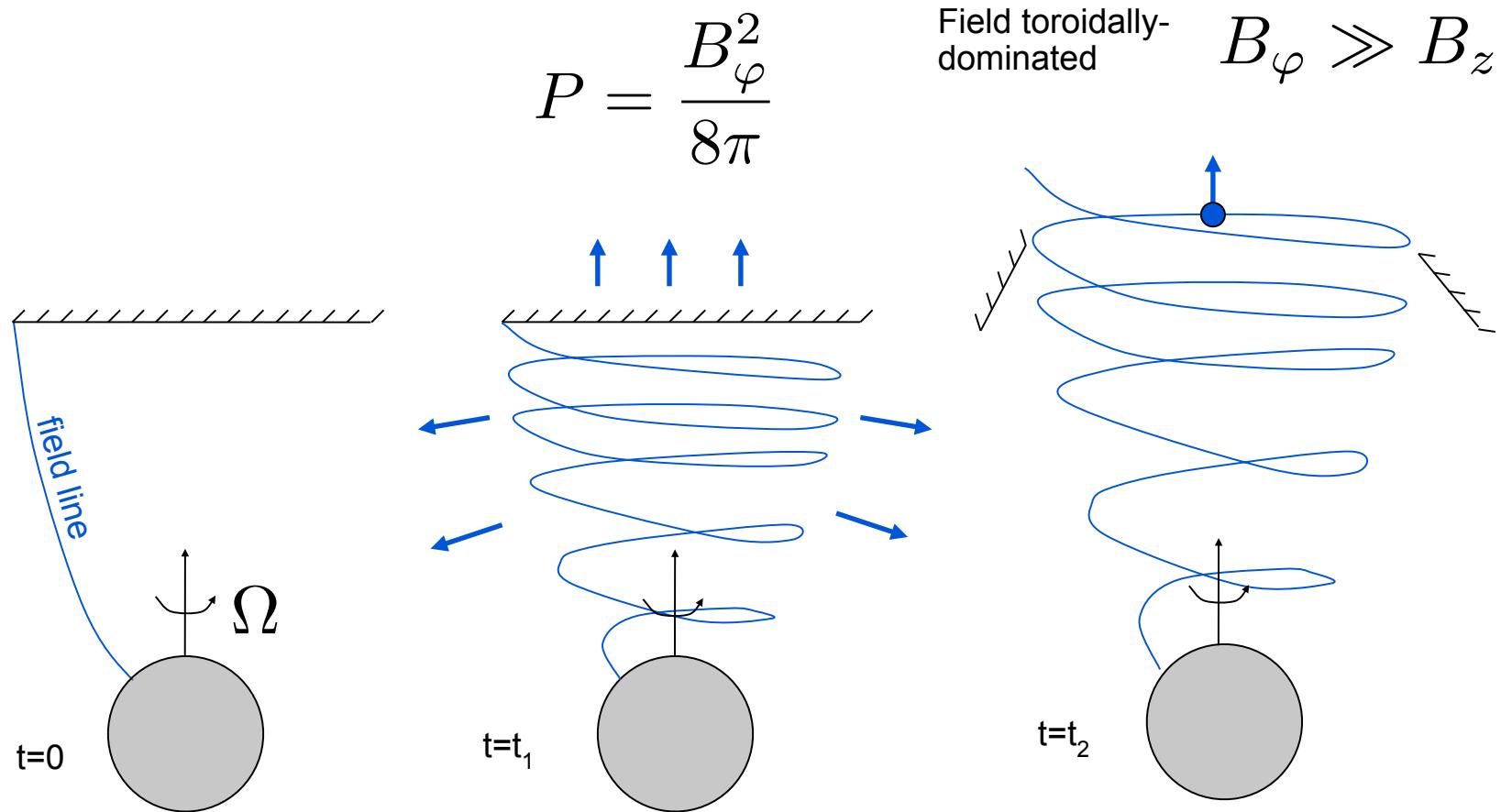


How Do Jets Work?

$$P = \frac{B_\varphi^2}{8\pi}$$



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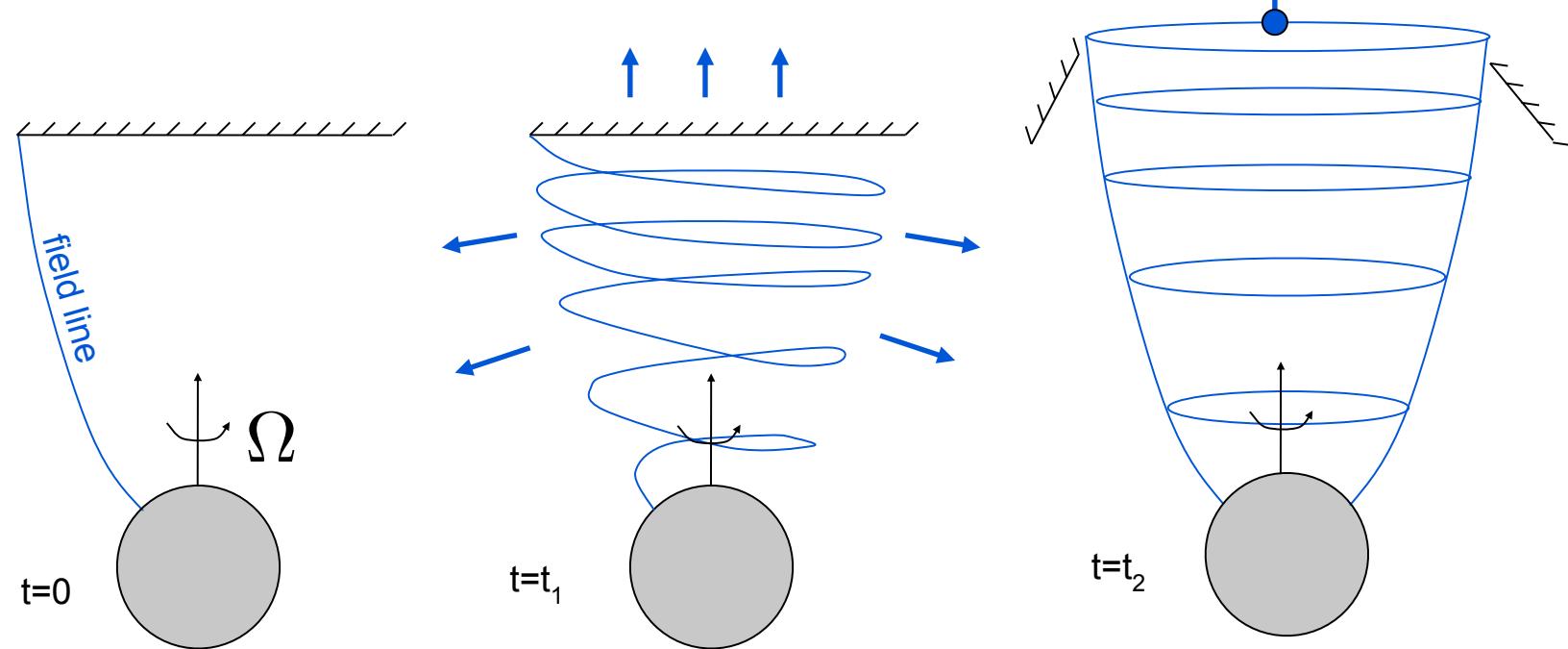


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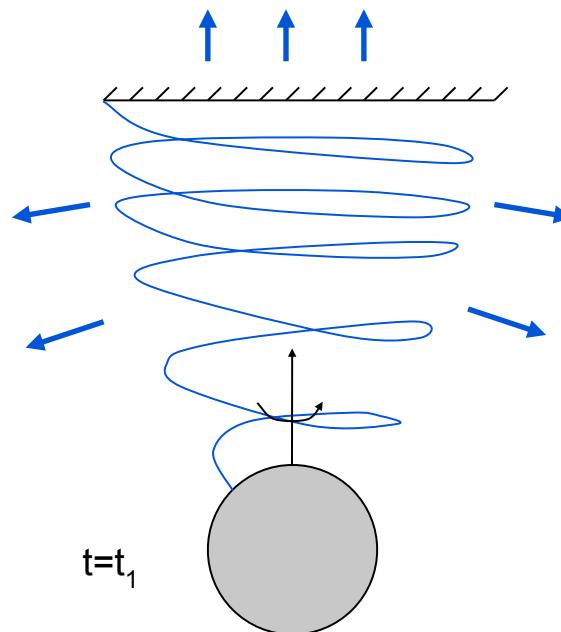
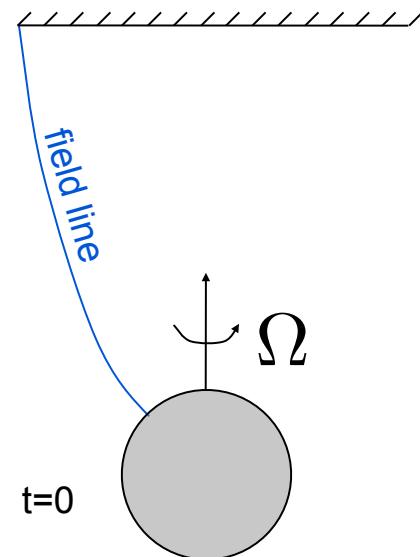
Field toroidally-dominated

$$B_\varphi \gg B_z$$

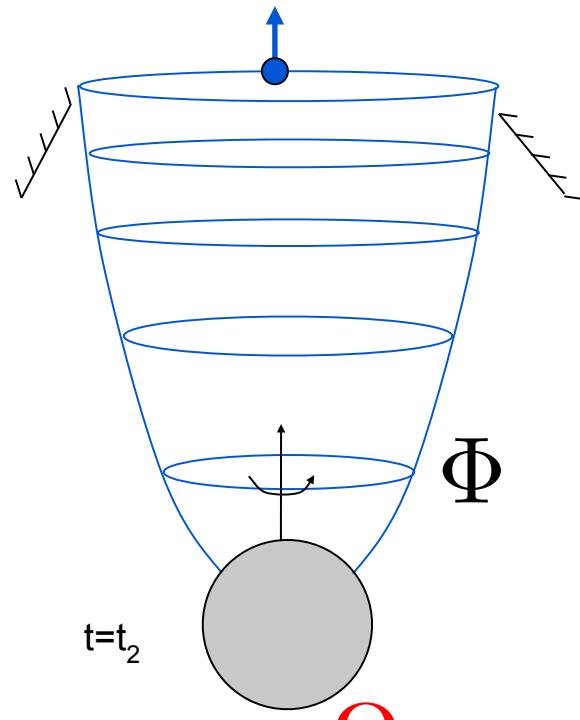


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Field toroidally-dominated $B_\varphi \gg B_z$



$$\Omega_H = \frac{ac}{2r_H}$$

What Sets BH Power?

- We understand well how BH power depends on Φ and Ω_H :

$$P_{\text{BZ}} = \frac{k}{c} \Phi^2 \Omega_H^2$$

(Blandford & Znajek 1977,
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$$P_{\text{BZ}} = \frac{k}{c} \Phi^2 \Omega_H^2 \times f(\Omega_H)$$

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$$1 + 1.4(\Omega_H r_g/c)^2 - 9.2(\Omega_H r_g/c)^4$$

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- But, what sets value of the proportionality factor,

$$\phi = \frac{\Phi}{\sqrt{\dot{M} r_g^2 c}},$$

(Gammie 2002,
Komissarov & Barkov 2009,
Penna et al. 2010)

and BH power efficiency,

$$\eta_{\text{BZ}} = \frac{P_{\text{BZ}}}{\dot{M} c^2} = k \phi^2 \left(\frac{\Omega_H r_g}{c} \right)^2 \times f(\Omega_H)$$

What Sets the Flux?

- Gravity limits BH field strength:

$$\frac{B^2}{8\pi} \lesssim \frac{GM\Sigma}{R^2} \quad (\text{Narayan+ 2003})$$

- Mass conservation in a disk:

$$\dot{M} = 2\pi R \Sigma \beta_r c$$

- BH magnetic field:

$$B_r^{\max} \sim 2 \times 10^4 \text{ [G]} (0.1/\beta_r)^{1/2}$$

$$\phi^{\max} \simeq 50 (0.1/\beta_r)^{1/2}$$

for $\dot{M} = 0.1\dot{M}_{\text{Edd}}$, $M = 10^9 M_{\odot}$

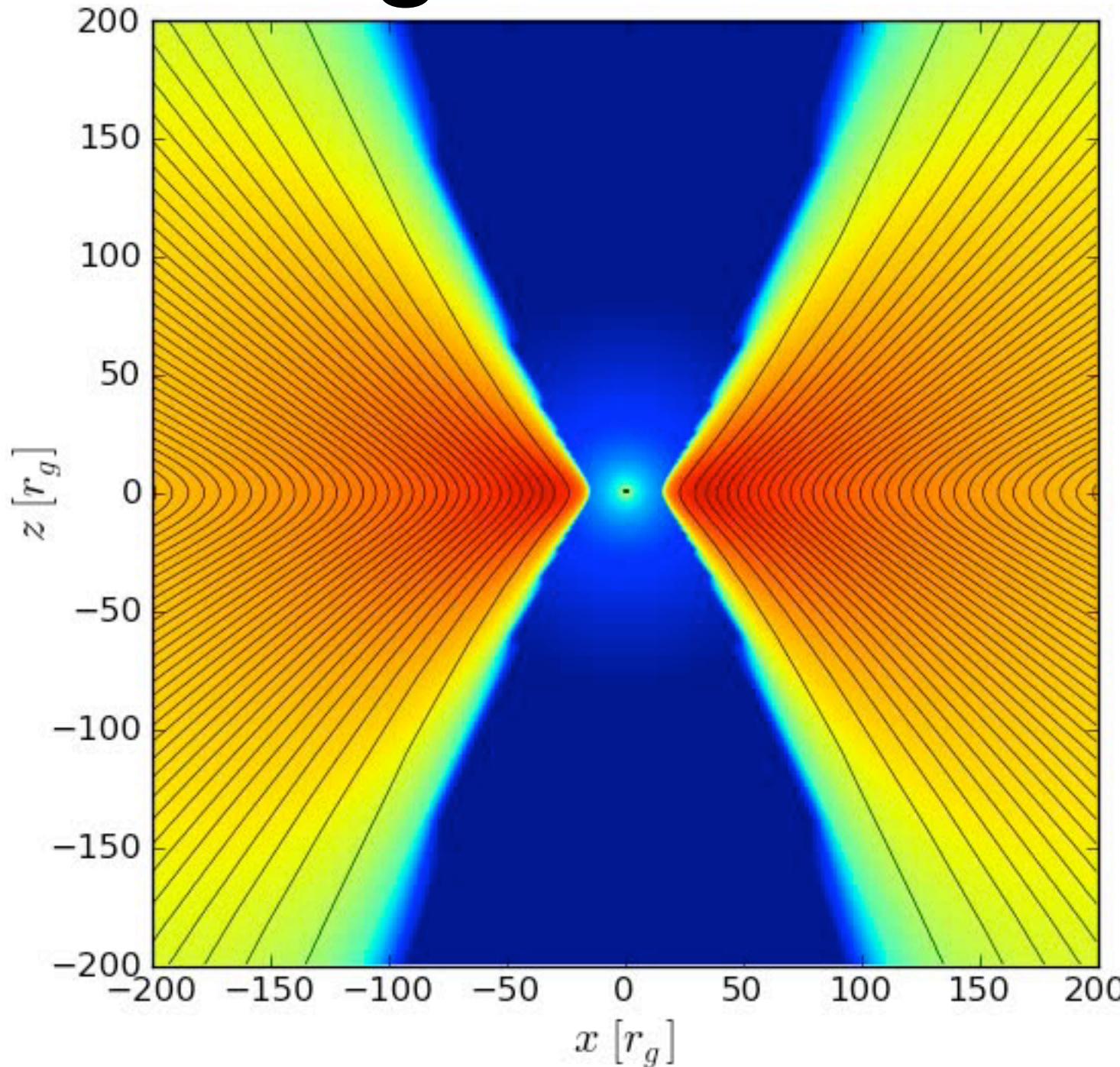
Maximum Black Hole Power = ?

- Jet power depends on magnetic field topology
(McKinney 2005, Beckwith, Hawley & Krolik 2008, McKinney & Blandford 2009)
 - Dipolar geometry gives powerful jet
 - Quadrupolar or toroidal gives weak or no jet
- GR MHD simulations give $\eta_{\text{BZ}} \lesssim 20\%$, even for nearly maximally spinning BHs (McKinney 2005, de Villiers et al. 2005, Hawley & Krolik 2006, Barkov & Baushev 2011)
- Can we obtain larger values of η ?
- Observations: some AGN have $\eta \gtrsim 100\%$
(Rawlings & Saunders 1991, Fernandes et al. 2010, Ghisellini et al. 2010, Punsly 2011, McNamara et al. 2011)

Maximum Black Hole Power = MAD Power

- Jet power increases with increasing BH magnetic flux, Φ .
- BH + large Φ = **magnetically-arrested accretion (MAD)**: BH is saturated with flux, and B-field is as strong as gravity
(Bisnovatyi-Kogan & Ruzmaikin 74, 76, Igumenshchev et al. 03, Narayan et al. 03, AT et al. 11, AT & McKinney 12a,b, McKinney, AT, Blandford 12).
 - Maximum η for each spin
 - Efficiency exceeds 100%
 - First example of net energy extraction from a BH
- New physics: high jet power, QPOs, mode of accretion...
- Advanced 3D GR MHD simulations with HARM code:
took over 2000 CPU-years! (Gammie et. al. 2003; AT et al. 2007)

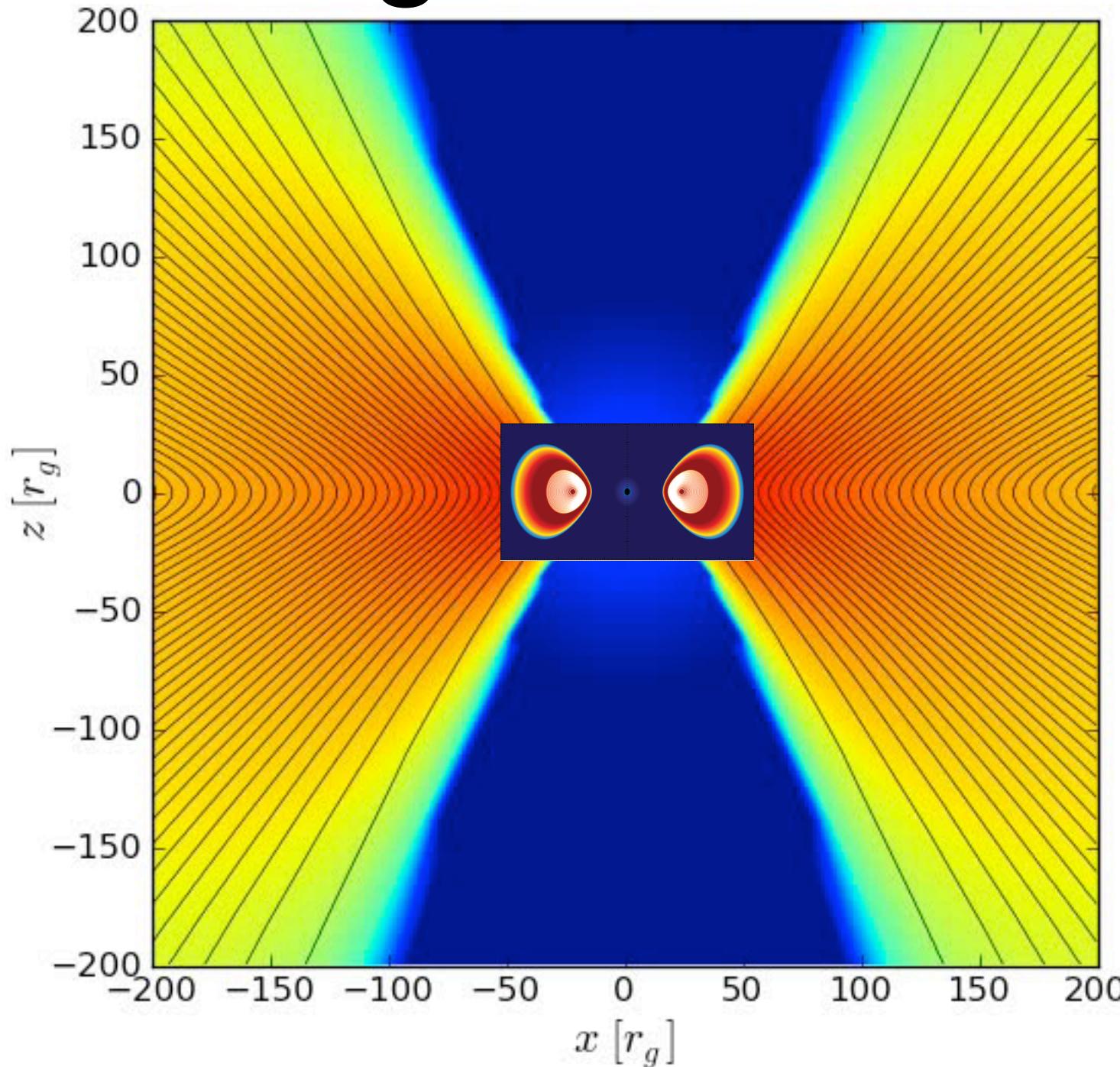
Much Larger Flux than Before



Our grid
extends
out to
 $10^5 r_g$

AT, Narayan,
McKinney 2011,
MNRAS, 478, L79

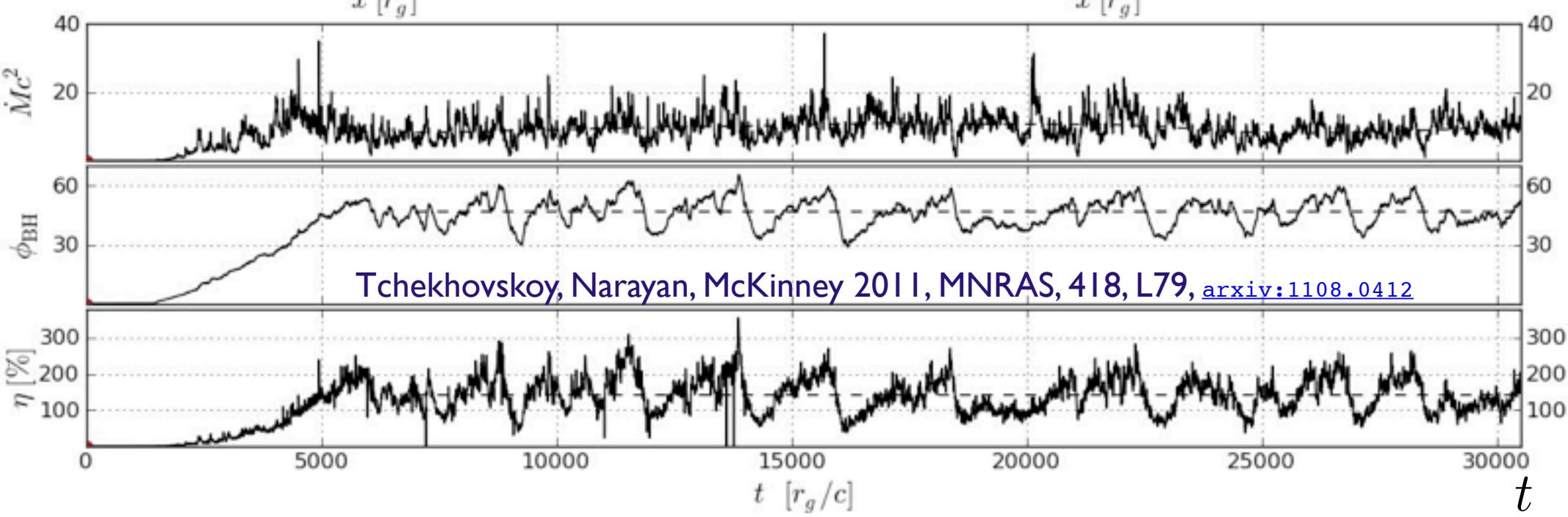
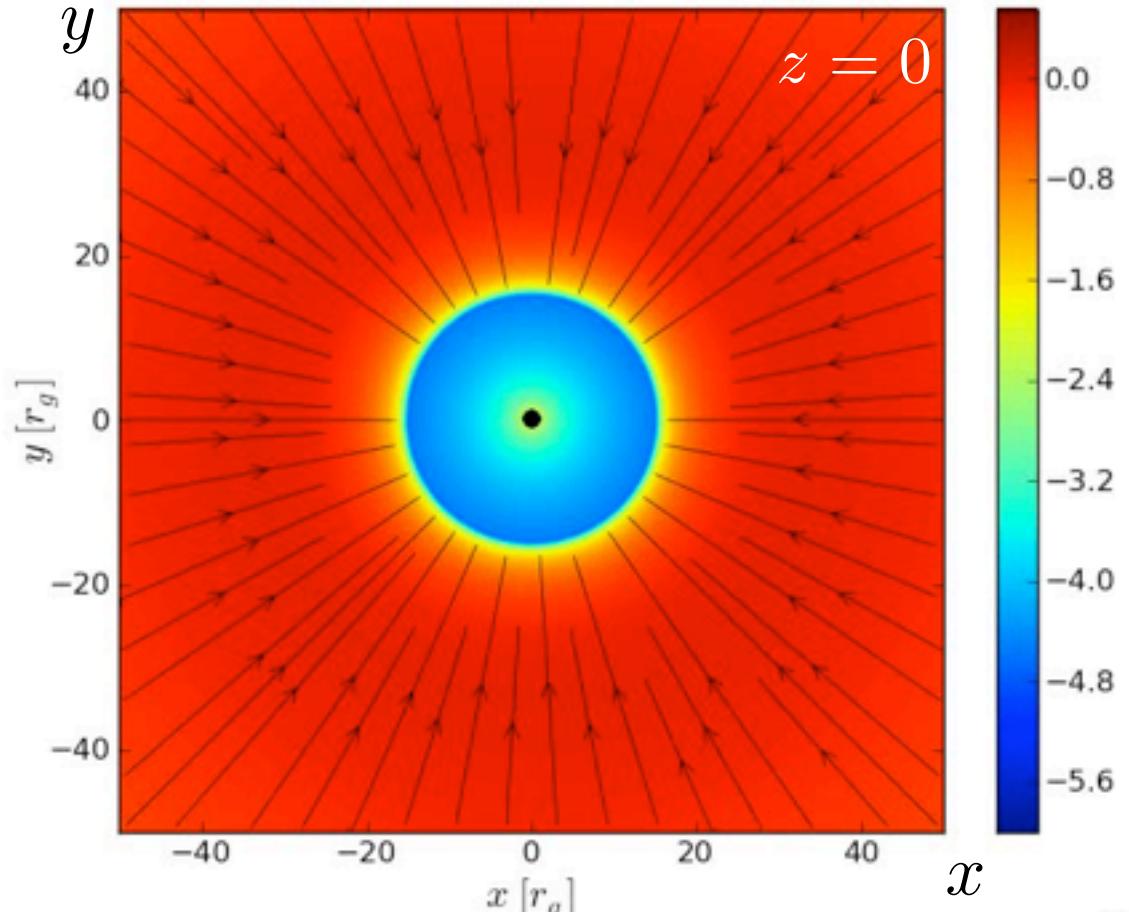
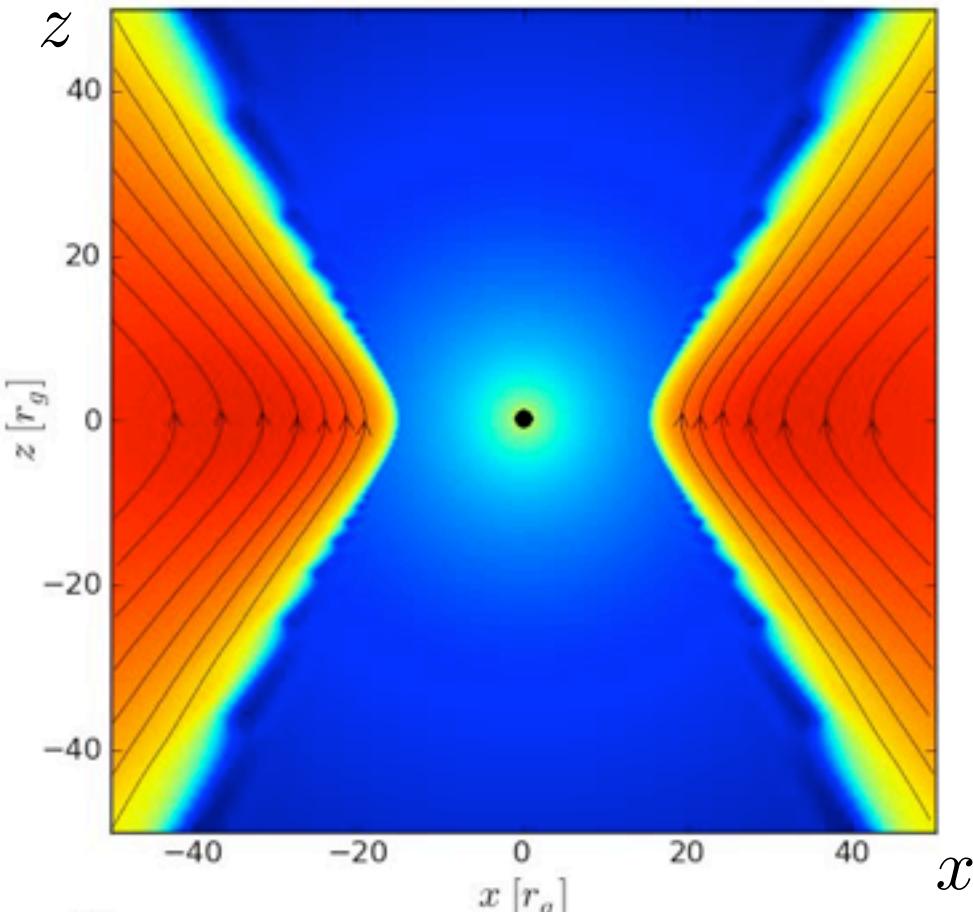
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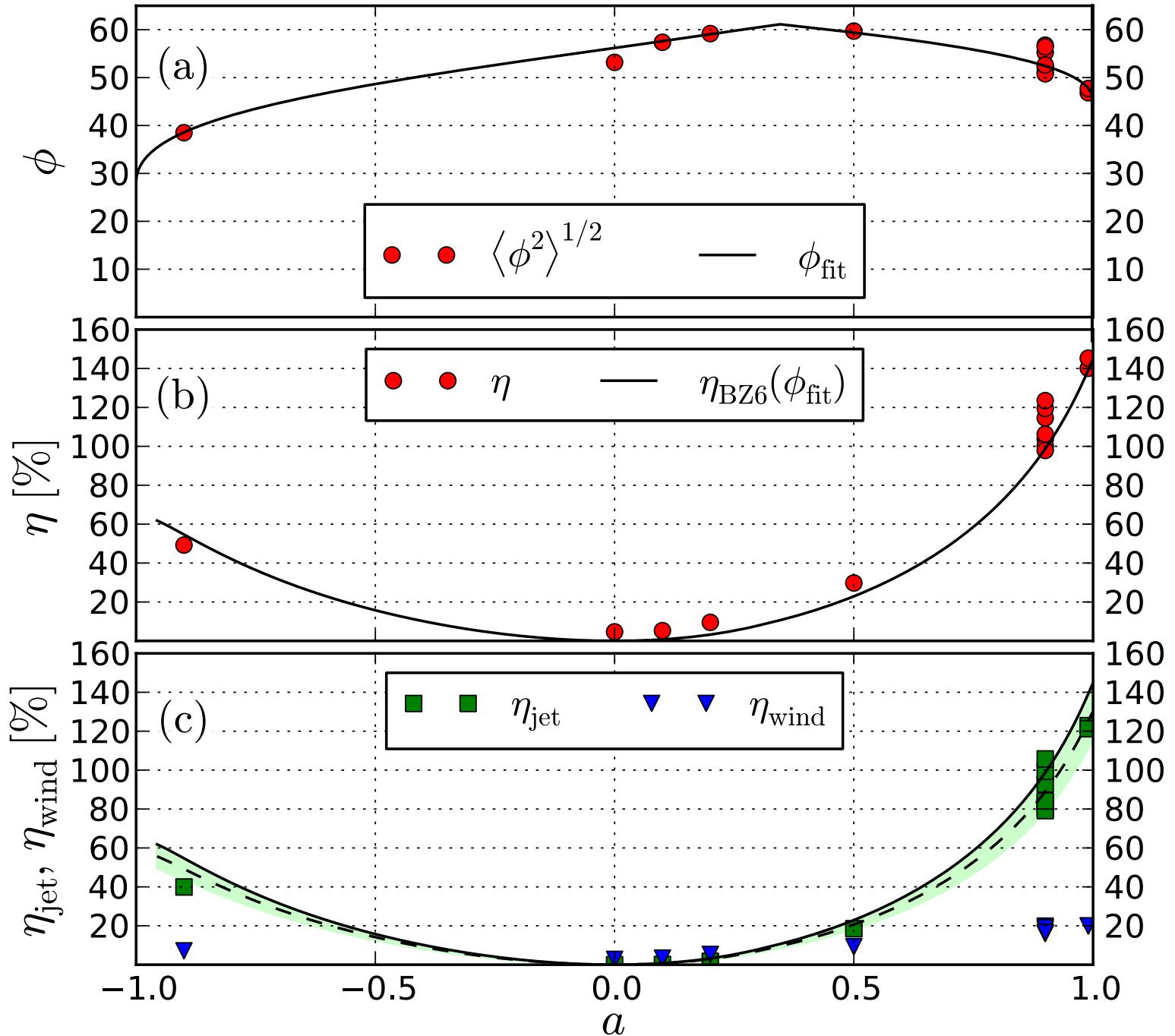
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Maximum jet power vs. spin ($h/r \sim 0.3$)

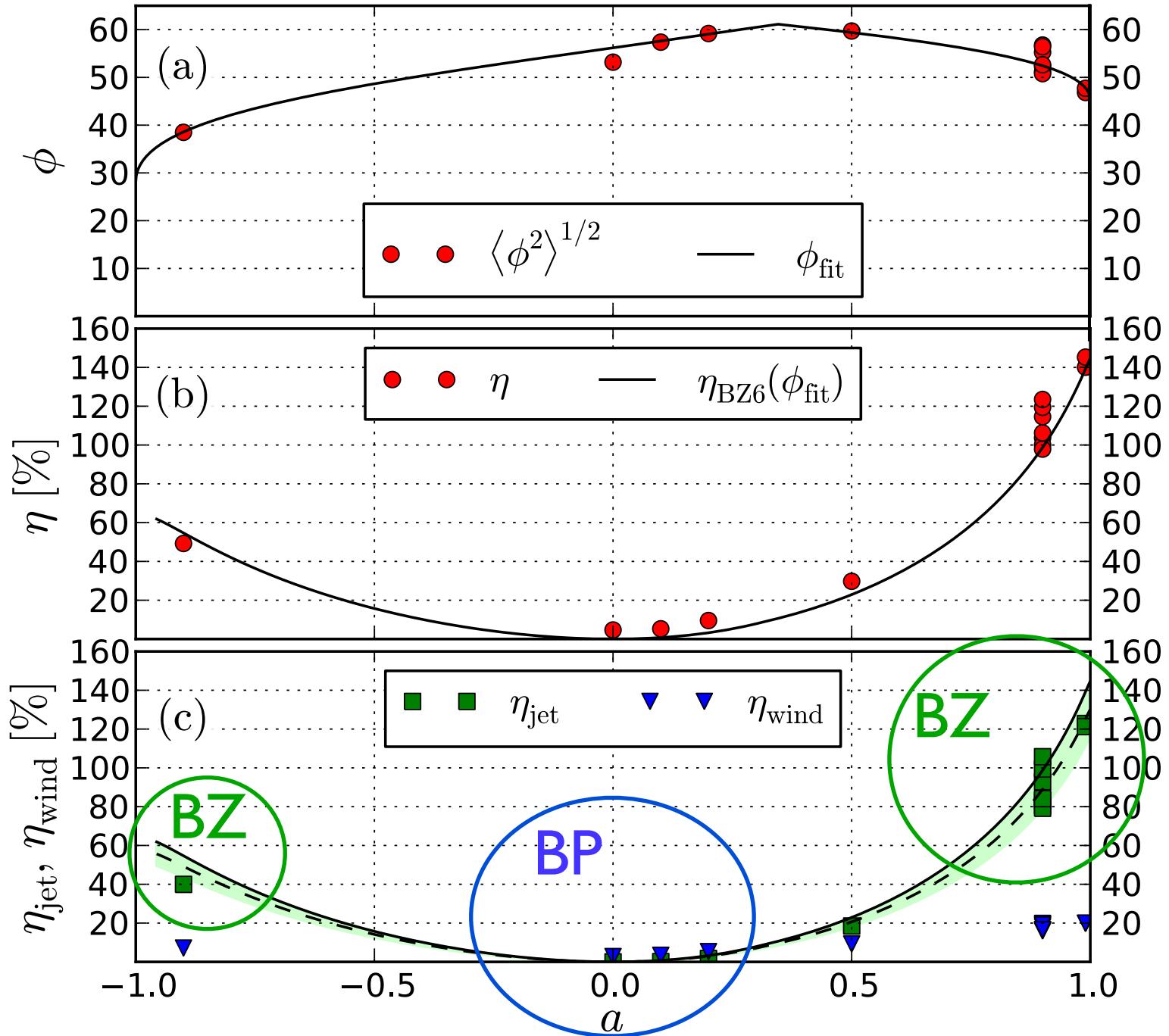
(AT, McKinney
2012a, MNRAS,
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$\eta > 100\%$ unambiguously shows that
net energy is extracted from the BH

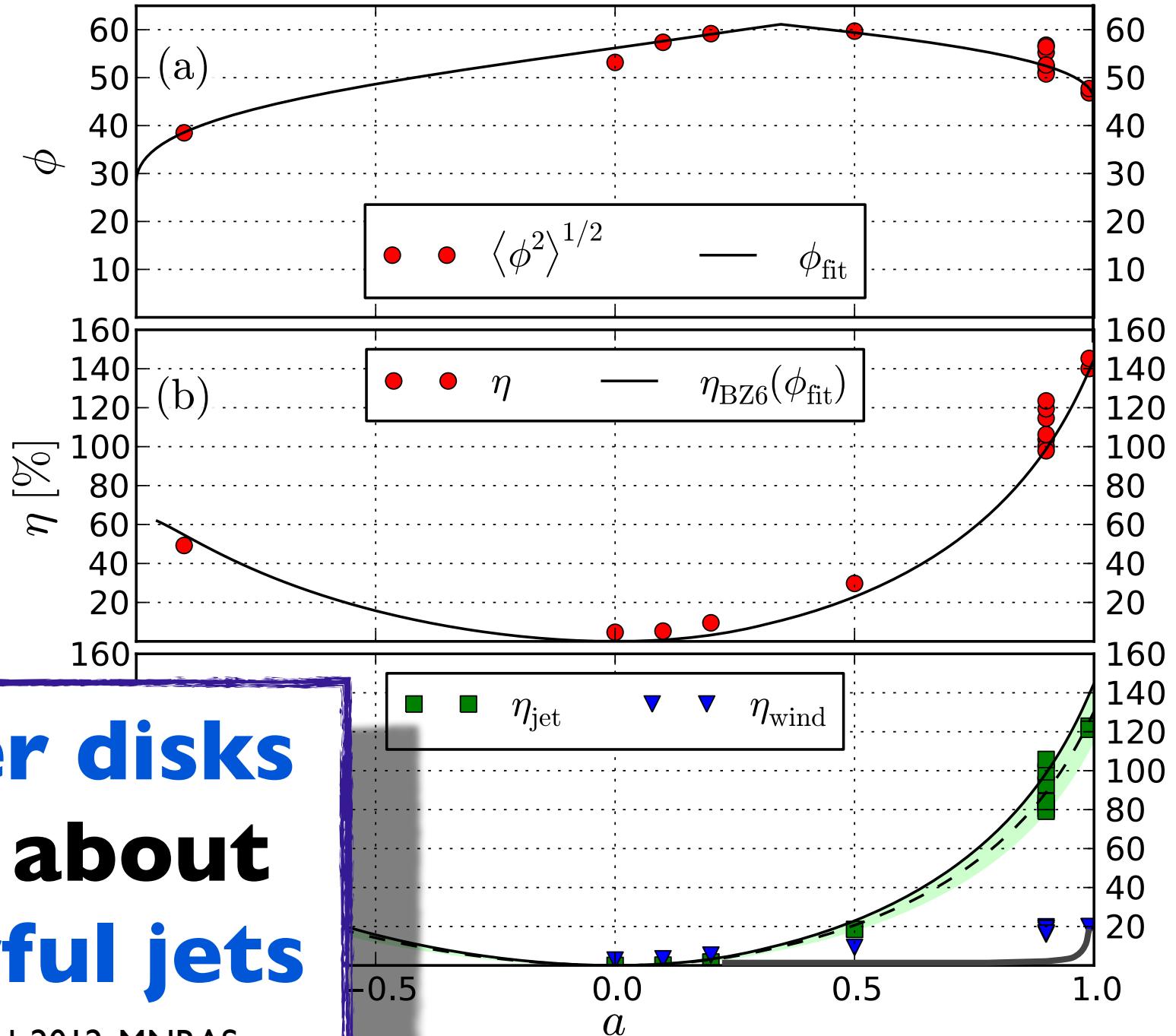
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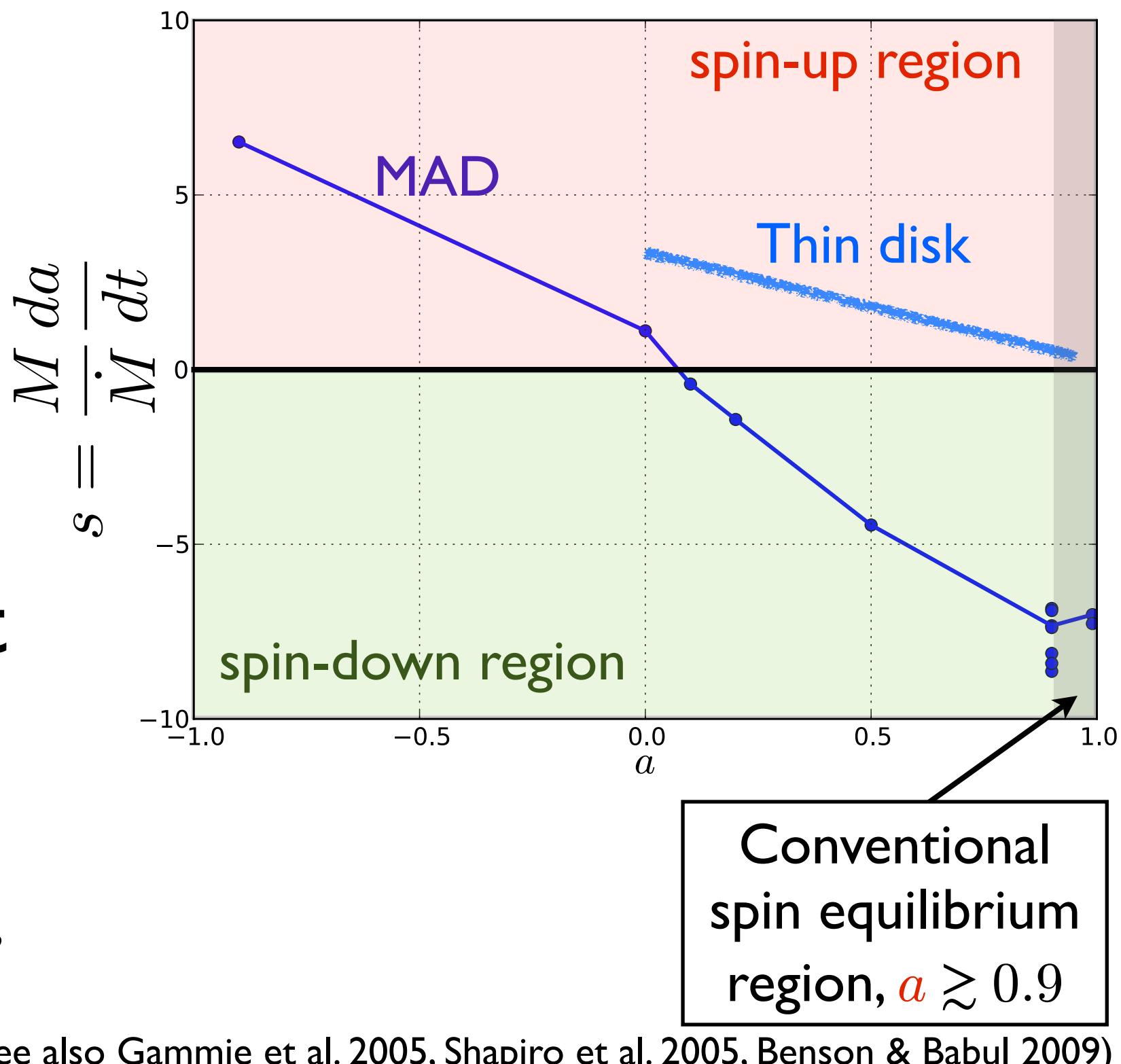
**2x thicker disks
produce about
2x powerful jets**

(McKinney, AT, Blandford, 2012, MNRAS,
submitted, arXiv:1201.4163)

% unambiguously shows that
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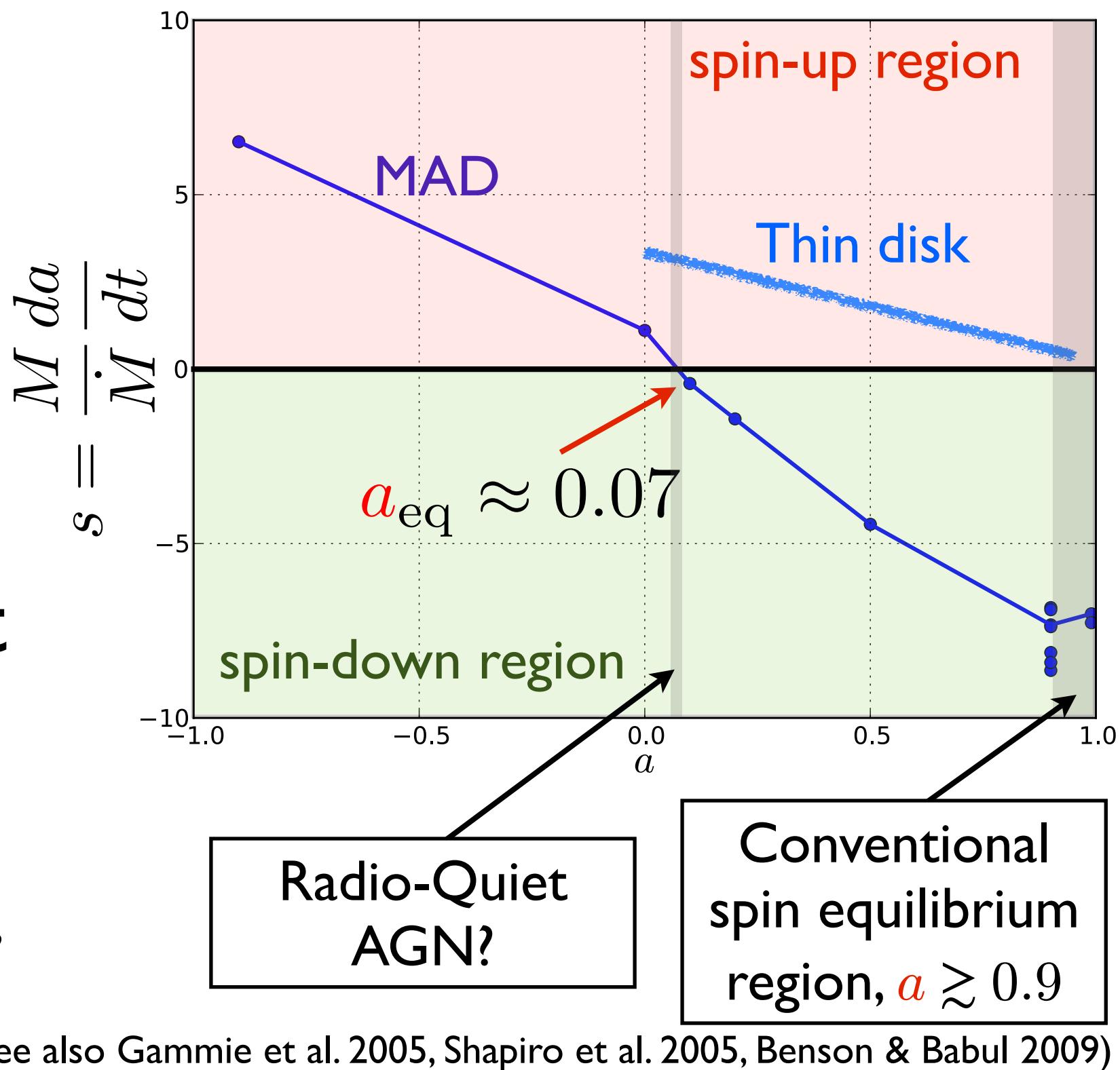
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Summary

- Simulations maximize $\Phi \rightarrow$ MAD state:
 - ▶ $\eta > 100\%$ for $a \gtrsim 0.9$
 - ▶ MADs slow BHs down to a “halt”, $a \lesssim 0.1$
- Relativistic jets can provide highly efficient feedback, efficiency up to \sim few $\times 100\%$
- Slower, mass-loaded disk winds can carry up to \sim few $\times 10\%$ -> UFOs? (Tombesi +2010,2011,2012; see talk by Cappi)