

*Normalized Intensity*

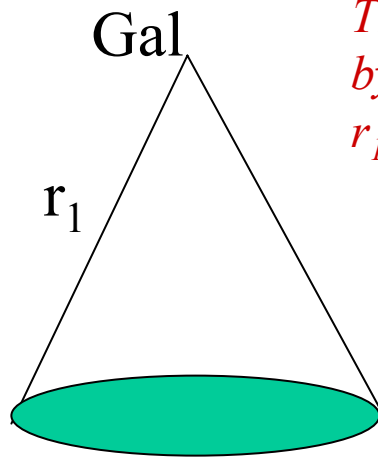
# Apparent magnitude

$$I(\lambda) d\lambda = \frac{\text{Energy}(\lambda)}{L} d\lambda$$

**Total Energy emitted at  $t_1$** 

$$E(\lambda) d\lambda \equiv dL = L I(\lambda) d\lambda$$

$$\text{or } J(\nu) d\nu = I(\lambda) d\lambda \text{ with } J(\nu) = \frac{1}{c} \lambda^2 I(\lambda)$$



*The emission at time  $t_1$  (constant) at position  $r_1$  (constant) – or by the time it reaches us it is distributed over a sphere of radius  $r_1$  so that the line element is*

$$ds^2 = -r^2 a^2(t) (d\theta^2 + \sin^2 \theta d\phi^2)$$

*The line element on the surface of a Euclidean sphere of Radius  $r a(t)$ .*

*The Area therefore is:  $4 \pi r_1^2 a^2(t_o)$  since it is the surface of a sphere  
But how much light?*

*I observe:  $\lambda_0 \quad \lambda_0 + \Delta\lambda_0 \quad \mathfrak{I}(\lambda_0)\Delta\lambda_0$*

*At departure,  $t_{em}$ ,  $\frac{\lambda_0}{1+z}$  and  $\frac{\lambda_0 + \Delta\lambda_0}{1+z}$*

*And the Energy emitted at  $t_{em}$  :*

$$L I\left(\frac{\lambda_0}{1+z}\right) \frac{\Delta\lambda_0}{1+z} \Delta t_{em}$$

$$\text{Energy per photon: } h \frac{c(1+z)}{\lambda_0}$$

$$\text{Number of photons: } L I \left( \frac{\lambda_0}{1+z} \right) \frac{\Delta \lambda_0}{1+z} \Delta t_{em} \frac{\lambda_0}{hc(1+z)}$$

$$L I \left( \frac{\lambda_0}{1+z} \right) \frac{1}{(1+z)^2} \frac{\lambda_0}{hc} \Delta \lambda_0 \Delta t_{em}$$

*The photons go over the area  $4\pi r_1^2 a^2(t)$  and are collected over a time  $\Delta t_0$  per unit area. Their Energy is  $h c/\lambda_0$  so that finally (multiplying the Number of photons by their energy) we have:*

$$\mathfrak{I}(\lambda_0) \Delta \lambda_0 = L I \left( \frac{\lambda_0}{1+z} \right) \frac{1}{(1+z)^2} \frac{1}{4\pi r_1^2 a^2(t)} \Delta \lambda_0 \Delta t_{em} \frac{1}{\Delta t_0}$$

$$\text{and as we have shown that: } \Delta t_{em} \frac{1}{\Delta t_0} = \frac{1}{(1+z)}$$

$$\mathfrak{I}(\lambda_0) = L I \left( \frac{\lambda_0}{1+z} \right) \frac{1}{(1+z)^3} \frac{1}{4\pi r_1^2 a^2(t)}$$

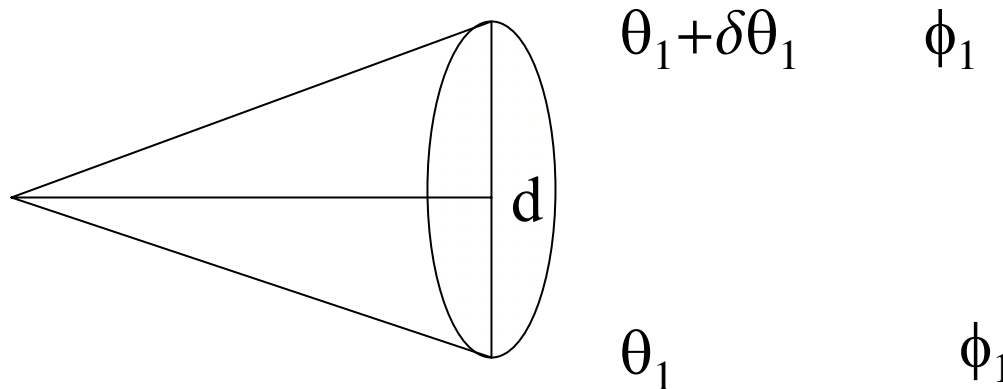
*Bolometric Luminosity and Bolometric Flux:*

$$\mathfrak{I}(\nu_0) = \frac{L J(\nu_0 \cdot (1+z))}{(1+z) 4\pi r_1^2 a^2(t_0)}$$

$$\begin{aligned} \mathfrak{I}_{bol} &= L_{em} \int_0^\infty I\left(\frac{\lambda_0}{1+z}\right) \frac{1}{(1+z)^3} \frac{1}{4\pi r_1^2 a^2(t)} d\lambda = \\ &= L_{em} \frac{1}{(1+z)^3} \frac{1}{4\pi r_1^2 a^2(t)} \int_0^\infty I\left(\frac{\lambda_0}{1+z}\right) d\left(\frac{\lambda_0}{1+z}\right) * (1+z) \\ \mathfrak{I}_{bol} &= \frac{L}{(1+z)^2 4\pi r_1^2 a^2(t_0)} \end{aligned}$$

# Angular size

- We have seen we are capable of measuring magnitudes only if we know  $r_1 a(t_0)$ .



*I choose a system of coordinates so that I measure the angle between two points in the coordinate  $\theta$ . Since I am talking about a proper separation in the RW metric I have:  $\Delta t=0$ ,  $\Delta r=0$ ,  $\Delta \phi=0$  and  $\Delta \theta \neq 0$ . That is I can write:*

$$ds^2 = -r_1^2 a^2(t_1) \Delta \theta^2 = -d^2$$

## Continue

On the other hand I measure at the time  $t_0$  and I will have:

$$d = r_1 a(t_1) (\Delta\theta_1)^2$$
$$\Delta\theta_1 = \frac{d}{r_1 a(t_1)} = \frac{d(1+z)}{r_1 a(t_0)}$$

*We need a Model to estimate  $r_1 a(t_0)$ . Note also that while  $r_1$  increases with  $z$  (I am looking further away)  $a(t_0)$  decreases.*